



Spin-up of UK Earth System Model 1 (UKESM1) for CMIP6

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Key Points:

- Earth system components and spin-up protocol of UKESM1 for CMIP6 outlined
- Ocean-only (5000y) and Land-only (1000y) phases used prior to fully-coupled finalising of spin-up (500y)
- Evaluation of spin-up protocol presented, including cross-component validation of piControl state and drift

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Abstract

For simulations intended to study the influence of anthropogenic forcing on climate, temporal stability of the Earth's natural heat, freshwater and biogeochemical budgets is critical. Achieving such coupled model equilibration is scientifically and computationally challenging. We describe the protocol used to spin-up the UK Earth system model (UKESM1) with respect to pre-industrial forcing for use in the 6th Coupled Model Intercomparison Project (CMIP6). Due to the high computational cost of UKESM1's atmospheric model, especially when running with interactive full chemistry and aerosols, spin-up primarily used parallel configurations using only ocean/land components. For the ocean, the resulting spin-up permitted the carbon and heat contents of the ocean's full volume to approach equilibrium over ~5000 years. On land, a spin-up of ~1000 years brought UKESM1's dynamic vegetation and soil carbon reservoirs towards near-equilibrium. The end-states of these parallel ocean- and land-only phases then initialised a multi-centennial period of spin-up with the full Earth system model, prior to this simulation continuing as the UKESM1 CMIP6 pre-industrial control (piControl). The realism of the fully-coupled spin-up was assessed for a range of ocean and land properties, as was the degree of equilibration for key variables. Lessons drawn include the importance of consistent interface physics across ocean- and land-only models and the coupled (parent) model, the extreme simulation duration required to approach equilibration targets, and the occurrence of significant regional land carbon drifts despite global-scale equilibration. Overall, the UKESM1 spin-up underscores the expense involved and argues in favour of future development of more efficient spin-up techniques.

40 **Plain Language Summary**

41 Earth system models (ESMs) are an important tool for understanding the Earth and
42 for projecting how climate change may affect natural and human systems. For simulations
43 of ESMs to separate anthropogenic influences on climate from the background state, the
44 stability of the unperturbed system is critical. However, achieving this equilibrium is both
45 scientifically and computationally challenging. Here, we describe how this was achieved
46 for one such model, UKESM1, for the 6th Coupled Model Intercomparison Project (CMIP6).
47 Due to the cost of the full model, especially when running with atmospheric chemistry
48 and aerosols, much of UKESM1's spin-up to equilibrium made use of ocean- and land-
49 only configurations. Millennial-scale spin-up phases of these component-only models were
50 used to initialise a final centennial-scale phase of the full model to reach pre-industrial
51 equilibrium targets. The stability and realism of UKESM1's spun-up state was then evalu-
52 ated across a broad range of properties. A number of lessons were drawn from this spin-
53 up including the extreme simulation duration required to reach equilibrium. A key conclu-
54 sion is the importance of developing efficient techniques to spin-up ESMs.

1 Introduction

To a first approximation, the behaviour of the Earth system (ES) is governed by the dynamics and interactions of the two geophysical fluids – the atmosphere and the ocean – that comprise the majority of the planet’s surface substrate. Despite a number of similarities, these two fluids diverge in many other respects, including a critical difference in the timescales of their internal dynamics. Features in the atmosphere form and dissipate over periods typically of the order of hours, days or weeks in duration, with a residence time for one of its most dynamic components – water – of only 8.9 days [van der Ent and Tuinenburg, 2017]. In contrast, while the ocean’s surface readily exchanges and interacts with the atmosphere over short timescales, its interior is structured by a vast thermohaline circulation that sluggishly transports water around its basins and into the abyssal deep. Depending upon its location, such water leaving contact with the atmosphere can take decades, centuries or even millennia to overturn completely and come back into contact with the atmosphere. For example, estimated from radiocarbon and from inverse models, the waters of the deep North Pacific have a ventilation age of 1200–1500 years (Gebbie and Huybers [2012]; Khatiwala et al. [2012]), with some model studies suggesting much longer timescales [Wunsch and Heimbach, 2008].

Consequently, with a ventilation timescale of more than a millennium, the ocean component of the Earth system has a long memory – one that can “remember” environmental perturbations far longer than other components such as the atmosphere and land surface [Ciais et al., 2013]. In addition, the ocean is the largest active reservoir of carbon in the Earth system, approximately 40000 petagrams carbon (Pg C) [Ciais et al., 2013]. Relative to the atmosphere – where the concentration of carbon dioxide (CO₂) has been of long-standing interest – this represents a store more than 50 times greater [Ciais et al., 2013]. A consequence of this is that even small imbalances in the air-sea exchange of CO₂ can lead to large changes in atmospheric CO₂ [Kwon et al., 2009]. Furthermore, biogeochemical processes within the ocean, such as those of the biological pump [Raven and Falkowski, 1999], can significantly alter seawater chemical composition, with implications for the wider carbon system when deep water parcels finally re-establish contact with the atmosphere.

The land system represents another significant store of carbon in the Earth system. On land, carbon is stored both in living biomass and in soil as decaying organic carbon.

87 Relative to the ocean, the total land reservoir is comparatively small (approximately 2200
88 Pg C vs. 40000 Pg C; Ciais et al. [2013]). Nonetheless, the more rapid timescale of in-
89 teraction between the land and atmosphere leads to the terrestrial carbon content being
90 strongly influenced by climate variability. Turnover timescales in the living and decaying
91 pools of carbon mean that, like the ocean, equilibration of the land system requires ex-
92 tended periods of model spin-up.

93 These significant reservoirs of carbon, and their relatively slow turnover times, whether
94 through sluggish ventilation or gradual decay processes, have important implications for
95 simulations of Earth system models (ESMs) aimed at studying the influence of human
96 perturbations on the system. Principally, in fully-coupled ESMs, where both the climate
97 and CO₂ are free to evolve, to robustly detect human perturbations requires the ocean and
98 land carbon stores be in temporal equilibrium before any human forcing is imparted. If
99 this temporal stability is not achieved then the slow equilibration trend of either carbon
100 reservoir could be confused with, and even influence, any human-induced trend, confound-
101 ing the detection of human forcing of the system. For instance, in a model with natural
102 land or ocean carbon pools outgassing, such drift will mask ingassing fluxes driven by the
103 steady accumulation of anthropogenic CO₂ in that atmosphere.

104 Separate to its carbon reservoir, ocean spin-up also serves to bring its physical state,
105 particularly ocean heat content, as well as the biogeochemical cycles of other elements,
106 into equilibrium. On land, spin-up serves to bring the various vegetation types into bal-
107 ance with their local climate (temperature, water and nutrient availability, etc.) and, through
108 ecological competition, with each other.

109 The desirability of a well-equilibrated ESM is typically offset by the computational
110 cost of achieving this. While most experimental simulations may only be years, decades
111 or centuries in duration, full spin-up typically requires of order one to tens of millennia
112 of simulation. In the case of the ocean, on top of the estimated ventilation timescales of
113 the ocean's "oldest" watermasses [Khatiwala et al., 2012], spin-up must further account for
114 biogeochemical "shuffling" of nutrients, such as the downstream effects of a model's bio-
115 logical pump on the nutrient concentrations of its deep, and then upwelling, watermasses.
116 Ocean physical properties are similarly affected, with the distinction that, in gradually
117 changing the seawater bulk properties, spin-up also alters the ocean density and potential
118 energy field, with consequences for the very circulation that is spinning everything up.

119 Consequently, techniques for model spin-up are diverse, reflecting spin-up aspira-
120 tions (i.e. physics or biogeochemistry or both), model cost (i.e. computational or wall-
121 clock or both) and the availability of suitable engineering solutions (i.e. does model code
122 permit particular accelerated modes of spin-up).

123 The most conventional, and arguably “best” approach, is simply to run the model for
124 a long period of time (simulated and wallclock). This ensures that spin-up is consistent
125 with the normal model operation, and avoids introducing any artifacts caused by spin-up
126 “shortcuts”. With continual advances in the power and availability of computational re-
127 sources, this approach should become less burdensome with time, with past models be-
128 coming easier to spin-up to equilibrium. However, our parallel increase in knowledge and
129 understanding favours increasingly well-resolved and more complex models, whose aspi-
130 rations foster a “Red Queen” effect within the modelling community (with some notable
131 exceptions; Cui et al. [2011]). That is, while computational gains should permit faster
132 spin-up, they actually favour increased realism, with the result that spin-up remains com-
133 putationally expensive despite these gains. Consequently, this “brute force” approach to
134 spin-up remains tantalisingly out of reach for state-of-the-art ESMs.

135 While ocean ventilation in the Earth system is relatively sluggish, ocean models are
136 usually computationally faster than their atmosphere counterparts, to which they are cou-
137 pled. Resolution may be comparable (as in UKESM1), but the absence of detailed radia-
138 tion schemes, typically fewer advected tracers, and automatically fewer grid cells because
139 of the occurrence of land, means ocean-only models typically exhibit greater wallclock
140 efficiency. Consequently, one spin-up approach called “online ocean-only” is to run a de-
141 coupled ocean component with appropriate surface boundary conditions, and simulate the
142 majority of ocean equilibration without the more expensive atmosphere. This approach is
143 facilitated by the atmosphere’s relatively rapid equilibration, such that it can readily both
144 provide surface forcing, and be “re-attached” to the ocean for a comparatively brief period
145 at the end of spin-up.

146 The online ocean-only approach also extends to land spin-up which, like the ocean,
147 can include elements (e.g. soil carbon) that require extended simulation periods. Much
148 like the ocean, the land can be driven by atmospheric forcing at its boundary, sparing the
149 cost of full atmospheric simulation. It differs in that the modelled system typically has re-
150 duced vertical resolution, and its prognostic variables (carbon reservoirs, vegetation types)

151 are not advected. As such, while the three-dimensional nature of the ocean means that it
152 remains moderately expensive to run even on its own, the reduced number of levels and
153 lack of advection means the land spin-up is computationally less of a burden.

154 A further approach, “offline ocean-only”, separates the spin-up of ocean physics
155 from that of biogeochemistry, by treating ocean circulation itself as just another part of the
156 forcing. Once circulation has stabilised, either in full Earth system or ocean-only mode, it
157 is used as a three-dimensional climatology to transport tracers of marine biogeochemistry.
158 In this way, the subsequent cost of calculating ocean physics is avoided, permitting a more
159 computationally-efficient spin-up. A superficially similar approach to “offline ocean-only”
160 is that of the transport matrix method (TMM; Khatiwala et al. [2005]; Khatiwala [2007]).
161 Rather than explicitly using a stored model circulation to drive biogeochemical tracers,
162 this method describes the spatial connectivity driven by ocean circulation as a sparse ma-
163 trix that can efficiently be used as a transport operator. While both of these approaches
164 serve to spin-up passive tracers at a somewhat reduced computational cost, both still re-
165 quire an equilibrium physical circulation in the first place, which in turn requires its own
166 spin-up. As we need to spin-up both the physical circulation and ocean biogeochemistry
167 of UKESM1, our ocean spin-up here makes use instead of the “online ocean-only” to do
168 both in tandem.

169 Note that the discussion above effectively assumes equilibration is always for the
170 good, essentially because of the resulting reduction in model drift. However, as imperfect
171 tools, models do not necessarily converge towards a state similar to that of the real Earth
172 system, and extended spin-up is liable to produce a divergent state relative to the true ob-
173 served state (while revealing model biases). Paradoxically, by reducing model drift while
174 increasing model bias, equilibration can seemingly reduce a model’s skill when evaluated
175 against observations [Séférian et al., 2016]. Conversely, by limiting spin-up, a model will
176 diverge less from its (typically) observationally-derived initial state, and its state will show
177 smaller biases (if greater drift). Nonetheless, the need for a stable control simulation from
178 which to initialise historical runs (and then future projections) is more important than such
179 considerations. Not least because the drift from an observation-based initial condition is
180 likely larger, over the first few hundred years of a simulation, than the anthropogenic sig-
181 nal we wish to detect.

182 Within the Coupled Model Intercomparison Project Phase 6 (CMIP6), the Diag-
183 nostic, Evaluation and Characterization of Klima (DECK) protocol [Eyring et al., 2016]
184 describes the baseline simulations that all participating models must undertake to “bench-
185 mark” their performance. An underpinning part of the DECK is the production of a pre-
186 industrial control (piControl) simulation from which model states can be drawn to ini-
187 tialise simulations for both the DECK and other Model Intercomparison Projects (MIPs).
188 While the DECK outlines certain boundary conditions for this piControl (e.g. atmospheric
189 CO₂ concentrations, orbital parameters, a mean solar cycle, etc.), it does not specify a par-
190 ticular methodology or duration for the production of this model state. This stems partly
191 from the variety of models participating in CMIP, and the resulting difficulty in defining
192 universal criteria for models that range widely in complexity, resolution and degree of in-
193 ternal coupling. Additionally, the potential computational cost of spin-up is a factor, with
194 participating groups varying in their access to compute resources. Some MIP protocols,
195 such as C4MIP [Jones et al., 2016], suggest equilibrium criteria for participating models,
196 but the DECK requirement of a multi-century piControl to shadow MIP simulations is
197 intended as a means to quantify (and control for) drift in CMIP6 simulations. This situa-
198 tion largely repeats that of CMIP5, where total spin-up durations varied widely from only
199 200 years up to almost 12,000 years [Séférian et al., 2016]. As well as this wide span of
200 spin-up durations, the CMIP5 models summarised by Séférian et al. [2016] also varied
201 widely in the spin-up methodology used. Models adopted various offline, accelerated of-
202 fline and component-only online approaches, often in unique combinations, prior to final
203 periods of fully-coupled simulation. However, in the absence of formalised guidance or
204 commonly-accepted spin-up procedures, the documentation of spin-up typically remains a
205 lower-priority activity. Nonetheless, a number of studies have examined aspects of spin-up,
206 such as how specifically to equilibrate (“spin-down”) from modern initial conditions to the
207 preindustrial state [Stouffer et al., 2004], quantifying the sources of drift or variability in
208 spun-up models Doney et al. [2006], and how drifts can be corrected without introducing
209 bias [Hobbs et al., 2016]. Furthermore, an increasing number of studies document the ap-
210 proaches used to spin-up ESMs (more comprehensive examples include: Watanabe et al.
211 [2011]; Séférian et al. [2013]; Lindsay et al. [2014]).

212 Here we document the spin-up procedure followed in preparing a pre-industrial con-
213 trol state of the new ESM, UKESM1, for the CMIP6 DECK and MIP experiments. The
214 manuscript begins with a brief description of UKESM1 and its main components, fol-

215 lowed by an extensive description of the procedure employed to equilibrate UKESM1 to
216 CMIP6 pre-industrial forcing. We then show the evolution of the model's state during
217 spin-up, from both the parallel ocean- and land-only spin-up activities, followed by the
218 final, fully-coupled model. The model's degree of equilibration and biases in its final state
219 are discussed, together with potential future avenues for addressing these. In addition to
220 the results presented in the main body of this manuscript, supplementary material includes
221 additional tables and figures to document the spin-up and performance of UKESM1.

222 **2 Methods**

223 **2.1 UKESM1 description**

224 UKESM1 is a new state-of-the-art ESM comprised of components that represent
225 both physical and biogeochemical aspects of the Earth's atmosphere, ocean, cryosphere
226 and land systems. It is built on the recent Hadley Centre Global Environment Model ver-
227 sion 3 Global Coupled (GC) climate configuration, HadGEM3 GC3.1 (Williams et al.
228 [2017]; Kuhlbrodt et al. [2018]). This physical core model is extended through the addi-
229 tion of ocean and land biogeochemistry, and interactive stratospheric-tropospheric trace
230 gas chemistry, which predicts atmospheric oxidant fields as input to the aerosol model as
231 well as a range of radiatively active gases (e.g. O₃, CH₄, N₂O). As well as including dy-
232 namics internal to their components, these Earth system additions couple where it is be-
233 lieved that they potentially feedback upon one another (either negative and damping, or
234 positive and amplifying), or where they impact the time-evolution of the physical climate
235 system. For example, atmospheric aerosols play a key role in mediating the transfer or
236 absorption of radiation within the atmosphere, and their occurrence and behaviour is an
237 outcome of interactions between chemical and physical processes in the atmosphere, ocean
238 and ice (Halloran et al. [2010]; Carslaw et al. [2010]; Quinn and Bates [2011]; Myhre
239 et al. [2013]; Kok et al. [2018]). Representing and understanding the nature of such link-
240 ages between components is of critical importance if models are to accurately represent
241 the true Earth system sensitivity to anthropogenic forcing.

242 Figure 1 presents a schematic diagram of the components included within UKESM1,
243 together with an indication of the nature of the coupling between them. In outline, at-
244 mosphere and land components are closely coupled together as a single, integrated exe-
245 cutable, use a common grid and time-step, and communicate their states directly at each
246 time-step without the need for a coupler. The same is true for the three ocean components
247 – dynamics, sea-ice and biogeochemistry – which are also coupled together as a single
248 executable. The two executables – atmosphere-land and ocean-ice-biogeochemistry – com-
249 municate once every 3 hours through interface layers, labelled OASIS3-MCT_3.0 (Valcke
250 [2013]; Craig et al. [2017]) in Figure 1.

251 The atmosphere of UKESM1 as represented by GA7.1 represents the physical dy-
252 namics of the atmosphere, including processes such as mass transport, radiative transfer,
253 thermodynamics and the water cycle. Coupled to the GA7.1 is the UK Chemistry and

254 Aerosols model (UKCA; Morgenstern et al. [2009]; O'Connor et al. [2014]), which rep-
255 represents stratospheric and tropospheric chemistry, as well as aerosols via the GLOMAP-
256 mode scheme [Mann et al., 2010], with dust represented by a binned scheme [Woodward,
257 2011]. UKESM1 differs from GA7.1 in its treatment of the natural emissions of monoter-
258 penes, dimethyl sulphide (DMS) and primary marine organic aerosols (PMOAs), which
259 are interactively calculated from elements of the land and ocean components, permitting
260 feedbacks between the biosphere and aerosol / cloud-radiative behaviour in UKESM1. A
261 further coupling that uniquely links the land to the ocean in UKESM1 is the production
262 of wind-borne mineral dust as a function of simulated climate and bare soil on land, and
263 which can fuel ocean productivity (and uptake of CO₂) by supplying bioavailable iron.
264 See Mulcahy et al. [2018], Sellar et al. [2019] and Archibald et al. [2020] for further de-
265 tails concerning atmospheric chemistry in UKESM1.

266 The physical ocean component of UKESM1 is represented by the Nucleus for Eu-
267 ropean Modelling of the Ocean model (NEMO; Madec et al. [2016]) for its dynamical
268 circulation, and by the Los Alamos sea-ice model (CICE; Rae et al. [2015]) for its marine
269 cryosphere. More complete descriptions of the NEMO and CICE configuration used in
270 UKESM1, including details of its sensitivity and resulting tuning, can be found in Storkey
271 et al. [2018], Ridley et al. [2018] and Kuhlbrodt et al. [2018]. Marine biogeochemistry is
272 represented by the Model of Ecosystem Dynamics, nutrient Utilisation, Sequestration and
273 Acidification (MEDUSA-2.1), which includes the cycles of nitrogen, silicon, iron, carbon
274 and oxygen. The version used in UKESM1 is identified as MEDUSA-2.1, to distinguish
275 it from its earlier parent model, MEDUSA-2, described in Yool et al. [2013]. Develop-
276 ments made for UKESM1 include: 1. replacement of its carbonate chemistry with the
277 MOCSY-2.0 scheme of Orr and Epitalon [2015]; 2. the addition of an empirical submodel
278 of surface seawater DMS concentration [Anderson et al., 2001]; 3. various code improve-
279 ments including adaptations for variable volume (VVL) and upgrading to utilise the XML
280 Input-Output Server (XIOS) adopted by NEMO [Meurdesoif, 2013].

281 The land component of UKESM1 is represented by the Joint UK Land Environment
282 Simulator (JULES; Best et al. [2011]; Clark et al. [2011]), which handles physics and in-
283 tegrated biogeochemistry. This is closely coupled with the Top-down Representation of
284 Interactive Foliage and Flora Including Dynamics model (TRIFFID; Cox [2001]; Jones
285 et al. [2011]), a dynamic global vegetation model that represents plant and soil dynamics
286 on land. Developments since CMIP5 include: 1. updating of plant growth and turnover

287 parameters to reflect the plant trait database, TRY [Kattge et al., 2011]; 2. an increase in
288 the number of plant functional types (PFTs) from five to thirteen, further permitting the
289 distinction of evergreen / deciduous plants and tropical / temperate evergreen trees [Harper
290 et al., 2016]; 3. the emission of volatile organic compounds (VOCs; e.g. Pacifico et al.
291 [2015]); 4. limitation on terrestrial primary production (and therefore CO₂ uptake) through
292 the availability of soil and plant nitrogen. Land-use by agriculture is represented in TRIF-
293 FID by reserving fractions of each grid cell for crops and pasture, with these fractions
294 prescribed as external forcing that can vary with time. The Greenland and Antarctic land
295 icesheets are represented via a sub-gridscale scheme described in Shannon et al. [2019].
296 For further details of UKESM1's land component, please refer to Sellar et al. [2019].

297 By default, UKESM1 has a relatively coarse horizontal resolution of N96 (approx-
298 imately 135 km) in the atmosphere and 1° (approximately 73 km) in the ocean. Verti-
299 cal resolution is 85 levels in the atmosphere (with a model top at 85 km), and 75 levels
300 in the ocean (with a maximum depth of 6 km), with, in both cases, high vertical res-
301 olution focused at the interface between the two fluids. This resolution corresponds to
302 the HadGEM3 N96ORCA1 configuration, a full description of which can be found in
303 Kuhlbrodt et al. [2018].

304 In the work described here, the fully coupled version of UKESM1 was only utilised
305 for a restricted (latter) portion of the full spin-up process. This was in part because of
306 its high computational cost, but also because this cost is largely associated with atmo-
307 spheric components that spin-up to equilibrium much more quickly than the ocean or
308 the land. The majority of the spin-up was therefore performed using parallel ocean-only
309 and land-only versions of UKESM1, forced at their surface boundary conditions by at-
310 mospheric output from a shorter coupled model simulation. More complete details of the
311 fully-coupled UKESM1, including an analysis of its pre-industrial and historical climate,
312 can be found in Sellar et al. [2019].

313 **2.2 UKESM1 spin-up**

314 Spin-up of UKESM1 utilised a combination of phases using coupled climate, ocean-
315 only, land-only and full Earth system coupled versions of the model (with and without
316 interactive atmospheric chemistry). The development cycle of the full model occurred in
317 parallel with spin-up activities, with the result that spin-up did not use a single version

318 of the model throughout. Periodically, model improvements and bugfixes were applied
319 between model phases. The last segment of spin-up did employ the final coupled version
320 of UKESM1.

321 Figure 2 presents an overview diagram of the spin-up procedure. Several primary
322 branches of parallel spin-up are shown, each focused on the equilibration of separate ES
323 components: ocean, land and atmospheric chemistry. Model states (i.e. restart files of
324 prognostic variables) were shared between these main branches during the full spin-up,
325 with the main points of restart state sharing specifically identified. To illustrate underlying
326 operational details, Supplementary Tables S1 to S3 present the chains of simulations per-
327 formed as part of the “ocean”, “land” and “atmosphere” branches of spin-up (respectively
328 the top, middle and bottom paths of Figure 2). Significant changes along these branches
329 are switches from component-only to coupled branches, the switch from prescribed, non-
330 interactive atmospheric chemistry (designated UKESM1-CN) to fully interactive chemistry
331 (UKESM1), and the adoption of component model states as initial conditions from other
332 branches. Since it is the longest branch in terms of total simulated years, attention focuses
333 here on the ocean branch summarised in Supplementary Table S1.

334 As noted previously, the largest active carbon and heat reservoir in the Earth sys-
335 tem is the ocean, and imbalances in this reservoir can have a large impact on simulation
336 drift. Consequently, the ocean spin-up branch was prioritised and operationally began first,
337 principally in ocean-only mode before switching to a fully-coupled mode with prescribed,
338 non-interactive atmospheric chemistry (designated UKESM1-CN). This was followed by
339 land spin-up, which also started in land-only mode and also subsequently transitioned to
340 UKESM1-CN mode. Finally, the fully-coupled model, complete with interactive atmo-
341 spheric chemistry (designated UKESM1), was spun-up. The ocean, land and atmosphere
342 states from these parallel branches were then finally combined into a UKESM1 simula-
343 tion (identifier u-av472; Supplementary Table S1 that led into the pre-industrial control
344 simulation (identifier u-aw310; Supplementary Table S1.) This piControl was then simu-
345 lated for a duration of more than 1000 years both to act as a control for numerous other
346 simulations in CMIP6, and to serve as a source of initial states for the CMIP6 Historical
347 ensemble of UKESM1. This latter ensemble forms the subject of the analyses of Sellar
348 et al. [2019] and Archibald et al. [2020].

349 UKESM1 is the successor ESM to CMIP5's HadGEM2-ES, and, as noted already,
350 this latter model is the source of some components in UKESM1. The spin-up procedure
351 adopted for HadGEM2-ES also parallels that used here for UKESM1, with periods of
352 ocean- and land-only spin-up followed by a final phase of fully-coupled simulation, al-
353 though with some significant differences [Collins et al., 2011]. Unlike UKESM1's ocean,
354 where observationally-derived initial conditions were used, spin-up of HadGEM2-ES was
355 initialised using an existing ocean state from the preceding HadGEM1 model used in
356 CMIP3 [Johns et al., 2006]. This included both physical and biogeochemical state vari-
357 ables, with the new biogeochemical variables introduced in HadGEM2-ES initialised with
358 climatological (silicic acid) or uniform (total iron) values. Similarly to UKESM1, this ocean
359 state was then spun-up in ocean-only mode under model atmospheric forcing for 400 y,
360 but with the advantage of starting from the previous spun-up state of HadGEM1. In the
361 case of the land component, HadGEM2-ES used an acceleration technique in which, af-
362 ter 3 y periods of coupled simulation, the model's land state was implicitly extrapolated
363 forwards by 100 y before returning to a further period of conventional coupled simula-
364 tion. This procedure was repeated 4 times, advancing the land state of HadGEM2-ES by
365 400 y. This approach did not fully equilibrate refractive soil organic material because of
366 the timescales its equilibration (e-folding of 50 y), and its sensitivity to sub-annual lit-
367 ter input. Spin-up of the model's soil carbon was instead achieved using 2000 y of of-
368 fline simulation of this reservoir, forced using monthly fields of litter inputs. The ocean
369 and land states obtained using these procedures were then used in a final period of fully-
370 coupled simulation under pre-industrial conditions for 280 y, to produce a piControl state.

371 ***2.2.1 Detailed spin-up approach***

372 The key motivating factor in our spin-up was minimising drift in the Earth system's
373 carbon cycle, and attention was strongly focused on the net air-sea CO₂ flux. Analysis by
374 Séférian et al. [2016] found the diverse array of spin-up protocols followed during prepa-
375 ration for CMIP5 resulted in models that exhibited large differences in simulated fields,
376 and potentially biased performance evaluations. Recognising this, Jones et al. [2016] sug-
377 gest a drift criterion of $\leq 10 \text{ Pg C century}^{-1}$ (i.e. a long-term average of $\leq 0.1 \text{ Pg C y}^{-1}$)
378 for net fluxes between the atmosphere, land and ocean reservoirs as part of the C4MIP
379 protocol [Jones et al., 2016]. Note that both the land and ocean components of UKESM1

380 only exchange carbon with the atmosphere component and not directly with each other
381 (e.g. via rivers).

382 To evaluate this particular target, as well as track a range of critical physical and
383 biogeochemical properties (e.g. ocean heat content, surface temperature, top-of-atmosphere
384 heat balance, Atlantic meridional overturning circulation, sea-ice cover, etc.), the spin-up
385 was monitored throughout using the Met Office Climate Model Monitoring tool (CMM)
386 and BGC-val tools [de Mora et al., 2018]. Running routinely in parallel with the spin-up
387 simulations, these tools greatly facilitated rapid decision-making during model develop-
388 ment, as well as identifying undesirable drifts or model errors.

389 The spin-up path began with a short physical climate simulation (run ID u-ai567;
390 the full list of run IDs is given in Supplementary Tables S1 to S3) using a prototype of
391 HadGEM3 GC3.1 (Williams et al. [2017]; Kuhlbrodt et al. [2018]) using CMIP6 pre-
392 industrial control forcing. This constitutes the physical core of UKESM1. This simula-
393 tion was initialised from rest (i.e. with zero u and v velocity fields), with observationally-
394 derived initial conditions for the ocean (EN4; Good et al. [2013]), and initial states for
395 the atmosphere and sea-ice drawn from a GC3.0 simulation Kuhlbrodt et al. [2018]. Af-
396 ter 30 years, the atmospheric state of this simulation was judged to be sufficiently spun-up
397 to serve as a source of forcing data for ocean-only configurations, and the simulation was
398 continued to provide a further 30 year period of forcing data.

399 The forcing data collected from this GC3.1 simulation (and for subsequent forcing)
400 consisted of 1.5 m air temperature, air humidity, 10 m wind velocities (U and V direc-
401 tions), surface downwelling short- and long-wave radiation, precipitation (rain and snow)
402 and aeolian dust flux at 3 hour frequency, and river runoff at monthly frequency. These
403 data fields are the same as those available in observationally-derived reanalysis forcing
404 datasets, such as CORE [Large and Yeager, 2009] and DFS [DRAKKAR Group, 2007],
405 although at higher temporal resolution for heat and freshwater fluxes. In addition to these
406 properties, fields of ocean surface temperature and salinity were collected from the same
407 GC3.1 simulation at monthly frequency for relaxation purposes.

408 Based on the variability found in the atmospheric component of GC3.1, a forcing
409 period of 30 years was selected as broadly representative of interannual patterns (but see
410 later). Test simulations using repeated cycles of this forcing did not find any significant
411 issues associated with the forcing “kick” imparted between cycles (i.e. upon reaching the

412 end of year 30, year 1 is simply recycled). This approach of recycling forcing has previ-
413 ously been used successfully with NEMO-MEDUSA [Couldrey et al., 2016].

414 After this initial spin-up phase with HadGEM3-GC3.1, a successor phase was pre-
415 pared using an ocean-only configuration of a UKESM1 prototype (run ID u-aj588). The
416 ocean physical state of this (i.e. ocean, sea-ice, icebergs) was initialised using the model
417 state at end of year 30 of the preceding GC3.1 simulation. The ocean biogeochemical
418 state was initialised using observationally-derived fields from the World Ocean Atlas 2009
419 (WOA09; Garcia et al. [2010a]; Garcia et al. [2010b]) and Global Ocean Data Analysis
420 Project v1.1 (GLODAPv1.1; Key et al. [2004]) climatologies. Fields of dissolved inor-
421 ganic nitrogen (DIN), silicic acid and dissolved oxygen were drawn from WOA09, while
422 pre-industrial DIC and alkalinity were drawn from GLODAPv1.1. Following Yool et al.
423 [2013], the fields of DIC and alkalinity from GLODAPv1.1 were modified to interpolate
424 over large regional lacunae in the original climatology (the revised GLODAPv2 climatol-
425 ogy was not fully available at this time; Lauvset et al. [2016]). Note that although older
426 climatologies were used to initialise run u-aj588, subsequent evaluation primarily uses
427 their revised and updated equivalents, World Ocean Atlas 2013 and GLODAPv2.

428 Once initialised, this ocean-only configuration was run under repeated cycles of
429 the initial atmospheric forcing data for a total of 1890 years (i.e. 63 cycles of 30 years;
430 run IDs u-aj588 and u-ak900). During this initial, extended period of spin-up, a differ-
431 ence in the bulk formulae for atmosphere to ocean momentum flux between the coupled
432 UKESM1 and the ocean-only configuration was found. Changing this calculation so the
433 ocean-only model mimicked the coupled model calculation led to the updated ocean-only
434 run u-an869. However, because of the long duration of the initial ocean-only phase, and a
435 consistent “direction of travel” in the carbon cycle, this subsequent phase (u-an869) used
436 the end state of the initial phase (u-aj588) for its initial condition. This new ocean-only
437 phase also allowed an update to the atmospheric forcing that was taken from a longer du-
438 ration spin-up of the UKESM1 prototype, again using a 30-year period.

439 This subsequent phase was run for a further 2905 years (96.5 cycles; including run
440 ID u-ar538). During this ocean-only phase, trial simulations of the full coupled ESM, ini-
441 tialised using ocean states drawn from the ocean-only spin-up, found that the two modes
442 were comparable – though not identical – in terms of their evolving ocean properties and
443 in net air-sea CO₂ flux, with these test coupled runs typically reaching an equilibrated

444 state within 150 years. During this ocean-only phase net air-sea CO₂ declined to less
445 than 0.1 Pg C y⁻¹, as desired. Upon reaching this CO₂ target, spin-up was switched from
446 ocean-only mode to coupled Earth system mode.

447 Having determined the ocean-only configuration had reached a near-equilibrium
448 state, these ocean conditions were used to initialise the coupled ESM, which ran for a
449 further 500 years of spin-up before the start of the CMIP6 piControl. The first 300 years
450 of this coupled spin-up used the faster UKESM1-CN configuration of the model to max-
451 imise the equilibration of the ocean and terrestrial biosphere in the available time. This
452 UKESM1-CN configuration differs from the full UKESM1 model by using prescribed
453 chemical oxidants taken from a parallel UKESM1 pre-industrial run, but is otherwise
454 identical; see Appendix A of Sellar et al. [2019] for details.

455 The fully-coupled model required some science changes during this final coupled
456 spin-up to address important biases, many of which emerged as a result of coupling com-
457 ponents which had previously been spun up separately. These changes are extensively de-
458 scribed in Section 3 of Sellar et al. [2019]. The magnitude and impact of these changes
459 decreased as the spin-up progressed, and the last 200 years were performed with the final
460 UKESM1 science settings, and with the full-complexity model configuration (e.g. interac-
461 tive atmospheric chemistry now included).

462 In parallel to the ocean spin-up, there were separate spin-up phases for terrestrial
463 biogeochemistry and atmospheric chemistry, prior to their introduction into the main spin-
464 up simulation (Figure 2). The land state, and in particular forest cover and the soil carbon
465 and nitrogen pools, takes many hundreds of years to equilibrate with the surface climate
466 and carbon fluxes. Some aspects of land cover, such as grass cover, equilibrate relatively
467 quickly, so initial priority was therefore given to improving simulation of slower equili-
468 brating aspects, such as forest cover and soil carbon, with subsequent tunings applied to
469 the grass plant functional types and snow-vegetation interactions. However, as the whole
470 system is interactive, changes in grass colonisation affects soil carbon and nutrients which,
471 in turn, feeds back on vegetation productivity. The land was initially spun up in a 1000-
472 year offline simulation of the JULES land surface model, driven by surface forcing from
473 a GC3.0 simulation, a prototype of fully-coupled UKESM1. This land-only phase was, it-
474 self, initialised using the land state from a land-only simulation run in excess of 10,000
475 years using CRU-NCEP reanalysis forcing as derived for the Global Carbon Project [Le

476 Quéré et al., 2018]. Prior to this, the land model was run in excess of 10,000 years us-
477 ing CRU-NCEP observation-based meteorology from the Global Carbon Project. Similar
478 to the ocean-only spin-up, this approach considerably reduces simulation cost. The land
479 system underwent a further 665 years of coupled UKESM1-CN spin-up, including the im-
480 plementation of final tunings, before being used to re-initialise the land state for the final,
481 200 year UKESM1 spin-up phase. Thus, the terrestrial BGC fields experienced 865 years
482 of coupled spin-up in total, following the initial 1000 year offline spin-up.

483 Atmospheric timescales (ranging from minutes to tens of years) are much shorter
484 than those of the land and ocean. Nevertheless, a coupled spin-up of 230 years was per-
485 formed prior to the resulting atmospheric state being combined with the evolving ocean
486 and land states to initialise a final 200 y period of coupled UKESM1 spin-up. This ex-
487 tended duration was required because of solar radiation and surface temperature differ-
488 ences between UKESM1-CN and UKESM1 that impacted the land carbon and nitrogen
489 pools. It also served to avoid any impacts on the model's climate which might arise if
490 the chemical tracers were far from equilibrium with the other components at initialisation.
491 The atmospheric chemical tracers therefore experienced 410 years of pre-industrial cou-
492 pled simulation during the spin-up (as did the rest of the atmosphere component).

493 In summary (with reference to Figure 2), the separate ocean and land spin-up states
494 were combined into a single model initial condition after, respectively (4800 + 230) years
495 of ocean spin-up and (1000 + 710) years of land spin-up. After a further 80 years of
496 coupled integration, atmospheric chemistry fields were also combined with the evolving
497 land and marine 3D fields, providing a final initial state for a further 200 years of coupled
498 UKESM1 spin-up. We deemed the spin-up to be complete when this adjustment amounted
499 to a multi-decadal land carbon flux of less than 0.1 Pg C y^{-1} (averaged over a century), as
500 recommended in the C4MIP experimental protocol noted already Jones et al. [2016].

501 This simulation then initialised all components of the UKESM1's piControl simula-
502 tion, from which pre-industrial initial states were drawn for the CMIP6 historical ensemble
503 (see the CMIP6 implementation paper of Sellar et al. [2019]).

504 **2.3 Analysing the UKESM1 spin-up**

505 The complexity of UKESM1 means a complete evaluation needs to be spread over
506 several dedicated studies. Such studies to date include atmospheric chemistry in Archibald

507 et al. [2020] and aerosols in Mulcahy et al. [2018]. The physical climate model that un-
508 derpins UKESM1, HadGEM3, is assessed at the same atmosphere-ocean resolution in
509 Kuhlbrodt et al. [2018]. Meanwhile, an overview of the entire model, analysed for the
510 near-present using the CMIP6 historical ensemble, is provided by Sellar et al. [2019].

511 Evaluation here is focused primarily on the spin-up period itself to analyse the equi-
512 libration pathway of key climate-relevant properties. The model state of the piControl
513 simulation – the end point of spin-up – is then briefly analysed to evaluate the scientific
514 performance of UKESM1. This is done across the point from which the piControl is first
515 used to initialise CMIP6 Historical simulations. Evaluation of the piControl focuses on
516 the slow timescale variables that need to be spun-up. The piControl continues beyond this
517 point as a reference simulation for all other CMIP6 experiments.

518 The selection of the piControl somewhat complicates model evaluation since most
519 target fields only have near-present day observations available. Such data contains signals
520 of ongoing anthropogenic climate change that are absent in the pre-industrial period that
521 the piControl aims to represent. In the case of the deep ocean, one focus of the spin-up
522 described here, these signals are relatively minimal or absent, but they are manifest in the
523 surface ocean, the atmosphere and the land, although natural or background processes are
524 arguably still dominant for numerous variables. As such, intercomparison with observa-
525 tions is still informative, so long as differences are appropriately interpreted. Sellar et al.
526 [2019] provides comparisons of UKESM1 at Historical period time-points aligned with
527 modern observations.

528 The specific observational datasets used for evaluation include:

- 529 • World Ocean Atlas 2013, for ocean physics (interior; Locarnini et al. [2013]; Zweng
530 et al. [2013]) and biogeochemistry (interior, surface; Garcia et al. [2014]; Garcia
531 et al. [2014]) fields
- 532 • Hadley Centre Sea Ice and Sea Surface Temperature [Rayner et al., 2003], for ocean
533 SST and sea-ice fields
- 534 • Sea-Viewing Wide Field-of-View Sensor (SeaWiFS; O'Reilly et al. [1998], for sur-
535 face ocean chlorophyll concentration

- 536 • Oregon State University Ocean Productivity group, for VGPM [Behrenfeld and
537 Falkowski, 1997], Eppley-VGPM [Carr et al., 2006] and CbPM [Westberry et al.,
538 2008] vertically-integrated primary production
- 539 • Global Ocean Data Analysis Project v1.1 [Key et al., 2004] and v2 [Lauvset et al.,
540 2016], for interior and surface carbonate biogeochemistry
- 541 • Estimating the Circulation and Climate of the Ocean (ECCO) V4r4 (Forget et al.
542 [2015]; Fukumori et al. [2019]), for ocean circulation
- 543 • Smeed et al. [2017] for RAPID time-series measurements of the Atlantic merid-
544 ional overturning circulation at 26°N
- 545 • Poulter et al. [2015] for plant functional type classification
- 546 • Loveland et al. [2000] for global land cover characteristics

547 In addition to the above, model-observation intercomparison makes use of multi-
548 annual periods throughout, rather than focusing on a single year. This aims to account for
549 both interannual variability in the case of synoptic observations for which we have good
550 observational data (e.g. satellite-derived surface fields), and the representative timeframes
551 associated with composite observational datasets that are assembled over time (e.g. point
552 samples of the ocean interior). Observational products differ in their availability and ref-
553 erence periods, but in general are available from the late 20th century, and typically used
554 here for the period 1995 to 2010.

3 Results

3.1 Time evolution of the spin-up simulations

In the following, time-series analysis focuses initially on the ocean spin-up branch because of its extended duration. Time-series analysis of the land- and atmosphere branches (per Figure 2) focuses on periods closer to the start of the piControl when these shorter branches have merged with the ocean branch.

The panels of Figure 3 (and Supplementary Figure S1) track key physical properties over the full duration of the ocean branch of the spin-up. The panels break the spin-up into sections coloured according to different run modes: two ocean-only phases, a coupled UKESM1-CN phase, a fully-coupled UKESM1 phase, and a final section that corresponds to the formal CMIP6 DECK piControl. Supplementary Table S1 presents the full list of run IDs associated with the spin-up period depicted, with some continuous phases actually split between several run IDs. In each case the panels in Figure 3 present the 30-year rolling averages of the properties, together with the corresponding 30-year interannual range.

In terms of volume-averaged ocean bulk properties, while – unsurprisingly – neither exhibit large interannual variability when averaged globally, both temperature and salinity experience long-term drifts across ocean-only and fully coupled phases, and it is noticeable that the trends in the ocean-only phase are reversed (at least somewhat) during the coupled phase.

In the case of volume-averaged temperature (Figure 3, row 1), an upward drift of approximately $+0.06^{\circ}\text{C ky}^{-1}$ during the ocean-only phase flips to a downward drift of approximately $-0.25^{\circ}\text{C ky}^{-1}$ when UKESM1 transitions to fully coupled simulation. For salinity (Supplementary Figure S1), a slight upward drift of $+0.0015\text{ PSU ky}^{-1}$ is approximately reversed in the transition between ocean-only and fully coupled phases.

In the ocean-only phase, the small salinity trends are related to strong surface salinity relaxation and water flux balancing, while in the fully coupled phase they reflect the conservation of water across the modelled Earth system. Meanwhile, modest drift in the heat content of the ocean-only phase is explained by the use of repeating surface forcing derived from a period of GC3.1 coupled spin-up simulation that exhibited a $+0.2\text{ W m}^{-2}$ global mean, top-of-atmosphere (TOA) radiation imbalance (downward directed) under

586 which the ocean warms. Disequilibrium creates a situation in which the upper boundary
587 of the ocean is consistently driven in one direction. The switch to fully coupled UKESM1
588 (in which the global mean net TOA radiation balance is effectively 0 W m^{-2}) permits a
589 correction of this, as the slightly-too-warm ocean can then properly exchange heat with the
590 dynamic overlying atmosphere. In short, the excess heat gained by the ocean during the
591 ocean-only spin-up is lost again during the coupled phase of spin-up.

592 The panels of Figures 3 and ?? showing the corresponding surface quantities indi-
593 cate more distinct behaviour for temperature and salinity. The former finds comparable
594 ocean-only and coupled phases once the bulk formulae are harmonised across the two
595 model configurations. The latter shows surface salinity returning to its observationally-
596 derived initial value after a prolonged period of lower salinity during ocean-only spin-up.

597 Unlike the full-ocean averages of temperature and salinity, northern and southern
598 sea-ice areas (Figure 3, row 2) are highly dynamic, with large interannual variabilities
599 across both ocean-only and fully coupled phases. In the coupled model, this variability in
600 sea-ice area shows marked multidecadal patterns. In the ocean-only phase, sea-ice trends
601 in both hemispheres quickly equilibrate ($< 100 \text{ y}$) under the repeating atmospheric forcing,
602 albeit to slightly different averages between the forcings used. In the fully coupled phase,
603 interannual variability is comparable in magnitude in the north, but noticeably larger in
604 the south, and does not appreciably settle down in either hemisphere during the (short)
605 duration of the fully coupled phase. Sea-ice area and its seasonality is discussed further
606 later.

607 Finally, row 3 of Figure 3 shows two important metrics of ocean circulation, the
608 Atlantic Meridional Overturning Circulation (AMOC) at 26°N and the Antarctic Circum-
609 polar Current (ACC) transport through Drake Passage. The AMOC characterises the pole-
610 ward flow of warm water in the North Atlantic, playing both an important role in heat
611 transport and the conditioning of water masses leading to deep water formation in the sub-
612 polar Atlantic [Smethie et al., 2000]. Since 2004, the AMOC has been well-observed by
613 the RAPID array at 26.5°N , with an annual average ranging between 14.6–19.3 Sv [Smeed
614 et al., 2017]. Drake Passage is a “pinchpoint” for the circular ACC that rings Antarctica,
615 with a role in framing the continent’s isolated and cold climate, and in climate variabil-
616 ity modes such as the Southern Annular Mode (SAM) [Majewski et al., 2009]. The ACC
617 is balanced by the meridional density gradient throughout the depth of the ocean, which,

618 in turn, is set up by the wind and buoyancy forcing at the surface [Meredith et al., 2011].
619 While not permanently instrumented like the AMOC, Drake Passage is intermittently sam-
620 pled, with its transport estimated at 173 ± 11 Sv [Donohoe et al., 2016].

621 UKESM1's pre-industrial AMOC is typically lower than that found by RAPID, but
622 it necessarily omits the present-day greenhouse gas (GHG) and aerosol forcing. By con-
623 trast, all UKESM1 runs over the historical period simulate an AMOC that strengthens by
624 approximately 2 Sv to a maximum of around 17 Sv in the 1990s Sellar et al. [2019]. This
625 is almost certainly linked to northern hemisphere aerosols changing the simulated inter-
626 hemispheric energy gradient, cooling the north relative to the south, with AMOC strength
627 responding to this.

628 During the first portion of the ocean-only phase ("Ocean 1" in Figure 3), both the
629 AMOC and Drake Passage transport are at the bottom end of their observed ranges, and at
630 significantly lower values than those found in the corresponding coupled UKESM1 precur-
631 sor that provided the atmospheric forcing. As noted earlier, investigation of this uncovered
632 a discrepancy in the bulk formulae used in the transfer of momentum between the atmo-
633 sphere and ocean, with the ocean-only model following that of CORE [Large and Yeager,
634 2009] and the coupled model following COARE 3.5 [Edson et al., 2013]. This was rec-
635 tified in the subsequent, longer portion of ocean-only spin-up where the coupled model
636 formulation was used ("Ocean 2" in Figure 3). Nonetheless, we retained the first portion
637 of spin-up because we judged that it achieved an ocean carbon state that was closer to
638 equilibrium than the initial state in spite of this discrepancy.

639 A potential issue in using distinct ocean-only and fully coupled spin-up phases is a
640 mismatch in the behaviour of the model between these phases. In ocean-only mode, the
641 ocean model experiences the atmosphere as unchanging forcing, in fully coupled mode,
642 the ocean model dynamically interacts with the overlying atmosphere, potentially modi-
643 fying the evolving atmospheric forcing. Broadly mirroring Figure 3, Figure 4 compares
644 the behaviour of both phases for physical properties across a 200 y period from the time
645 point at which the coupled phase branches from the ocean-only phase with the coupled
646 ocean initialised using the ocean-only state. Although the ocean-only phase is generally
647 equilibrated at this point, to more clearly evaluate the significance of this transition, it was
648 continued past this branch point.

649 In the case of temperature (Figure 4, row 1) and salinity (Supplementary Figure
650 S2), the volume-integrated panels show the differences between the equilibrium values
651 to which the two phases have converged, or are converging. At the surface, SST in the
652 coupled phase remains close to that of its ocean-only parent (with larger interannual vari-
653 ability), while SSS quickly (< 100 years) equilibrates at a slightly higher value (though
654 with similar interannual variability; see Supplementary Figure S2).

655 In the case of sea-ice (Figure 4, row 2), the coupled phase very quickly shows larger
656 interannual to interdecadal variability, but the longer-term behaviour takes a more ex-
657 tended period to manifest (visible in Figure 3). The difference between the two spin-up
658 phases is more obvious in the case of southern hemisphere ice, where the cyclic 30-year
659 forcing period used in the ocean-only phase precludes the large multidecadal variability
660 exhibited by the coupled model.

661 The pattern of increased variability in southern sea-ice closely corresponds to that
662 of variability of coupled mode Drake Passage transport in row 3 of Figure 4. Here, high
663 transport is associated with reduced sea-ice area, and vice versa, with the two properties
664 connected via periodic deep ocean mixing off Antarctica influencing the latitudinal gradi-
665 ent of the ocean density field across the Southern Ocean and the Antarctic sea-ice extent
666 in a coherent manner (as discussed by Latif et al. [2013]; de Lavergne et al. [2014]). In
667 the ocean-only phase, ACC interannual variability is low (< 10 Sv), but grows quickly as
668 the coupled phase begins, reaching almost 40 Sv within the time period shown. As Fig-
669 ure 3 shows, this magnitude largely persists during the piControl simulation, with strong
670 centennial-scale variation in Drake Passage transport.

671 The change between spin-up phases is more slight for the AMOC, and after an ini-
672 tial shock (< 100 years), the AMOC between the two phases remains similar (Figure 3).
673 And, unlike the relationship between Drake Passage transport and southern sea-ice, the
674 AMOC's relationship with northern sea-ice is less clear. As noted, AMOC strength is
675 related to the inter-hemispheric energy gradient Marshall et al. [2014]. Poleward heat
676 transport driven by a strong AMOC might be expected to correlate with increased melt
677 of northern sea-ice. However, AMOC strength at 26°N does not show such a clear corre-
678 lation because the underlying relationship is more complex. For example, northern hemi-
679 sphere cooling, relative to that in the south, can act to both directly increase Arctic sea-
680 ice, and intensify the meridional energy gradient, leading to a strengthening of the AMOC,

681 which can negate the direct cooling influence on sea-ice decrease. Furthermore, AMOC
682 strength is also influenced by buoyancy and freshwater fluxes out of the Arctic, which im-
683 pact the occurrence of deep water formation Liu et al. [2019].

684 Across the ocean's physical parameters, the ocean-only and coupled phases do show
685 similar patterns and magnitudes of variability. There is limited evidence of strong shocks
686 as the model's state branches, and most properties quickly adjust, although the ocean's
687 different equilibrium heat content between the phases manifests in the change of inflection
688 of long-term drifts. Nonetheless, the prescribed 30 y atmospheric forcing in the ocean-
689 only phase clearly prevents the model from reproducing the longer-term modes visible in
690 the fully coupled phase, most noticeably in the circulation of the Southern Ocean.

691 Remaining with the ocean, Figure 5 shows corresponding spin-up time-series for
692 several key marine biogeochemistry metrics, over both the full period of the ocean branch
693 spin-up, and for the same 200 y overlapping period for the ocean-only to fully coupled
694 transition.

695 As already noted, the most significant features of this spin-up branch lie with how
696 the two periods of ocean-only spin-up differ in response to a change to the formulation
697 of surface momentum exchange. In the integrated primary production panel, addressing
698 this discrepancy results in a global increase of 15%. The mechanism for this large in-
699 crease lies in the increased momentum transfer, which can be seen in row 2 of Figure 5
700 to deepen average mixed layer depth (47 m to 50 m), leading to elevated surface DIN con-
701 centrations that fuel productivity. This change between the ocean-only phases is markedly
702 larger than the 4% decrease in primary production driven mainly by an 8% decrease in ae-
703olian deposition of iron as the 30 year cycle of dust forcing becomes dynamic in the fully
704 coupled simulation. It is also noticeable that the model's productivity response to such
705 transitions requires a longer period to equilibrate than seen for the earlier physical prop-
706 erties. Here, periods of at least several hundred years, and approaching 1000 y, are nec-
707 essary for the model to reach a new quasi-steady state. Nonetheless, despite only a 30 y
708 cycle in forcing during the ocean-only stages, the interannual variability of productivity is
709 similar, though slightly greater, to that in fully coupled UKESM1. This can also be seen
710 in intercomparison of the 100 y sections of ocean-only and fully coupled simulation.

711 Bar a short initial period of ingassing, net air-sea exchange of CO₂ is, on average,
712 outgassing across the entire spin-up and into the piControl. Interannually, both ingassing

713 and outgassing occur, but the long-term trend is to steadily outgas as equilibrium is ap-
714 proached. As already mentioned, an initial target for average net air-sea flux was 0.1 Pg C y^{-1} ,
715 and this was reached after around 3500 y, during the ocean-only stage of spin-up. By the
716 start of Historical ensemble simulations, an annual average outgassing flux of around -
717 0.04 Pg C y^{-1} had been reached, with a multi-centennial range of approximately -0.35 to
718 0.25 Pg C y^{-1} . Due to the repeating 30 y cycle of surface forcing (e.g. wind-driven piston
719 velocity), progress towards this equilibrium is steadier during the ocean-only phase (with
720 the exception of the jump between ocean-only stages), although the range of interannual
721 variability is very similar between ocean-only and fully coupled phases. The continuous
722 outgassing of CO_2 from the ocean is indicative of a model bias in ocean carbon content
723 and is discussed in more detail later.

724 The bottom two panels of Figure 5 show how productivity and air-sea exchange vary
725 interannually across the transition between the ocean-only and fully-coupled phases. While
726 there is a slight offset in primary production between these phases, the modelled variabil-
727 ity is otherwise largely consistent. The same is true for air-sea CO_2 flux, for which both
728 phases oscillate interannually around near-zero net flux. Overall, and much as with the
729 model's physics, the differences between the two spin-up modes are relatively minor.

730 Together with the physical responses shown in Figure 4, these results indicate that
731 the coupled model largely adjusts to a new equilibrium after around 150 y, even when its
732 ocean is initialised from the end of an ocean-only simulation.

733 Finally, switching to the terrestrial system, Figure 6 illustrates the spin-up path-
734 way of the carbon reservoirs in living biomass and soil, and the global fractional cover
735 of three major aggregate land surface types: forests, grasslands and bare soil. Unfortu-
736 nately, due to archiving issues not all of the model data are available. Any data gaps,
737 however, occur earlier than 500 years before the piControl, so do not affect our evalua-
738 tion of the model's final equilibrium state. The land spin-up branch began with the land
739 surface model being run offline under 30-year cycles of meteorological data drawn from
740 a prototype of UKESM1. The land-model was run for approximately 1000 years offline
741 until the globally-averaged vegetation cover and carbon and nitrogen pools (results not
742 shown) reached a quasi-equilibrium state. These were then used to initialise the coupled
743 model at timepoint #1 of Figure 6. As is evident from the immediate drift the land-model
744 run offline has a different stable state to the coupled model, reflecting the importance of

745 land-atmosphere coupling and differences in meteorology to the forcing used offline. Sub-
746 sequently, after further 30 years at timepoint #1 (Figure 6), the turnover rate parameters
747 were scaled to lower values to increase the size of the soil carbon and nitrogen pools.
748 At the same time, the soil carbon and nitrogen pools were rescaled to be consistent with
749 these turnover rate changes. The spin-up of the vegetation fractions continued with an ad-
750 ditional tuning applied to the rate at which grass can expand and colonise bare ground,
751 which was reversed at timepoint #2 (Figure 6). The result was a rapid decrease in grass
752 cover and concurrent increase in bare soil over a 10-year period. Over the course of the
753 next 600 years the spin-up continued with some changes made to parameters control-
754 ling snow-vegetation interaction and the rate of grass colonisation resulting in a close to
755 stable global mean state at the start of the final spin-up at timepoint #3 (Figure 6). The
756 drifts in soil and vegetation carbon over the course of the piControl were -0.07 and 0.0025
757 Pg C y^{-1} respectively over the first 1000 years of the piControl, well within the C4MIP
758 acceptable range. The drift in tree cover is also small at less than 0.5% over 1000 years.

759 Although the global drift may be small, it can be more significant at the regional
760 level, particularly if some regions are compensating for changes in carbon or vegetation
761 cover in other regions. Figure 7 shows the drift in soil carbon across seven major biomes
762 for the final part of spin-up and the piControl. The drift in most biomes is less than 0.001
763 Pg C y^{-1} with the exception of the tundra, boreal and desert regions. Tundra and boreal
764 regions lose soil Carbon at -0.012 and $-0.017 \text{Pg C y}^{-1}$ of carbon per year respectively.
765 This reflects the particularly long residence of soil carbon in these regions and therefore
766 the greater time required for the pools to equilibrate. The desert regions continue to accu-
767 mulate carbon at 0.002Pg C y^{-1} responding to changes made during the spin-up phase to
768 grass colonisation rates and therefore the flux of litter to the soil carbon. The pools also
769 demonstrate some long-term variability. For instance across the 250 years corresponding
770 to the first Historical ensemble member, the Savanna biome accumulates 3Pg C despite
771 having lost carbon during the preceding spin-up period. Consequently, we recommend that
772 any analysis accounts for both the ongoing drift in terrestrial carbon pools and the multi-
773 annual variability. Further, our regional drifts imply that benchmarking global drift pools
774 is a necessary but possibly insufficient condition for evaluating the equilibration of land
775 carbon models. Future spin-up efforts may wish to execute longer spin-ups in order to
776 equilibrate all regions and ecosystems.

777 The atmosphere adjusts rapidly (over days or weeks) to a changed external forcing,
778 such as associated with different sea surface temperatures or incoming top of the atmo-
779 sphere (TOA) solar radiation. We therefore, are generally not overly concerned with the
780 spin up of atmospheric variables when bringing ESMs into balance with a pre-industrial
781 forcing. Nevertheless, one of the primary constraints used to evaluate ESM simulations of
782 the pre-industrial period is that, averaged over sufficiently long timescales, the global mean
783 net TOA radiation balance should be zero (0 W m^{-2}). This was a leading constraint used
784 in developing UKESM1. A consequence of a zero TOA net radiation balance is that, also
785 averaged over sufficiently long timescales, the global mean energy content of the climate
786 system should be temporally stable. As ocean heat content constitutes the overwhelming
787 majority of energy in the climate system Trenberth et al. [2014], this constraint equates
788 to a temporally stable global mean ocean heat content. Observational constraints of the
789 absolute value of the pre-industrial ocean heat content are not available, neither are con-
790 straints on the component, solar (shortwave) and Earth-emitted (longwave), TOA radi-
791 ation fluxes. Observational estimates of global mean TOA radiation components and ocean
792 content do exist for present-day conditions, albeit with non-negligible uncertainties (Loeb
793 et al. [2009]; Stephens et al. [2012]; Cheng et al. [2017]). However, both estimates in-
794 clude an anthropogenic component. Observational estimates of pre-industrial (more cor-
795 rectly the very early industrial period, e.g. 1850 to 1900) global mean surface air temper-
796 ature (GSAT) do exist (e.g. HadCRUT4, Morice et al. [2012]; GISSTMP, Lenssen et al.
797 [2019]), although observation coverage is limited during this early period. With this in
798 mind, our primary targets for the UKESM1 pre-industrial atmosphere are: a near-zero
799 global mean net TOA radiation balance and a temporally stable GSAT, close to obser-
800 vational estimates for the 1850-1900 period. As a consequence of these two constraints,
801 temporally stable Arctic and Antarctic mean sea ice amount and volume is also a useful
802 constraint.

803 Figure 8 summarizes these quantities over the final 500 years of the coupled spin
804 up. The first 300 years of this run uses UKESM1 with offline atmospheric chemistry (re-
805 ferred to as UKESM1-CN), with ozone and chemical oxidants prescribed from an earlier
806 pre-industrial simulation of UKESM1 with interactive chemistry. The latter 200 years are
807 with interactive atmospheric chemistry enabled. The simulation is initialized at year minus
808 500, in the ocean by fields derived from the ocean-only spin up (u-ar538) and on land and
809 in the atmosphere using fields from the parallel land-spin up run of UKESM1-CN (shown

810 in Figure 8). At year minus 200, interactive atmospheric chemistry is activated and the re-
811 quired chemistry fields initialised based on output from the parallel “UKESM1 spin-up for
812 atmospheric chemistry” run shown in Figure 8. All other prognostic fields are propagated
813 from the UKESM1-CN spin up run (u-ar783 → u-au835, years -500 to -200) into the final
814 UKESM1 spin-up (years -200 to 0).

815 The primary spin-up characteristic in Figure 8 is an increase (of 1 Wm^{-2}) in both
816 global mean TOA net downward solar (SW) and outgoing longwave (LW) radiation, lead-
817 ing to a near-zero, global mean net TOA radiation budget at year 0 (the start of the pi-
818 Control simulation). This shift clearly occurs at the point when UKESM1-CN switches to
819 include interactive chemistry (UKESM1) and results from two differences between these
820 model configurations: (i) in the manner marine-emitted DMS is processed through to
821 cloud condensation nuclei (CCN) in the model atmosphere, and (ii) in the simulation of
822 stratospheric ozone. Both differences influence the absorption and reflection of solar radi-
823 ation and therefore net TOA solar radiation. As a result of these differences, and to retain
824 a near-zero global mean net TOA radiation balance as we transitioned from offline to in-
825 teractive chemistry, it was necessary to introduce a small tuning to the parameterization of
826 seawater DMS in UKESM1 (see Sellar et al. [2019] for more details). This tuning acted
827 to reduce seawater DMS in biologically inactive regions of the global oceans, reducing
828 the average cloud droplet number in marine clouds and thereby reducing the simulated at-
829 mospheric solar reflectivity and retaining the desired near-zero net TOA radiation balance
830 (-0.09 W m^{-2} downward; Sellar et al. [2019]).

831 From the start of the UKESM1 piControl (year 0 on Figure 8), slower timescale
832 variability in GSAT appears to largely disappear, in concert with a reduction in the vari-
833 ability of Antarctic sea ice. In the early part of Figure 8 these variables exhibit an inverse
834 correlation, driven by variability in ocean overturning in the far Southern Ocean (as dis-
835 cussed earlier). While it is tempting to conclude the final coupled tuning reduced this in-
836 ternal variability, we note that similar timescale variability does intermittently reappear in
837 later periods of the full UKESM1 piControl. Finally, long-term mean GSAT during the
838 first 200 years of the piControl is around 287.5K (13.35°C), suggesting a cold bias of
839 $0.3\text{-}0.4^\circ\text{C}$ in UKESM1 compared to available observational estimates for the period 1850-
840 1900 (Morice et al. [2012]; Lenssen et al. [2019]).

841 **3.2 Equilibrium state**

842 We now analyse the equilibrium state that results from the confluence of all three
843 spin-up branches, focusing on the UKESM1 piControl simulation from the point at which
844 CMIP6 Historical ensemble members begin to be drawn. This time-point occurs early in
845 the piControl simulation but at a point where UKESM1 was judged to be sufficiently equi-
846 librated. The strategy for using the piControl as the source for the Historical ensemble,
847 and additionally its role in controlling for model drift, is described in detail by Sellar et al.
848 [2019]. We use a decadal climatology of the piControl from this point throughout.

849 **3.2.1 Ocean**

850 Figure 9 compares simulated sea surface temperature (SST) with the HadISST observation-
851 derived product, HadISST, for the period 1870-1879 [Rayner et al., 2003]. This period is
852 chosen as it is closest to that which the piControl simulation aims to represent (1850), but
853 note that HadISST is a data-assimilated reanalysis product with relatively limited observa-
854 tional constraint for this time period (but see Supplementary Figure S3).

855 Northern and southern summer periods are shown, together with the differences be-
856 tween the model and observations. In general terms, the model shows very similar pat-
857 terns to those observed, both geographically and seasonally. Nonetheless, the difference
858 plots show persistent biases in the model, including a warm Southern Ocean, warm up-
859 welling regions and strong cold bias in the western North Atlantic. The latter feature is
860 the most pronounced of a series of zonal, dipole-like biases in the North Atlantic, which
861 include warm biases off the eastern seaboard of North America and in the Irminger / Ice-
862 land basins, and a cold bias in the Greenland-Iceland-Norwegian (GIN) Sea.

863 In the case of the Southern Ocean, this warm SST bias is primarily driven by a cor-
864 responding positive bias in downward shortwave radiation that originates in cloud biases
865 (i.e. in cloud amount and albedo). The warm SST biases in upwelling regions are prin-
866 cipally a result of the relatively coarse resolutions of UKESM1's ocean and atmosphere
867 components. In the ocean, the model cannot represent the necessary small-scale features
868 of coastal upwelling, while in the atmosphere, coastal wind forcing cannot be resolved.
869 The root of the North Atlantic dipole bias has a similar cause, with resolution causing the
870 separation of the Gulf Stream from the eastern seaboard of North America to occur too
871 far south, resulting in a path that is too zonal Kuhlbrodt et al. [2018].

872 Figure 10 illustrates the seasonal extents of northern and southern sea-ice, again
873 compared to the same period of HadISST (but see Supplementary Figure S4). In the Arc-
874 tic, modelled sea-ice extent is always greater than that observed, with the excess ranging
875 between 1–4 10^6 km² seasonally, but greatest around the annual sea-ice minimum. By
876 contrast, in the Antarctic, it is the simulated minimum sea-ice extent that most closely
877 matches that observed, but modelled sea-ice growth is conspicuously weaker, leading to
878 a maximum extent only around two-thirds of that observed. These patterns of sea-ice
879 biases generally align with the SST biases in the GIN Sea (cooler) and Southern Ocean
880 (warmer).

881 Switching to focus on the ocean interior, Figure 11 illustrates the biases of temper-
882 ature and salinity within the ocean, again compared to the WOA. In each case, the plots
883 present so-called “thermohaline circulation” sections that centre the zonal averages of both
884 major basins around the interconnecting Southern Ocean (see Figure 11 for more details).

885 In the case of potential temperature, UKESM1 shows a general warm bias in the
886 upper 3 km, a smaller cold bias below this. This pattern differs between basins, with the
887 Atlantic showing a much stronger bias, particularly in the upper 1 km of the subtropics
888 (30°S–30°N), where it can exceed 4°C. The corresponding region of the Pacific generally
889 shows a cool bias, although with a more complicated structure. Despite the marked warm
890 bias in its surface waters noted earlier, the Southern Ocean shows generally weaker biases,
891 particularly in its main Pacific sector. Similarly, the salinity biases in UKESM1 broadly
892 track those of temperature, with a similar strong positive bias in the subtropical Atlantic,
893 and a negative bias in the subtropical Pacific.

894 Moving to ocean circulation, Figure 12 shows the global streamfunctions of merid-
895 ional overturning circulation (MOC) for the model and the observationally-derived Esti-
896 mating the Circulation and Climate of the Ocean (ECCO; Forget et al. [2015]; Fukumori
897 et al. [2019]) product. Qualitatively, the simulated MOC broadly follows that observed,
898 with a strong positive cell focused in the upper water column, driven by the AMOC, and
899 a weaker negative cell at depth, driven by Antarctic Bottom Water formation. Compared
900 to that in ECCO, the model exhibits a slightly weaker Deacon Cell centred around 50°S
901 [Döös and Webb, 1994], indicative of weaker surface wind stress over the model’s South-
902 ern Ocean. The maximum strength of the model’s deep AABW cell is also weaker than
903 estimated in ECCO. Supplementary Figure S5 shows the corresponding simulated MOC

904 patterns for the Atlantic and Indo-Pacific sections. In the Atlantic strong northward flow in
905 the surface of the Atlantic is balanced by the production at high latitudes of NADW that
906 flows southward at depth, and overlies an Antarctic Bottom Water (AABW) cell driven
907 from the Southern Ocean. However, while the strength of this northward flow is similar to
908 that observed as noted earlier (cf. Figure 3), circulation of the large AABW cell underly-
909 ing the NADW in the Atlantic is very weak (especially when compared to the correspond-
910 ing cell in the Pacific). As noted in Figure 11, this cell is characterised by cool and fresh
911 biases that are indicative of a Southern Ocean origin. While this pool is fed in part by
912 Antarctic Bottom Water (AABW), it shows biogeochemical properties that are suggestive
913 of a more sluggish transport than typical for this watermass (see later).

914 Switching to marine biogeochemistry, Figure 13 shows model-observation intercom-
915 parison of three fields of key properties: surface nitrogen nutrient, surface chlorophyll and
916 vertically-integrated primary production. In each case, these fields are annual averages,
917 with the nutrient field drawn from the present-day climatology of the World Ocean Atlas
918 2013 [Garcia et al., 2014], and the lower two fields from the period 2000-2009.

919 Row 1 shows DIN, the main limiting nutrient for biological productivity across the
920 World Ocean. The general pattern of higher values in upwelling regions and at higher
921 latitudes, particularly the Southern Ocean, and low values in ocean gyre regions is sim-
922 ulated. However, the model does display a number of marked biases: concentrations are
923 markedly elevated in and around equatorial Pacific upwelling, particularly in the adjacent
924 South Pacific. These discrepancies stem from upwelling of excessively DIN-rich deep
925 water, as is clearer from ocean interior concentrations (see later). These patterns of mis-
926 match in MEDUSA-2.1 are very similar to those found previously by Yool et al. [2013]
927 with MEDUSA-2 (despite a considerably longer period of spin-up).

928 Row 2 shows the observed and modelled surface chlorophyll concentrations [O'Reilly
929 et al., 1998]. While, again, MEDUSA-2.1 captures the broad patterns of the observed
930 field, agreement is much weaker, and the model displays a number of strong biases. Most
931 clearly, simulated Southern Ocean concentrations are noticeably higher, with elevated con-
932 centrations also seen in Equatorial and subtropical Pacific concentrations, in keeping with
933 the corresponding DIN excess availability. Away from the Pacific, oligotrophic gyre con-
934 centrations are much more biased downwards, with relatively large regions of very low

935 chlorophyll in the subtropical Atlantic. Spatial patterns in this basin are also somewhat
936 aberrant with a pronounced patchiness that is absent in the observations.

937 As noted by Yool et al. [2013] (and Kwiatkowski et al. [2014]), chlorophyll is gen-
938 erally poorly represented in marine biogeochemistry models, with models frequently per-
939 forming much better for fields of other bulk properties (nutrients, carbon) and productivity.
940 As well as its observed high dynamic range (note the plot log scale), Yool et al. [2013]
941 suggest that this may stem from the strong plasticity (in reality and in models) of chloro-
942 phyll:carbon ratios relative to other quantities, and the resulting high dynamic range.

943 Finally, row 3 shows vertically-integrated net primary production, with observations
944 represented by the simple average of three observationally-estimated products, VGPM
945 [Behrenfeld and Falkowski, 1997], Eppley-VGPM [Carr et al., 2006] and CbPM [West-
946 berry et al., 2008]. Although this is empirically derived from satellite chlorophyll (as well
947 as other fields), MEDUSA-2.1's agreement with it is greater than for surface chlorophyll.
948 The observed patterns of low and high values are generally reproduced, again with biases.
949 These include excessive productivity in the Southern Ocean throughout the year, more
950 latitudinally-focused productivity in the equatorial Pacific, and noticeably low productivity
951 in the North Atlantic.

952 As discussed, one of the drivers for UKESM1's spin-up is equilibration of the ma-
953 rine carbon reservoir. Figure 14 compares the observed and simulated surface concentra-
954 tions of DIC and total alkalinity [Lauvset et al., 2016]. DIC here is an observation-based
955 estimate of pre-industrial DIC as this biogeochemical property has changed significantly
956 since the beginning of the industrial revolution (especially so in the surface; Lauvset et al.
957 [2016]). Modelled DIC concentrations generally show good agreement with the observed
958 estimates, although model DIC is somewhat elevated in the Southern Ocean, while biased
959 low in the Indonesian Archipelago and parts of the Arctic (although data availability re-
960 mains somewhat limited in this region). While the patterns of model surface alkalinity are
961 similar to those observed, alkalinity is generally lower in the model, and there are notice-
962 able biases, particularly in the North Pacific and, again, the Indonesian Archipelago.

963 Since alkalinity acts in part as a buffering capacity for DIC, generally lower sea sur-
964 face alkalinity will reduce the amount of DIC in surface waters, and, in turn, the ocean
965 interior. The interior impacts can be seen in Figure 15 and Supplementary Figure S6,
966 which respectively show DIC and alkalinity along thermohaline transect sections (see ear-

967 lier). These show both lower alkalinity in the upper 1 km of the water column, and the
968 correspondingly lower DIC throughout the ocean interior. This is most obvious in the
969 southward-moving NADW and in the deep waters (> 1 km) of the North Pacific, where
970 model bias can exceed 100 mmol C m⁻³.

971 3.2.2 Land

972 Figure 16 presents a biome-based evaluation of land cover at the end of UKESM1
973 spin-up against two observationally-derived estimates, the IGBP and CCI products. The
974 upper panel provides a geographical perspective of where different biomes are located,
975 while the lower panel shows the fraction of each land cover type that occurs in each biome
976 for the model and from the data products, keeping in mind that vegetation type, amount
977 and geographical distribution are dynamically predicted in UKESM1. Overall, the model
978 performs well, with simulated biomes largely capturing their observed compositions, al-
979 though there are some biases. With the exception of the high latitude biomes, such as bo-
980 real forest and (especially) tundra, UKESM1 underestimates the observed fractional cover
981 of C3 grasses. In tropical forest, it is largely replaced by broadleaf trees, while in extra-
982 tropical forests and deserts its low bias is countered by a high bias towards C4 grasses.
983 In the grassland and savanna biomes, where grasses are found to dominate, C3 grasses
984 are displaced by forest, mostly by broadleaf trees in savanna and needleleaf trees in grass-
985 lands. C4 grasses, meanwhile, are typically biased positive, while modelled shrubs show
986 mixed biases with observations across the biomes. In terms of bare ground (i.e. no veg-
987 etation cover), UKESM1 only shows a bias in the tundra biome where C3 grasses are
988 overly abundant. As the IGBP and CCI products are assembled from present-day obser-
989 vational datasets, they include vegetation cover changes driven by human influences. At
990 least in part, the land “biases” identified in UKESM1 are consistent with these changes in
991 land cover, and therefore indicate the problem of evaluating a pre-industrial climate state
992 using present-day observations.

993 Figure 17 complements Figure 16 by illustrating the geographical patterns of frac-
994 tional cover of each land cover type. As noted, UKESM1 simulates excessive broadleaf
995 forest cover and extent in the tropics, particularly in South America, but also in Africa and
996 southeast Asia. However, UKESM1 does not include fire feedbacks, which may explain
997 some of this overestimate. Inclusion of an interactive treatment of fire in other models
998 [Hantson et al., 2016], as well as in our land surface scheme when driven by observed

999 climate improves this type of vegetation bias [Burton et al. , 2019]. The geographical
1000 range of needleleaf forest is generally modelled well, though fractional cover is elevated
1001 in northern boreal forests, and there is an anomalous Asian forest in the vicinity of Ti-
1002 bet. As well as being biased low, the extent of UKESM1's C3 grasses is noticeably cir-
1003 cumscribed, with almost no grasslands in southeast Asia, and reduced extents in Europe
1004 and the Americas. However, as already mentioned, UKESM1 does erroneously simulate
1005 solid C3 grass cover across northern Siberia. The geographical cover of C4 grasses is bet-
1006 ter than for C3 grasses, but there is a marked positive bias in Australia, as well as west-
1007 erly displacements of their abundance in both the northern and southern Americas. Shrub
1008 cover is generally underestimated across the world, with exceptions only in Asia, again
1009 around northern Siberia and Tibet. Finally, UKESM1's patterns of bare soil generally map
1010 well to those observed. The exceptions lie in high northern latitudes, where tundra areas
1011 have excessive C3 and shrub cover in UKESM1.

1012 **4 Discussion**

1013 The new UKESM1 model has been jointly developed by NERC and the UK Met
1014 Office to represent the fully-coupled Earth system with a state-of-the-art level of real-
1015 ism. UKESM1 succeeds its CMIP5 predecessor, HadGEM2-ES, and while incorporating
1016 evolved forms of components from this earlier model, is almost a wholly new model. A
1017 particular effort was made to ensure that coupling between ES components and physical
1018 climate components was fully prognostic, enhancing the utility of UKESM1 for investi-
1019 gating future coupled ES feedbacks. As part of its preparations towards use in CMIP6,
1020 UKESM1 requires the production of a pre-industrial control state that can be used to ini-
1021 tialise the DECK and other MIP experiments. Critically, this piControl should exhibit a
1022 near-steady state climate so that forced trends introduced to the model Earth system in var-
1023 ious experiments are clearly distinct, and not confounded by model drift.

1024 The primary components of the Earth system are its major reservoirs of heat and
1025 carbon – the atmosphere, the ocean and the land. The physical sizes of these components,
1026 and the timescales of the major processes that govern them, both physical and biogeo-
1027 chemical, mean that equilibration to achieve a steady state is necessarily prolonged relative
1028 to the perturbation experiments typically performed in CMIP6. Furthermore, the full ESM
1029 is computationally expensive to run, with a turnaround time of only a few simulated years
1030 per wallclock day. However, the most expensive component of UKESM1, the atmosphere
1031 and its attendant chemistry, is also the fastest to equilibrate. Consequently, the strategy
1032 adopted here was to spin-up the slow equilibrating components, the ocean and the land,
1033 decoupled from the atmosphere, and to only bring the full model together once much of
1034 their time evolution was complete.

1035 In the case of the ocean, a period of 4800 years of ocean-only simulation was re-
1036 quired to achieve a net air-sea CO₂ flux within the threshold suggested by the C4MIP
1037 community. In the case of the land, a corresponding period of 1000 years was used to
1038 bring the modelled reservoirs of carbon into the same net balance with the atmosphere.
1039 During both of these phases, the individual component models were run in forced mode,
1040 under an atmospheric dataset (bulk properties, heat and freshwater fluxes, winds) derived
1041 from simulations of precursor versions of UKESM1 run with preindustrial forcing. Sub-
1042 sequently, the model states from both component-only spin-up branches were combined
1043 with the atmosphere, and UKESM1's spin-up was finalised in fully-coupled mode, first

1044 with prescribed atmospheric chemistry (UKESM1-CN), and subsequently with interactive
1045 chemistry also activated.

1046 The branch of ocean spin-up found relatively rapid stabilisation ($\ll 1000$ y) of
1047 near-surface physical variables and major circulation metrics. Interior temperature and
1048 salinity, however, exhibit prolonged drift, which changes sign from ocean-only to coupled
1049 phases, although in the case of temperature, of very low magnitude compared to simulated
1050 trends over the Historical period [Sellar et al., 2019]. Biogeochemical processes, such as
1051 productivity and surface nutrients, typically had somewhat slower stabilisation (≈ 1000 y).
1052 Net ocean surface carbon flux, like interior temperature, essentially exhibited steady de-
1053 crease over the spin-up, slowly reaching a net flux below the target threshold. Examination
1054 of interior carbon concentrations shows that this slow decline is a function of the marine
1055 biogeochemistry model “favouring” a slightly lower total carbon inventory, driven, at least
1056 in part, by a bias towards lower sea surface alkalinity (cf. Halloran et al. [2015]). A no-
1057 table, if unwelcome, feature of the ocean-only phase of spin-up was a bulk formulae dis-
1058 crepancy that initially resulted in lower momentum transfer between the atmosphere and
1059 ocean in the ocean-only configuration compared to the coupled model. While this clearly
1060 affected the absolute magnitudes of properties across the model (Figures 3 and 5), this
1061 was amended without any significant lasting impact on the broad state of the ocean, with
1062 the subsequent revised period of ocean-only spin-up ultimately coming to more closely re-
1063 semble that of the final, fully-coupled model. The transition between the ocean-only and
1064 fully-coupled phases was found to introduce a “kick” across the model, although the im-
1065 mediate effects of this were typically found to quickly (≈ 100 simulated years) settle, fol-
1066 lowed by slower evolution to a slightly different final coupled state, for some predicted
1067 variables.

1068 Overall, the ocean-only spin-up compares well with that in fully-coupled mode.
1069 Inevitably, the modes and scale of variability exhibited in the ocean-only configuration
1070 are reduced compared to that of the fully coupled model (e.g. Drake Passage transport),
1071 partly because of the limited variability in the short period of atmospheric forcing used,
1072 but mostly because the absence of coupled responses between the ocean and atmosphere.
1073 In terms of model biases, several were noted in the ocean’s physical and biogeochemical
1074 spun-up pre-industrial state, most significantly the carbon deficit already noted.

1075 The inclusion of an interactive nitrogen cycle in UKESM1 has made for very a slow
1076 spin-up because of the interaction between soil and vegetation. The mineralisation of soil
1077 inorganic Nitrogen fertilises vegetation and encourages growth and the turnover and qual-
1078 ity of vegetation litter affects the soil state. The spin-up of the model has through neces-
1079 sity gone hand in hand with the application of tunings and fixes as the model advances to
1080 being frozen and ready without the possibility of a long (in excess of 500 years) spin-up
1081 post-freeze. The computationally-efficient offline JULES model, forced using surface level
1082 atmospheric fields, was used initially for approximately 1000 years of spin-up. However,
1083 when subsequently coupled directly to the atmosphere the model's behaviour was found
1084 to differ, primarily because in the coupled model the change in vegetation state is able to
1085 feedback on the climate. The result is that an extended period of online spin-up is still re-
1086 quired. Furthermore, the model shows some long periods of variability making it hard to
1087 assess the degree of drift whilst model integrations are proceeding.

1088 One of the further challenges is deciding on an appropriate pre-industrial state given
1089 the general lack of observational data for the 1850s. We generally rely on present day data
1090 such as the landcover (Figure 16) and make informed assessments around the expected
1091 level of change over the past due to land-use change and the role of climate.

1092 In UKESM1, we have achieved a near spun-up state for the ocean and land carbon
1093 pools well within the requirements of C4MIP for making assessments of carbon budgets
1094 for climate targets. However, as is shown in Figure 7 there can be significant regional
1095 drifts, which in some cases may oppose each other and give the impression of a better
1096 steady state. In UKESM1, the land is generally losing carbon driven by soil carbon losses
1097 in the Boreal and Tundra biomes. These are the regions with the slowest carbon residence
1098 times and therefore the most difficult to equilibrate. The high latitude losses are slightly
1099 offset by the small positive drift we see in the Savana biome. While both the land and
1100 ocean components are losing carbon to the atmosphere, its fixed pre-industrial CO₂ con-
1101 centration masks this. However, in the fully-coupled emission-driven model, these net
1102 fluxes to the atmosphere would result in a positive drift in atmospheric CO₂. The spin-
1103 up of the emission-driven model is a separate activity that takes advantage of the more
1104 completely spun-up state from a later time point of the piControl.

1105 In terms of the major Earth system quantities pertinent to anthropogenic change, the
1106 duration of spin-up in UKESM1 allowed these to reach quasi-equilibrium. After tuning

1107 (see Sellar et al. [2019]), net top-of-atmosphere radiation balance reached -0.09 W m^{-2} by
1108 the conclusion of UKESM1 spin-up, compared to the perturbed present-day net flux of ap-
1109 proximately $0.61\text{--}0.81 \text{ W m}^{-2}$ [Johnson et al., 2016]. Surface ocean temperature drift for
1110 the same spin-up period was $0.016 \text{ }^\circ\text{C decade}^{-1}$, as compared with observation-based esti-
1111 mated ranges during the historical period of 0.042 to $0.054 \text{ }^\circ\text{C decade}^{-1}$ (1880–2012) and
1112 0.072 to $0.124 \text{ }^\circ\text{C decade}^{-1}$ (1979–2012) [Hartmann et al., 2013]. Exchange of CO_2 be-
1113 tween the atmosphere, land and ocean, reached net fluxes of 0.020 and $-0.039 \text{ Pg C y}^{-1}$
1114 with the land and ocean respectively over the final century of spin-up, well below the
1115 C4MIP target of 0.1 Pg C y^{-1} sought [Jones et al., 2016].

1116 In the preceding analysis of equilibration, the focus has largely concerned the ex-
1117 changes of carbon between the land, ocean and atmosphere components of the model.
1118 Table 1 presents the linear trends in ocean properties at different depth horizons for the
1119 final 500 y periods of both the ocean-only and fully-coupled spin-up phases. While carbon
1120 fluxes fall below C4MIP’s drift criterion (see Figures 5 and 6), it is clear that the ocean’s
1121 state is still drifting, and that these drifts have generally increased with the transition from
1122 the long duration ocean-only phase to the much shorter duration fully-coupled phase.

1123 As already noted, drifts in ocean temperature between these phases are a response,
1124 respectively, to a heat flux imbalance in ocean-only forcing, and the subsequent equilibrat-
1125 ing response when fully-coupled. Trends in nitrogen and iron nutrients have levelled off
1126 during the long ocean-only phase, but have grown into the fully-coupled phase as dust-
1127 forcing both changes and becomes more dynamic. Opposite sense trends result in these
1128 two nutrients, and are much larger in the upper ocean where the change in iron is affected.
1129 Meanwhile, although the air-sea flux continues to equilibrate, drift in the ocean’s surface
1130 carbon cycle is affected by a more dynamic hydrological cycle that increases surface alka-
1131 linity (tracking salinity; Jiang et al. [2014]), buffering higher DIC concentrations.

1132 In general, model drift is greater at the surface than at depth, although this varies
1133 between properties, most obviously dissolved oxygen. Here, surface concentrations are
1134 essentially controlled by the temperature-dependent solubility of this gas, while interior
1135 concentrations are affected by remineralisation of sinking organic material. As noted pre-
1136 viously, the fully-coupled phase has slightly lower production of organic material because
1137 of reduced dust deposition and greater iron limitation. In turn, this translates to elevating
1138 interior oxygen as less oxygen is consumed.

1139 Overall, Table 1 underscores the difficulty in equilibrating ESMs, especially where
1140 spin-up modes such as ocean-only incompletely capture the behaviour of the fully-coupled
1141 model.

1142 A number of lessons can be drawn from the experience of the spin-up of UKESM1.

1143 Inevitably, biases occur across the model components, but a particularly marked
1144 bias is that of the ocean's dissolved inorganic carbon pool. As illustrated in Figure 15,
1145 UKESM1 shows a general deficit in ocean DIC concentration, together with patterns of
1146 bias that align with those in nutrients and oxygen. Some of these biases stem from defi-
1147 ciencies in modelled circulation, but MEDUSA-2.1's biogeochemistry plays a key role in
1148 others. While some minor tuning of model parameters took place during the development
1149 of UKESM1, no tuning to improve these interior ocean biases was undertaken, principally
1150 because of the timescales necessary (simulated and wallclock) to equilibrate changes to
1151 identify improvement [Yool et al., 2013]. As noted earlier, there are offline and acceler-
1152 ated simulation modes that can assist with this, although none were mature enough within
1153 the infrastructure of UKESM1 to be used during CMIP6 preparations. As such, a key
1154 lesson, and future aspiration for UKESM1, is the adoption of techniques for more rapid
1155 model equilibration, to facilitate both the identification and tuning-out of such biases.

1156 Focusing on the component-only phases of spin-up, an obvious lesson lies in ensur-
1157 ing the interface exchanges between model components and the atmosphere are calculated
1158 in a manner consistent with that of the fully-coupled model. As the ocean results show,
1159 and drawing also from land-only preparations, inconsistency favours alternative steady
1160 states, with the potential to favour different evolutions of heat and carbon between the
1161 component-only and coupled configurations. Given the ultimate aim is a spun-up model
1162 state consistent with the coupled model, a requirement is that the component models be-
1163 have as close to the coupled model as possible. It is also worth remarking that, since we
1164 expected our ocean-only phase to differ from that of the coupled model because of the
1165 absence of ocean-atmosphere interactions, the source of the differences noted in the first
1166 ocean-only spin-up phase took time to be discovered. Ideally, the relationship between the
1167 fully-coupled and component-only versions of an ESM should be formally examined, both
1168 in terms of code and coupling (e.g. the same parameterisations being used with the same
1169 input properties in the same ordering, etc.), and in the resulting simulation dynamics.

1170 Another factor that our spin-up experience identified is the selection of the atmo-
1171 spheric forcing itself. Here, our ocean-only phases were driven using atmospheric forcing
1172 from periods of GC3.1 piControl simulation in which there was a top-of-atmosphere im-
1173 balance. This imbalance ($+0.15-0.2 \text{ W m}^{-2}$) is in the GC3.1 run throughout its piControl.
1174 This led to our ocean-only model consistently warming during spin-up, admittedly only a
1175 small absolute amount, but large enough that the final, fully-coupled spin-up phase could
1176 be seen reversing this in response. Since stable ocean heat content is one of the key tar-
1177 gets of spin-up, this points to the need for careful selection of atmospheric forcing, again,
1178 with the aim to be as consistent as possible with the atmosphere in the target coupled
1179 model. On a similar note, another feature of the atmospheric forcing used is its duration
1180 and character. Initial experience with limited-duration GC3.1 simulations found only mod-
1181 est variability in the ocean and atmosphere, both in terms of the absolute magnitude of
1182 variability and its temporal profile. Consequently, a relatively short, multi-decadal period
1183 was selected for use in ocean-only spin-up. However, as results shown here illustrate, the
1184 model clearly exhibits variability of much larger magnitude, and with much longer peri-
1185 ods, most clearly in UKESM1's Southern Ocean, where Drake Passage transport exhibits
1186 strong centennial-scale cycles. While a forced ocean-only model is unable to respond in
1187 the same way as the ocean in a fully coupled simulation because of the absence of ocean-
1188 atmosphere interactions, the short periods of forcing used here are not necessarily repre-
1189 sentative of what the full model can produce. Overall, an assessment of the flux biases of
1190 downward atmospheric forcing, and the role of slow timescale variability in atmospheric
1191 forcing on both, ocean- and land-only spin-up, requires further analysis.

1192 An item that is not apparent in the earlier description of UKESM1's spin-up was its
1193 interaction with the model's development cycle. Since UKESM1 includes numerous new
1194 model components and developments relative to its CMIP5 predecessor, HadGEM2-ES,
1195 it required a lengthy period of development. The timescales associated with CMIP6 and
1196 with the throughput of the fully-coupled model (approximately 4 simulated years per 1
1197 wallclock day) meant that development and spin-up necessarily occurred in parallel. While
1198 this meant that the spin-up was not "clean" (i.e. was not made using a single identical
1199 model throughout), and that it was not without inconsistency as problems were ironed
1200 out (e.g. the ocean-only bulk formulae issue), this mode of operation maximised the time
1201 available for spin-up. It avoided the need to wait for code freezing of a final version, and
1202 permitted the addition of features that would otherwise not have been included. A number

1203 of coupling interactions in UKESM1, in particular, became possible because of this flexi-
1204 ble approach to the development of UKESM1 from new components and its spin-up. The
1205 alternative approach of finalising first would necessarily have either delayed UKESM1's
1206 participation in CMIP6 or required the use of a less complete ESM.

1207 Of particular value in the development, tuning and spin-up of UKESM1 was the
1208 availability of the BGC-val evaluation suite [de Mora et al., 2018]. Focused on the ocean
1209 component, this tool automated the analysis of simulations, providing a range of plots cov-
1210 ering geographical, depth and globally-integrated properties, as well as comparisons with
1211 observational fields where available. Initially used on a run-by-run basis, BGC-val became
1212 invaluable in the intercomparison of multiple runs, and in monitoring the spin-up to iden-
1213 tify and avoid runtime or model bias problems. While most ESM groups will already have
1214 access to such tools because of the role they can play, we would encourage new entrants
1215 to the field to acquire (by adoption or development) such a tool.

1216 **5 Conclusions**

- 1217 • The UKESM1 model was spun-up using a combination of component-only phases
1218 for land and ocean, followed by a period of fully-coupled simulation
- 1219 • Component-only phases were spun-up under atmospheric forcing derived from sim-
1220 ulations of coupled climate precursors of UKESM1
- 1221 • Model states from parallel ocean (≈ 5000 year) and land (≈ 1600 year) spin-up branches
1222 were united with the atmosphere and, later, the full atmosphere chemistry and aerosol
1223 component (≈ 240 year)
- 1224 • The resulting pre-industrial control has a top-of-atmosphere heat balance of less
1225 than -0.09 W m^{-2} and net atmosphere-ocean and atmosphere-land CO_2 fluxes of
1226 less than 0.1 Pg C y^{-1}
- 1227 • Although equilibrated at global scale, analysis of land carbon fluxes indicated that
1228 regional shifts were significant, implying that longer spin-up periods are required to
1229 ensure regional as well as global equilibration
- 1230 • Issues encountered during spin-up included consistency of the interfaces of component-
1231 only models, the duration and variability of the atmospheric forcing, including its
1232 overall consistency with atmospheric forcing in the target coupled model, and the
1233 important role played by rapid-turnaround evaluation tools
- 1234 • While some tuning of UKESM1 was undertaken during spin-up, the slow turnover
1235 of the ocean component and conventional spin-up modes used here limited its scope,
1236 supporting the future tailoring of accelerated spin-up techniques to UKESM1 to re-
1237 duce ocean biases, as well as achieve better equilibration

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1256 **6 Code and data availability**

1257 All simulations used in this work were performed using version 10.9 of the UM,
1258 version 5.0 of JULES, NEMO version 3.6, CICE version 5.1.2 and OASIS3-MCT ver-
1259 sion 3.0. Model output from the NEMO ocean model was handled using the XML Input-
1260 Output Server (XIOS) library [Meurdesoif, 2013].

1261 Identifiers for simulations in the ocean branch of spin-up are listed in Table ??, with
1262 u-aw310 the final piControl simulation and the source of the majority of model output
1263 shown in this paper.

1264 All simulation data used in this study are archived at the Met Office and are avail-
1265 able for research purposes through the JASMIN platform (www.jasmin.ac.uk). For fur-
1266 ther details please contact UM_collaboration@metoffice.gov.uk referencing this pa-
1267 per.

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Property	Units	0 m	500 m	1000 m	2000 m
Temperature	$^{\circ}\text{C ky}^{-1}$	-0.004	-0.000	0.010	0.023
		0.464	0.390	0.243	-0.194
Salinity	PSU ky^{-1}	0.001	0.001	0.003	0.002
		0.110	0.130	0.037	-0.040
DIN	$\text{mmol N m}^{-3} \text{ky}^{-1}$	0.046	0.069	0.045	-0.007
		2.871	-2.024	-1.075	0.098
Silicic acid	$\text{mmol Si m}^{-3} \text{ky}^{-1}$	0.049	0.246	0.350	0.047
		0.813	-2.437	-4.965	-0.797
Iron	$\text{mmol Fe m}^{-3} \text{ky}^{-1}$	-0.001	-0.000	0.000	0.000
		-0.281	-0.109	-0.020	-0.004
DIC	$\text{mmol C m}^{-3} \text{ky}^{-1}$	0.215	-0.213	-1.818	-3.615
		4.993	-10.514	-5.758	-0.897
Alkalinity	$\text{meq m}^{-3} \text{ky}^{-1}$	0.328	0.406	0.446	-0.017
		11.609	6.884	-1.076	-2.720
Oxygen	$\text{mmol O}_2 \text{m}^{-3} \text{ky}^{-1}$	0.012	-0.221	-0.372	-0.245
		-1.902	12.298	9.173	9.379

1848 **Table 1.** Global mean drift rates for key ocean properties at different reference depths for the final 500 year
1849 periods of the ocean-only (upper row) and coupled (lower row) phases. Drift rates calculated as the linear fit
1850 across the final 500 year periods, and shown as ky^{-1} .

1851 **Figure 1.** Schematic diagram of the components of UKESM1 and the associated code structuring and
1852 coupling relationships. Circular arrows indicate couplings between closely associated component executables,
1853 while large arrows indicate coupling between separate component executables (principally the atmosphere-
1854 land and ocean).

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1855 **Figure 2.** Simplified schematic diagram of the main simulation branches involved in the spin-up of
1856 UKESM1 components, and their approximate durations. The uppermost branch centres on ocean spin-up,
1857 the middle branch on land spin-up, and the lower branch on atmospheric chemistry spin-up. Colours indicate
1858 distinct configurations. Branches effectively occurred in parallel, and the main lines of state sharing between
1859 branches are indicated in solid black arrows.

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1860 **Figure 3.** Time-series plots of the full spin-up period. Colours indicate different phases, with two ocean
1861 phases followed by a UKESM1-CN phase and then a full UKESM1 phase, prior to the start of the piControl.
1862 Solid lines indicate 30-year rolling averages of the properties, with the shaded areas denoting the corre-
1863 sponding 30-year range of annual averages. Row 1 shows the evolution of ocean-average volume and surface
1864 temperature. Row 2 shows the evolution of ice area in the northern and southern polar regions. Row 3 shows
1865 the evolution in circulation strength for the AMOC and Drake Passage. The time axis is indexed such that the
1866 end of spin-up (and the start of the piControl) is at zero, with total spin-up duration (per Table ??) indicated
1867 by the negative extent of the time axis.

1868 **Figure 4.** Time-series plots of the 200 y period after the ocean-only (blue) phase branches to start the cou-
1869 pled (red) phase. Panel ordering follows that of Figure 3. Row 1 shows the evolution of ocean-average volume
1870 and surface temperature. Row 2 shows the evolution of ice area in the northern and southern polar regions.
1871 Row 3 shows the evolution in circulation strength for the AMOC and Drake Passage. The time axis indicates
1872 the time (in years) since the branching occurs, with a preceding 50 year period with ocean-only phase only.

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1873 **Figure 5.** Following Figure 3, rows 1 and 2 show time-series plots of the full spin-up period, with colours
1874 indicating different spin-up phases. Solid lines indicate 30-year rolling averages of the properties, with
1875 the shaded areas denoting the corresponding 30-year range of annual averages. The panels show globally-
1876 integrated net primary production (Pg C y^{-1}) and globally-integrated net air-sea flux (Pg C y^{-1}). Following
1877 Figure 4, row 3 shows the corresponding time-series plots of the same properties for the final 100 y periods of
1878 the ocean-only (blue) and coupled (red) phases. The time axis shows both phases running in parallel whereas
1879 they ran in series.

1880 **Figure 6.** Time-series of globally-integrated land component properties during the land branch of spin-up.
1881 The upper two panels show soil and vegetation carbon (in Pg C), while the lower three panels show the frac-
1882 tional cover of total land area associated with tree, grass and bare soil. Gaps in the time-series were caused
1883 by data archiving failures. The uppermost panels include a grey zone to indicate C4MIP's "drift cone" of
1884 0.1 Pg C y^{-1} Jones et al. [2016]. The numbers indicated with "#" are referenced in the text and Table ?? . The
1885 period shown follows on from after the initial land-only spin-up phase, using UKESM1-CN (from -865 y; #1)
1886 and then UKESM1 (from -210 y) prior to starting the piControl at 0 y.

1887 **Figure 7.** Time-series of soil carbon integrated across biomes for the final spinup and piControl. In order to
1888 illustrate the scale of drift, the panels include a grey zone that indicates the 250 year period of the first Histori-
1889 cal ensemble member (1850–2100). Additional branch dates for subsequent Historical ensemble members are
1890 indicated by dashed lines. The period shown is from the final period of UKESM1 spin-up (-110 y to 0 y) prior
1891 to start of the piControl (0 y onwards).

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1892 **Figure 8.** Time-series covering the final 500 years of the UKESM1 coupled spin-up and the first 200 years
1893 of the coupled UKESM1 piControl. The figure plots: years -500 to -200 UKESM1-CN (run IDs; u-ar783
1894 and u-au835) followed by, years -200 to 0 UKESM1 (run IDs; u-av472, u-av651 and u-aw310), followed by,
1895 years 1 to 200 UKESM1 piControl. Panel 1 shows global mean top of atmosphere (TOA) net downward short
1896 wave (SW) radiation. Panel 2 shows the corresponding global mean TOA outgoing long wave (LW) radiation.
1897 Panel 3 shows the resulting balance of global mean TOA net radiation. Panel 4 shows global mean 1.5 m air
1898 temperature. Finally, panel 5 shows Arctic (blue) and Antarctic (black) sea-ice extent. In each panel, thick
1899 lines are an 11-year running mean, and the thin lines are annual mean values. Radiation values are in W m^{-2} ,
1900 with positive values indicating a downward-directed flux for net SW down and net radiation, and an upward-
1901 directed flux for outgoing LW. Temperature is in units of degrees Kelvin, and sea-ice extent expressed as 10^6
1902 km^2 .

1903 **Figure 9.** Observational (HadISST) and simulated sea surface temperature for northern (top; JJA) and
1904 southern (medium; DJF) summer. Differences (simulated - observed) for both seasons shown in bottom row.
1905 Temperature (and difference in temperature) in °C. HadISST data from the period 1870-1879.

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1906 **Figure 10.** Observed (red; HadISST) and simulated (blue) seasonal cycles of sea-ice extent (> 15% cover)
1907 for the Arctic (left) and Antarctic (right). Sea-ice extent in 10^6 km². HadISST data from the period 1870-
1908 1879.

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1909 **Figure 11.** A “thermohaline circulation” section of biases in modelled potential temperature (top) and
1910 salinity (bottom). The section tracks southwards “down” the Atlantic basin from the Arctic to the Atlantic
1911 sector of the Southern Ocean, before tracking northwards “up” the Pacific basin from the Pacific sector of the
1912 Southern Ocean to the Bering Straits. The aim is to capture the stereotypical transport of deep water from its
1913 formation as a “young” water mass in the high North Atlantic through to end as an “old” water mass in the
1914 North Pacific. Dotted lines mark the “boundaries” of the Southern Ocean at 40°S in each basin. Biases in
1915 potential temperature are in °C, and in practical salinity units (PSU) for salinity. Observational data from the
1916 World Ocean Atlas climatology for the period 1995-2004.

1917 **Figure 12.** Observationally-derived (top) and simulated (bottom) meridional overturning circulation
1918 (MOC) for the global ocean. The model circulation shown is based on the decadal-averaged streamfunc-
1919 tion. MOC in Sv, with both plots including Gent-McWilliams components [Gent and McWilliams, 1990].
1920 Observational data from the ECCO V4r4 ocean circulation reanalysis for the period 1992-2017.

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1921 **Figure 13.** Annual average observational (left) and simulated (right) fields of surface dissolved inorganic
1922 nitrogen (top; mmol N m^{-3}), total surface chlorophyll (middle; mg m^{-3}), and vertically-integrated net pri-
1923 mary production (bottom; $\text{g C m}^{-2} \text{d}^{-1}$). Note that total surface chlorophyll is shown on a logarithmic scale.
1924 Observational data are from the World Ocean Atlas (DIN; climatology from 1981-2010), SeaWiFS (chloro-
1925 phyll; climatology for the period 2000-2009) and the VGPM, Eppley-VGPM and CbPM products (NPP;
1926 climatology for the period 2000-2009).

1927 **Figure 14.** Observational (left) and simulated (right) annual average surface dissolved inorganic carbon
1928 (top) and total alkalinity (bottom). DIC in mmol C m^{-3} , alkalinity in meq m^{-3} . Observational data are
1929 from the GLODAPv2 climatology, from the pre-industrial period for DIC, and normalised to year 2002 for
1930 alkalinity.

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1931 **Figure 15.** A “thermohaline circulation” section of observed (top) and modelled (bottom) zonal average
1932 dissolved inorganic carbon. Figure 11 explains the format of this section. Concentrations in mmol C m^{-3} .
1933 Observational data are from the GLODAPv2 climatology, for the pre-industrial period.

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1934 **Figure 16.** Observationally-derived geographical map of major land biomes (top), together with a compari-
1935 son of the land cover type found in each biome for the model and two observational products, IGBP and CCI
1936 (bottom). Each biome appears as a separate triplet of bars, with the colour composition of the bar relating to
1937 the vegetation cover types indicated in the key. The observational IGBP product is derived from year 1992
1938 data, while the CCI product is derived from year 2000.

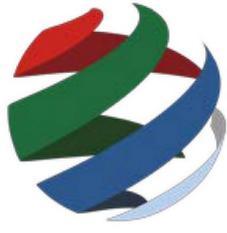
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1939 **Figure 17.** Geographical maps of fractional cover associated with each land cover type for the model (left)
1940 and two observational products, IGBP and CCI (middle and right, respectively). In each case, increasing
1941 colour intensity denotes greater fractional cover. The observational IGBP product is derived from year 1992
1942 data, while the CCI product is derived from year 2000.

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Figure 1.

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Unified Model atmosphere
UKCA-GLOMAP chemistry & aerosol

JULES
land physics
JULES
biogeochemistry
TRIFFID
dynamic vegetation

OASIS3-MCT
coupler

CICE
sea-ice
NEMO
ocean physics
MEDUSA
marine biogeochemistry

Figure 2.

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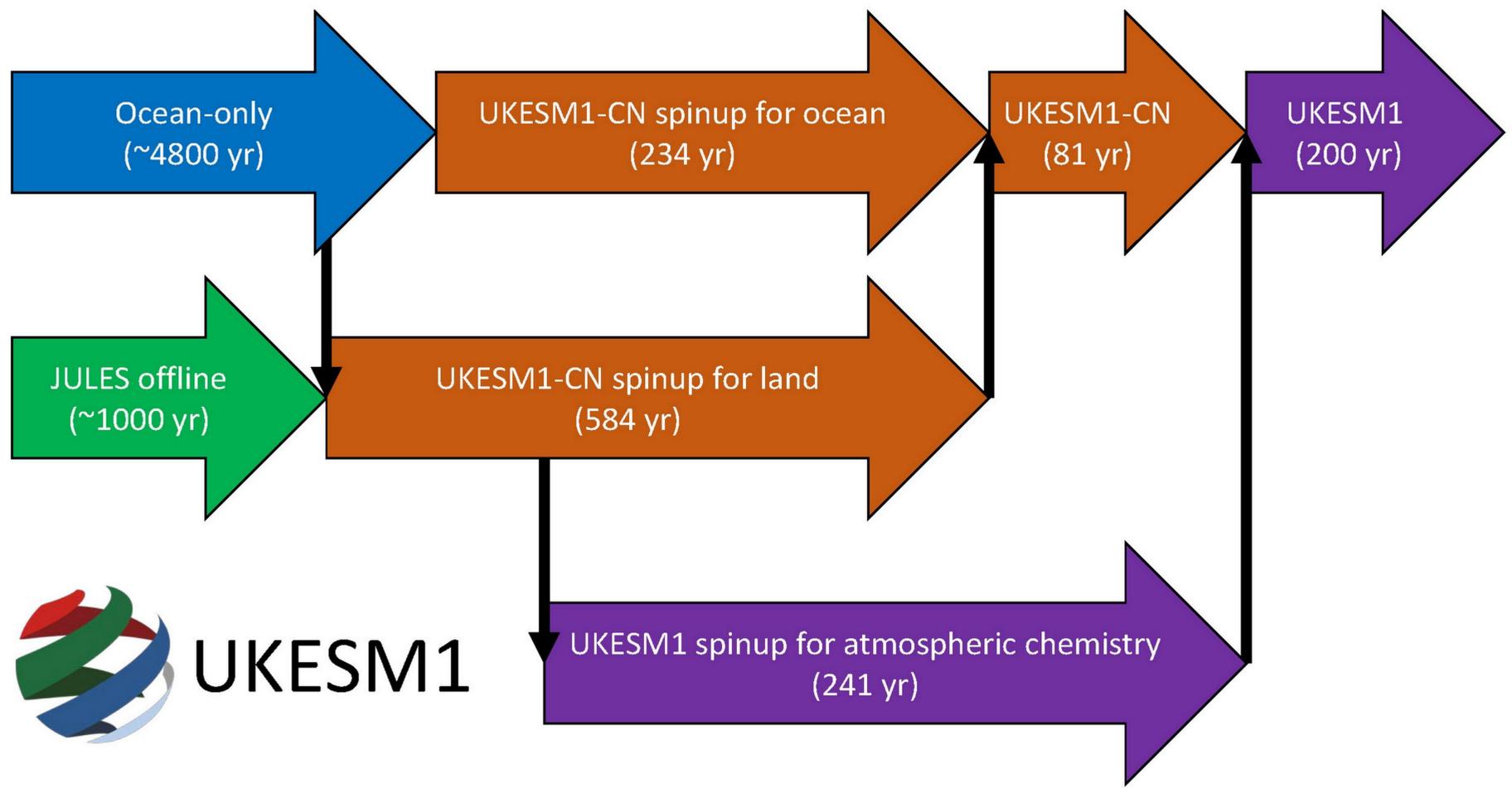


Figure 3.

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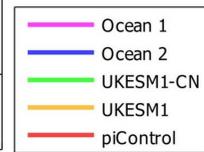
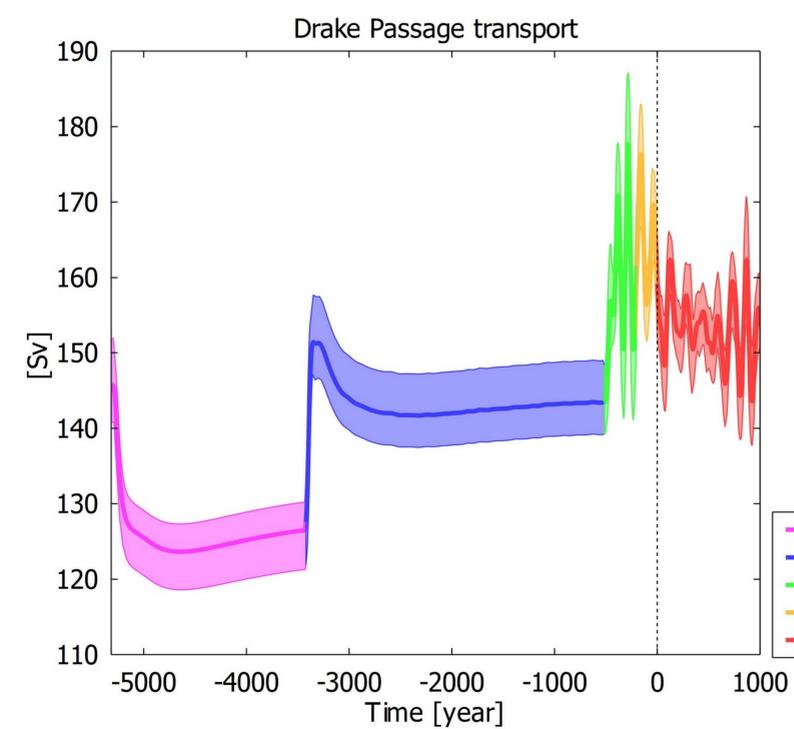
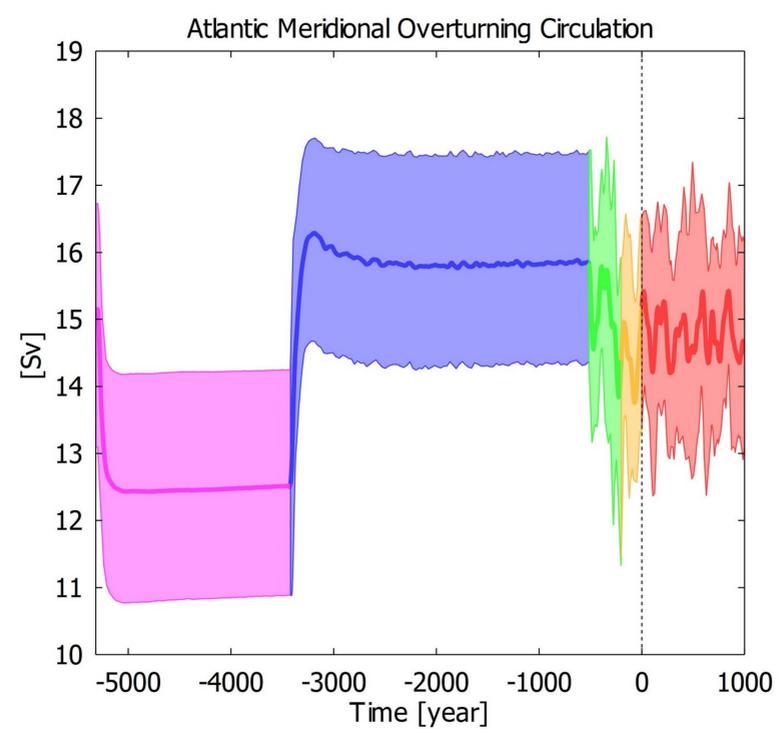
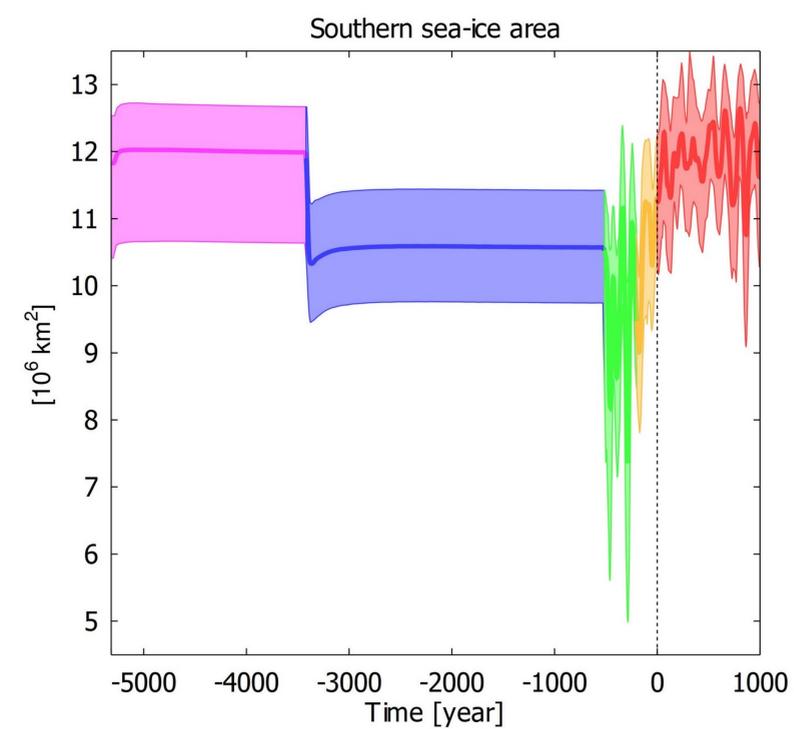
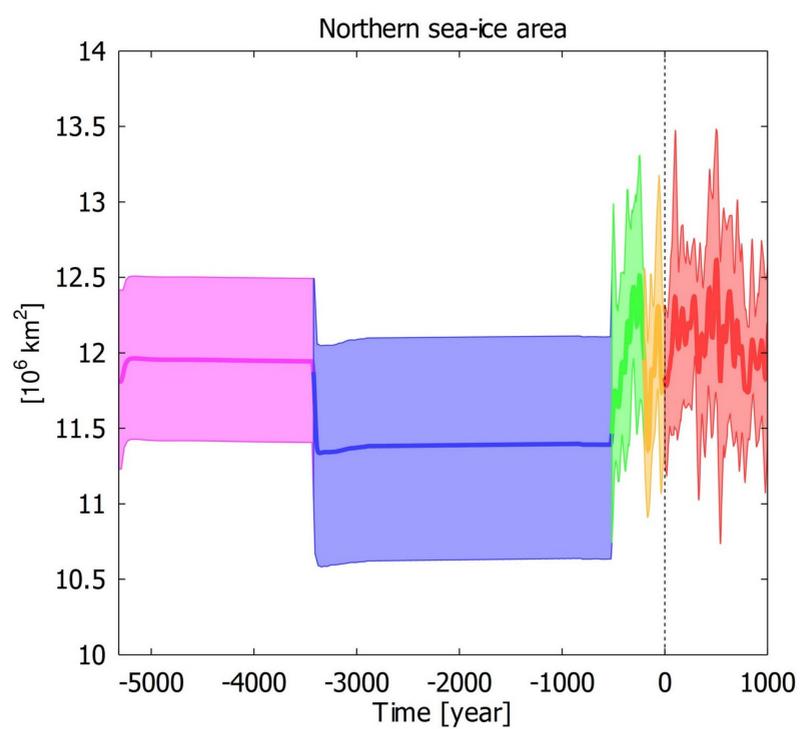
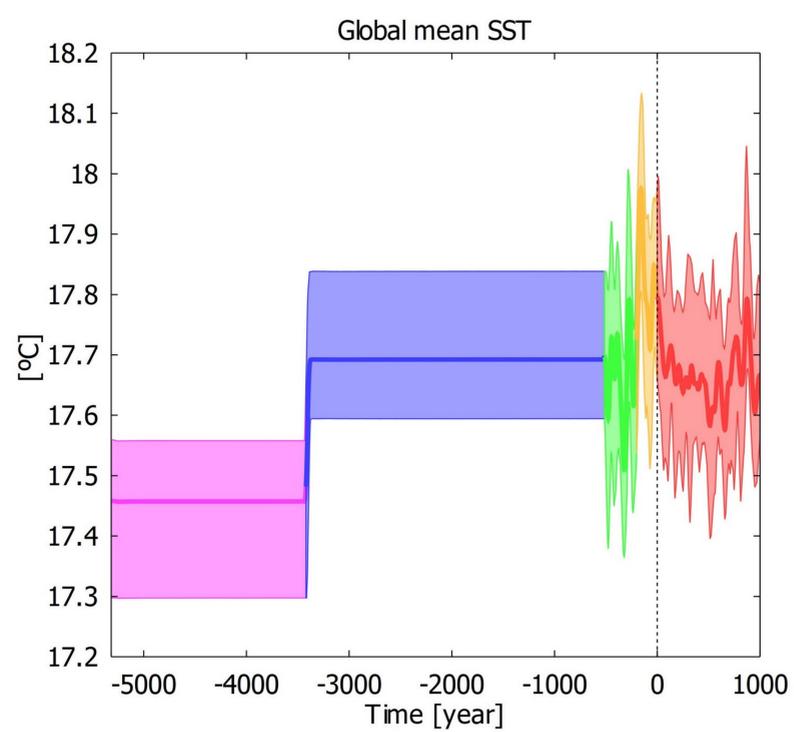
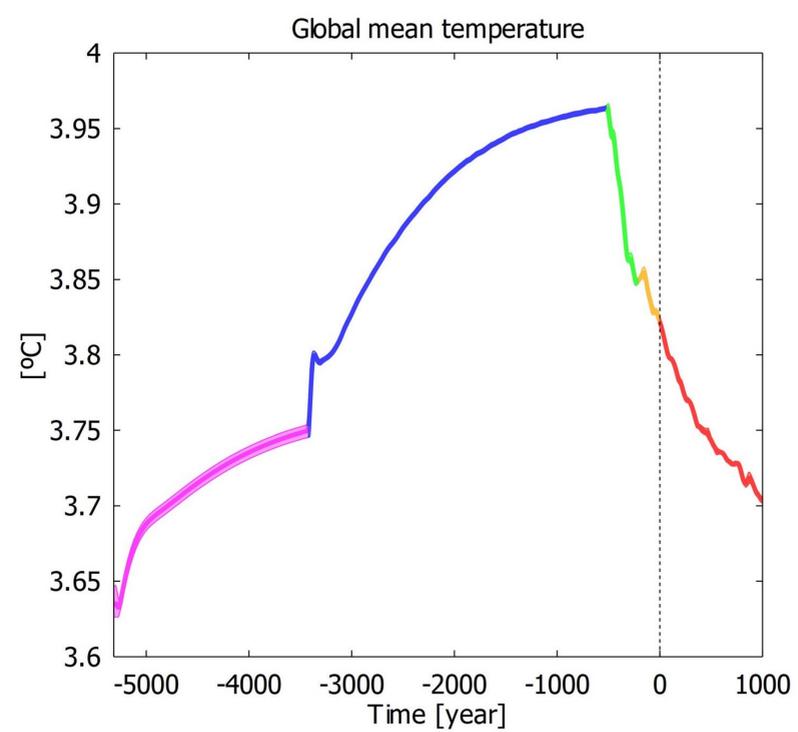


Figure 4.

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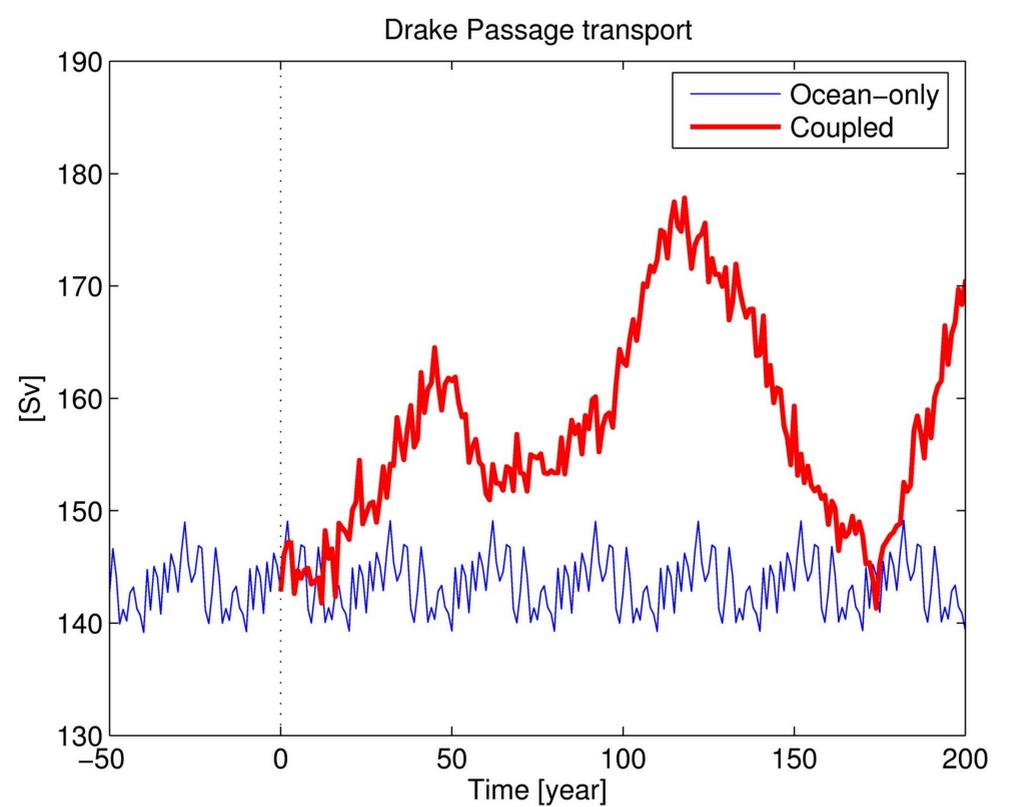
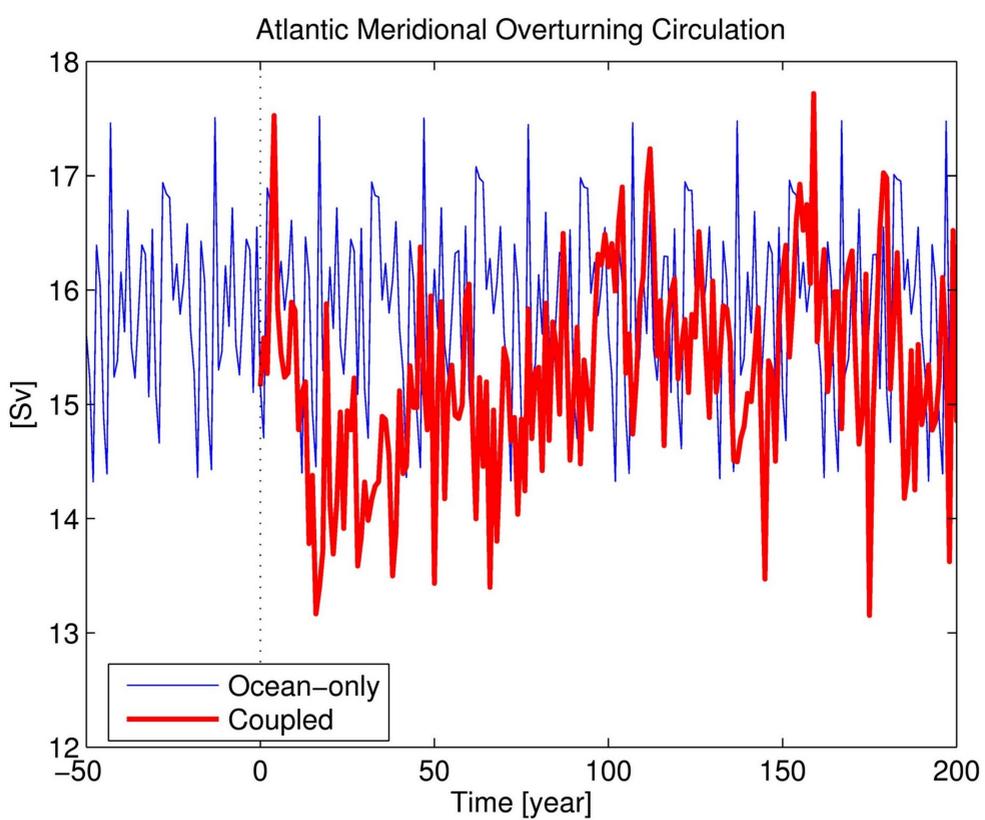
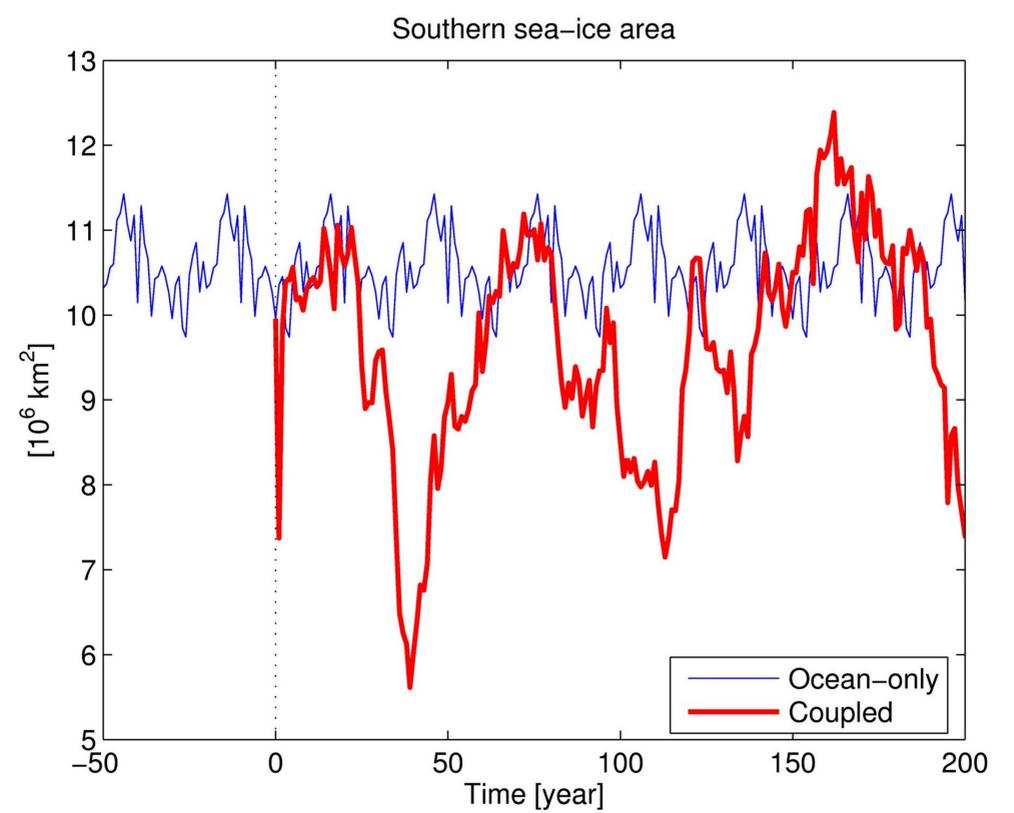
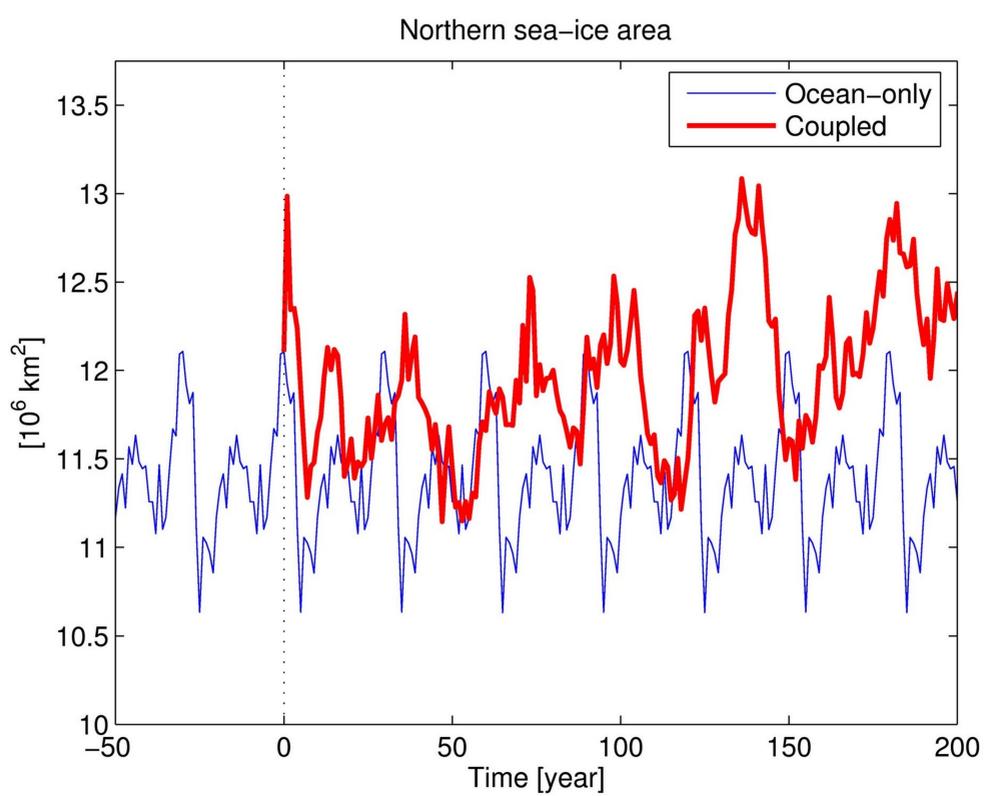
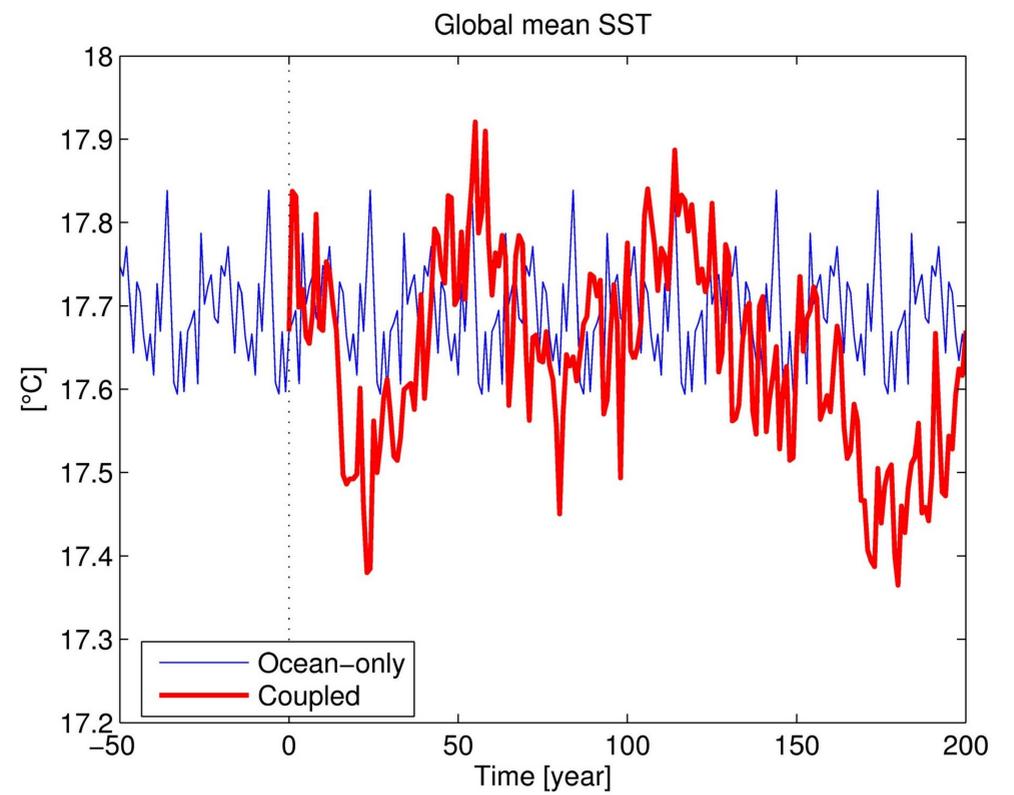
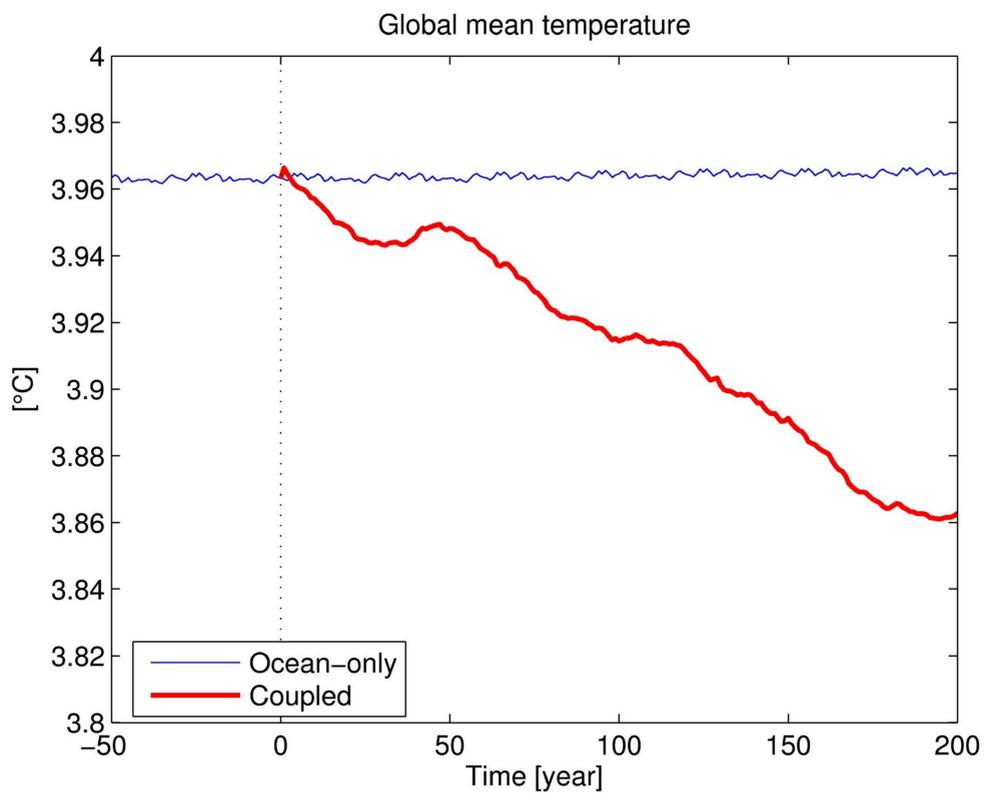


Figure 5.

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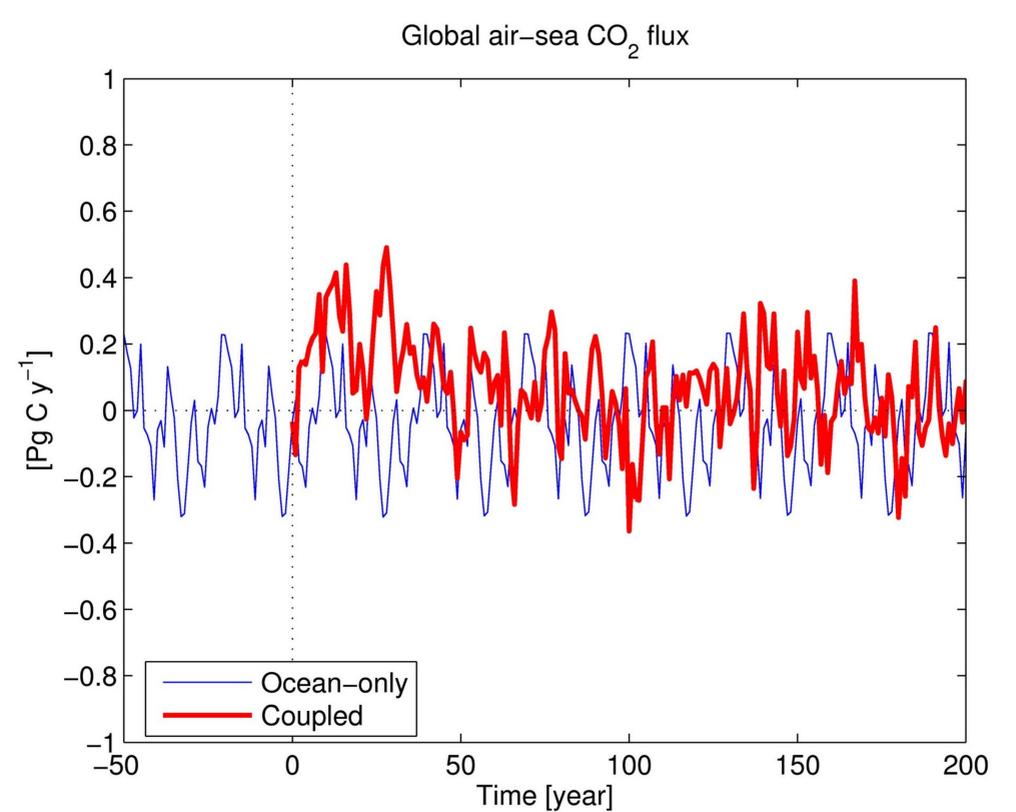
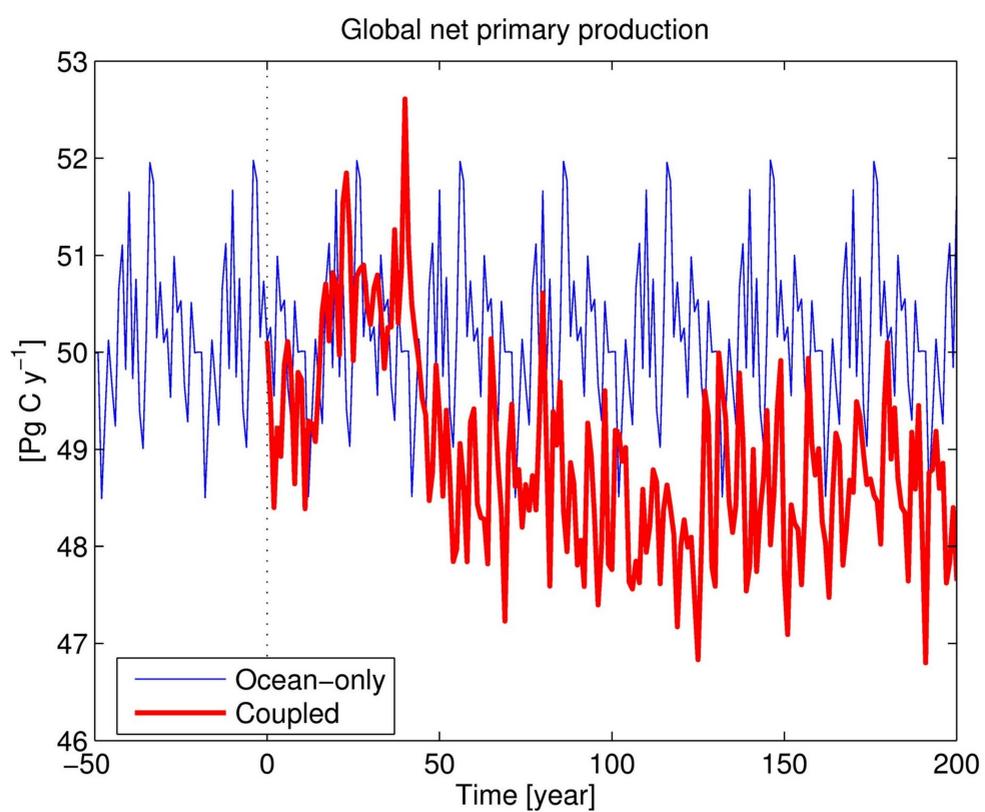
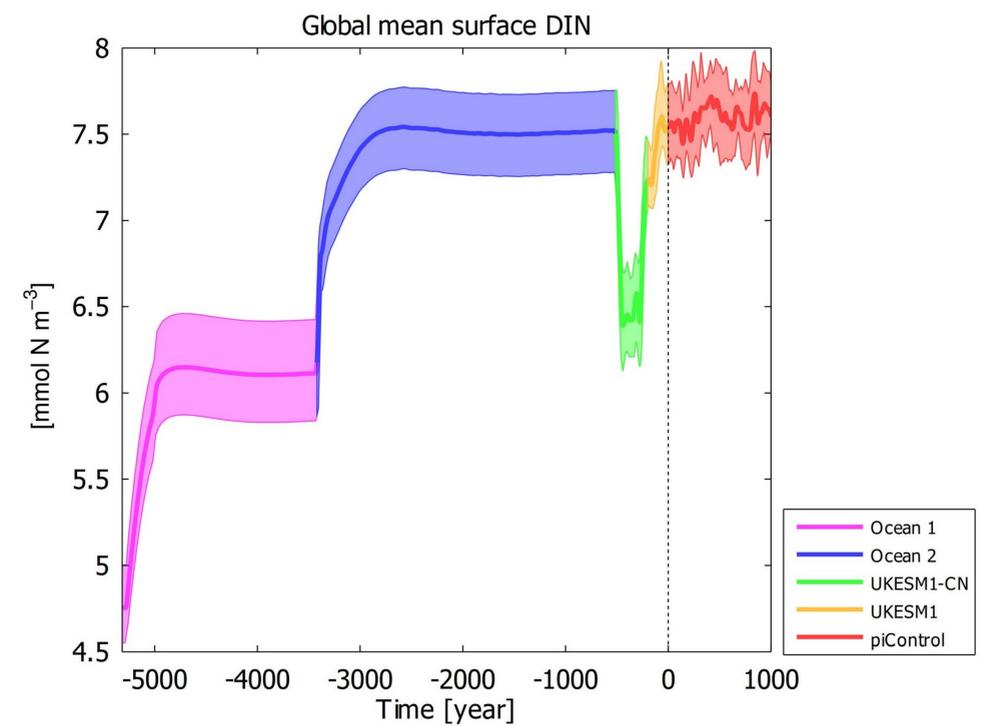
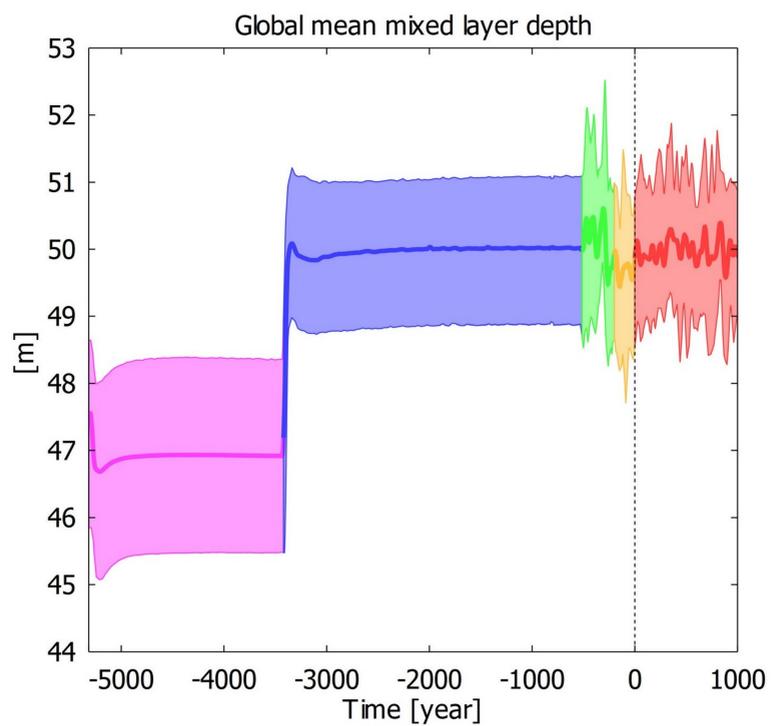
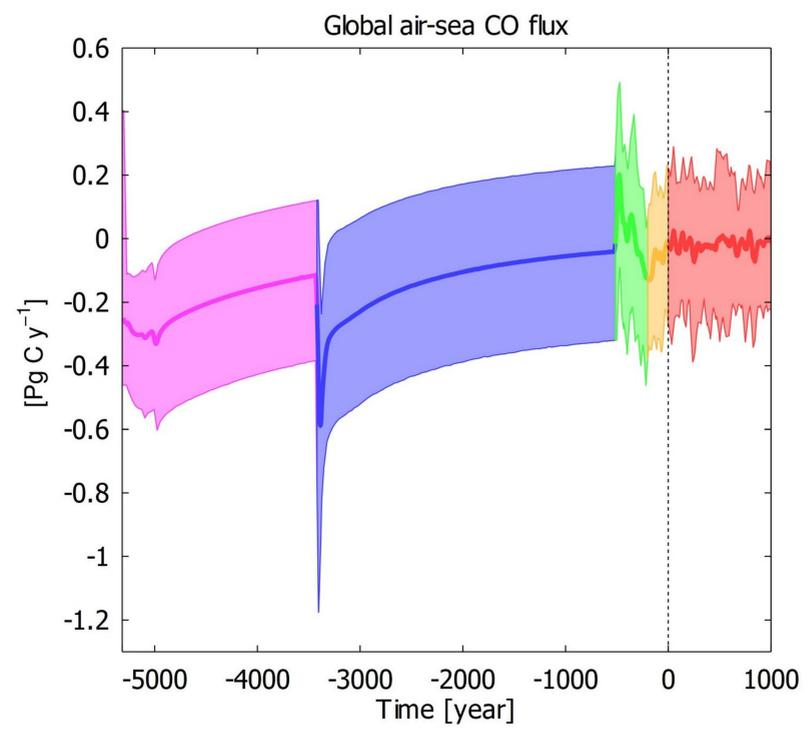
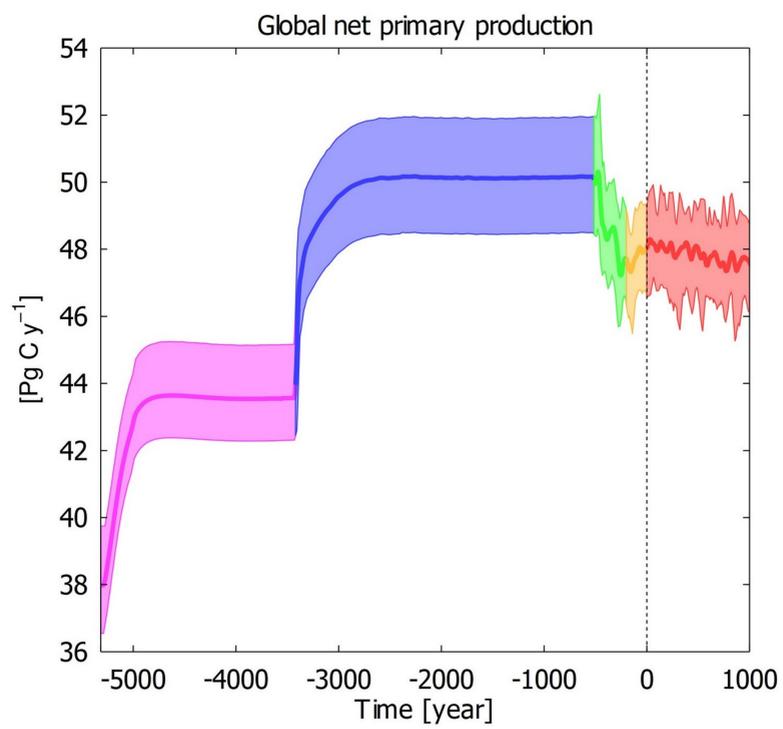


Figure 6.

Accepted Article

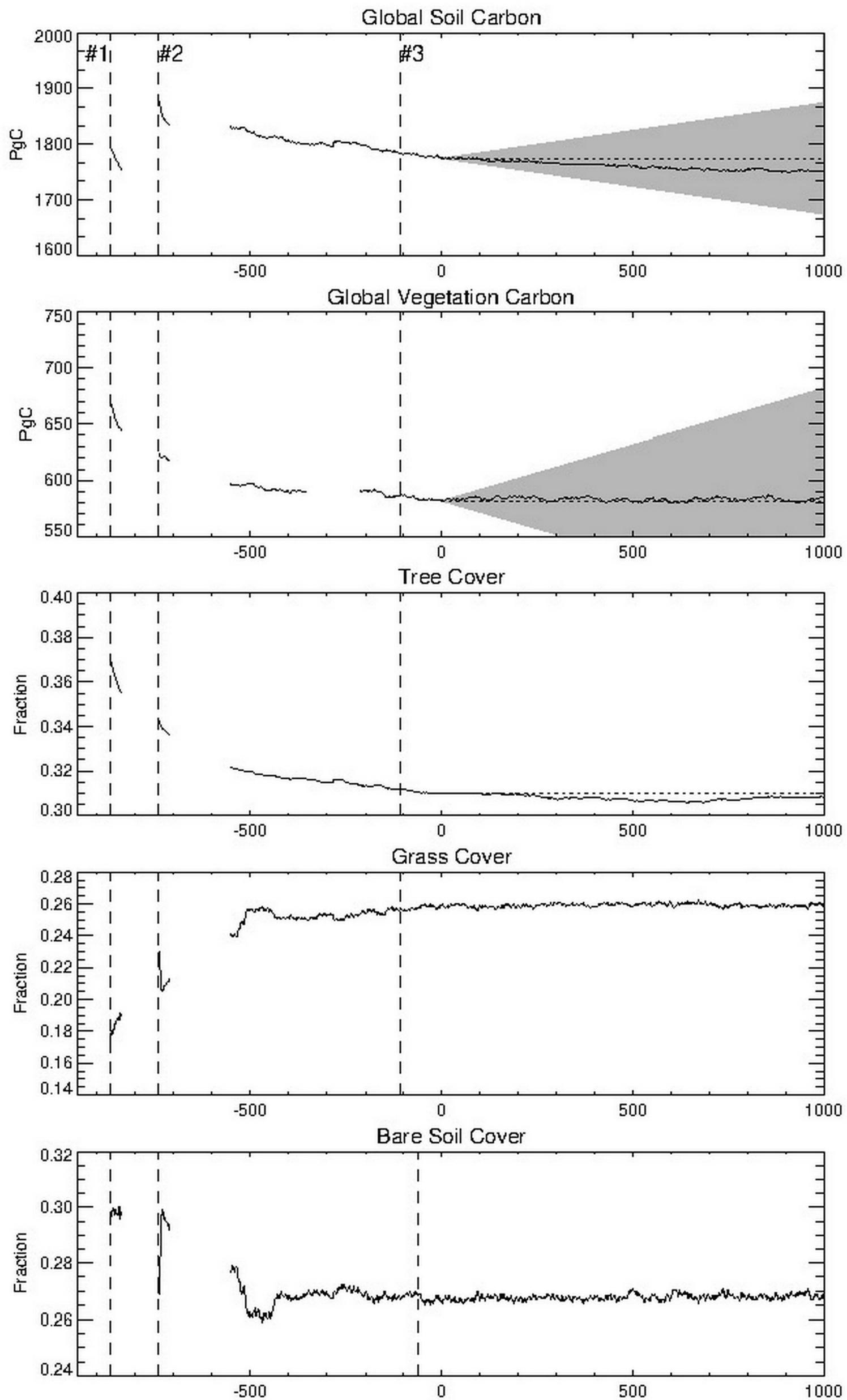


Figure 7.

Accepted Article

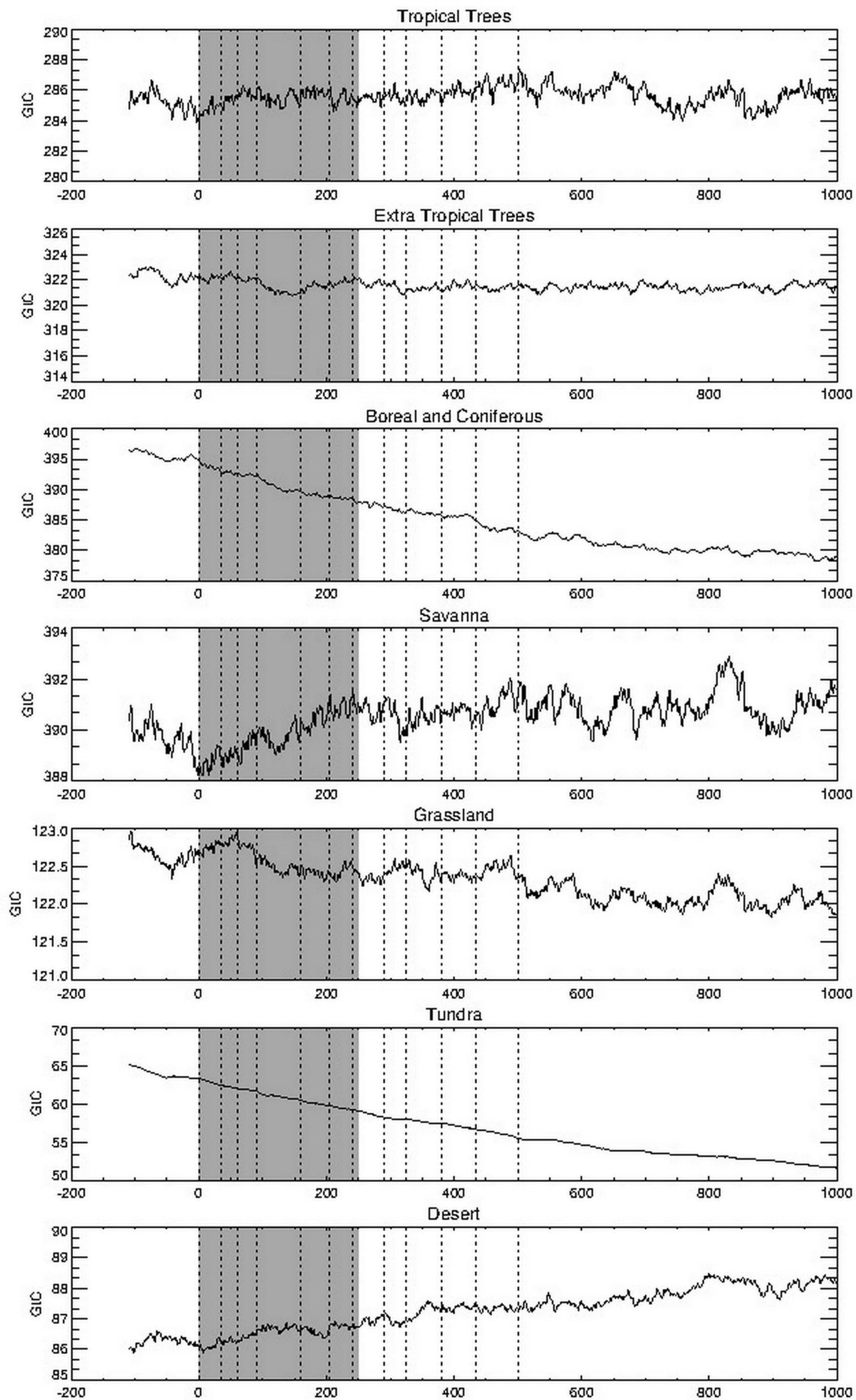


Figure 8.

Accepted Article

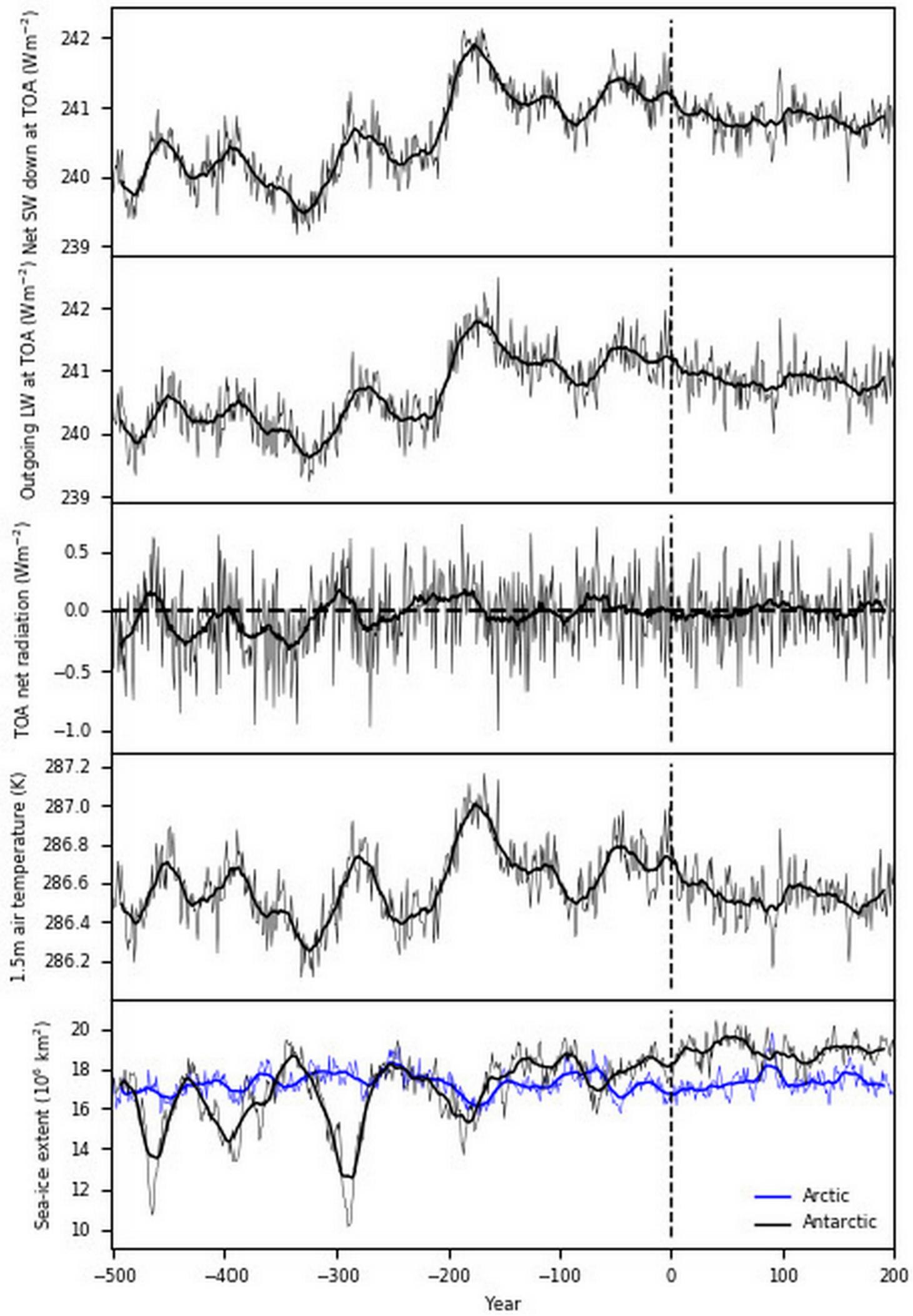
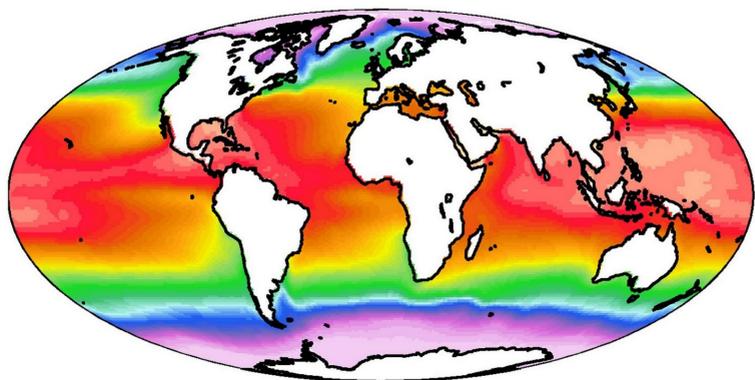


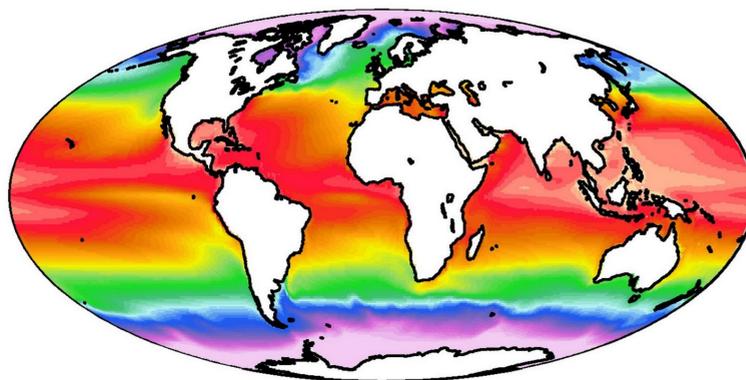
Figure 9.

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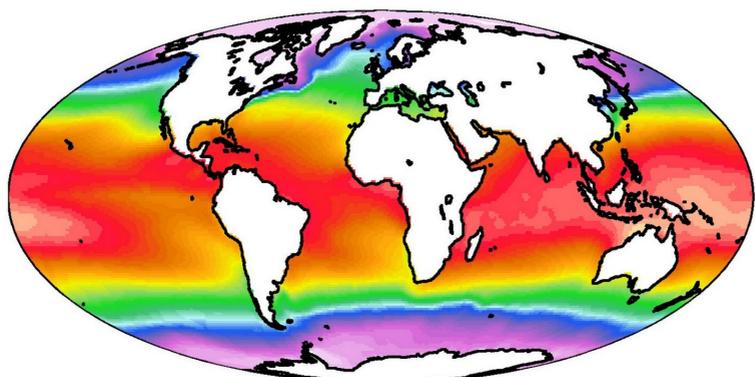
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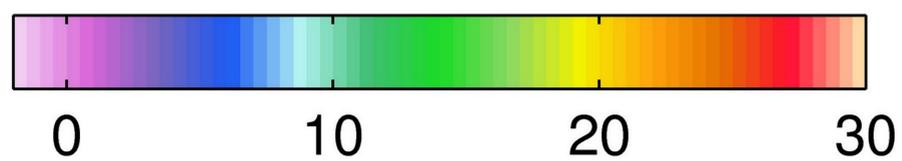
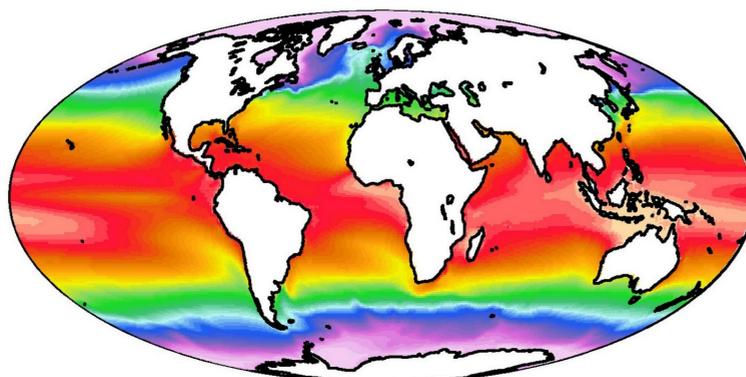
Simulated, JJA



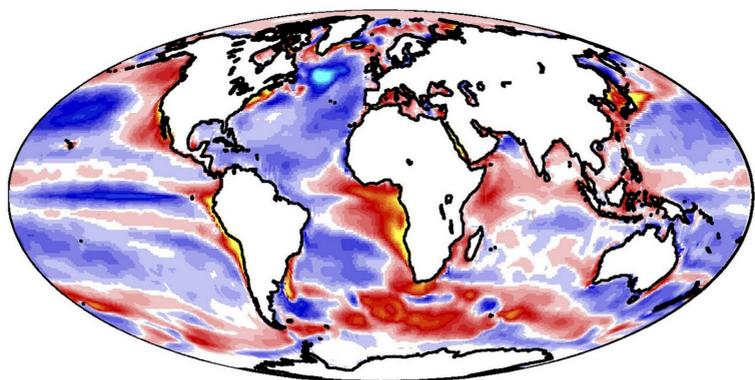
Observed, DJF



Simulated, DJF



Difference, JJA



Difference, DJF

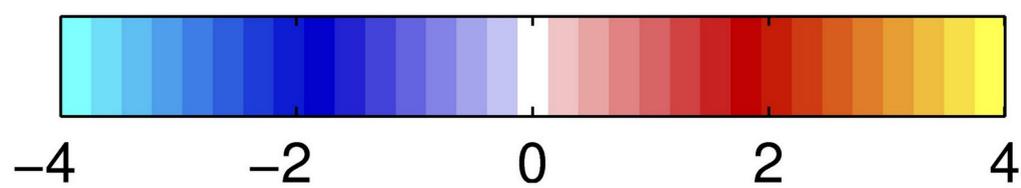
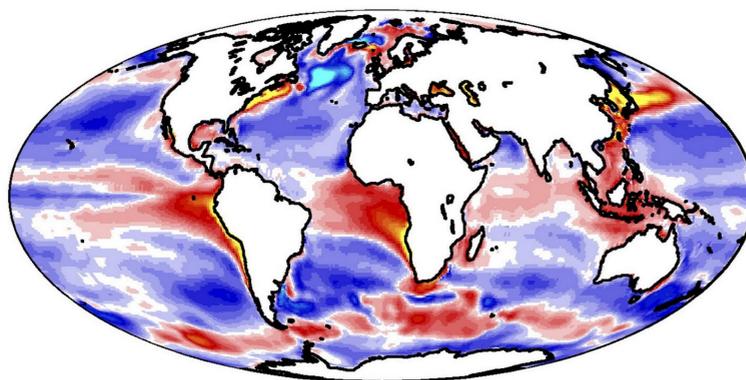
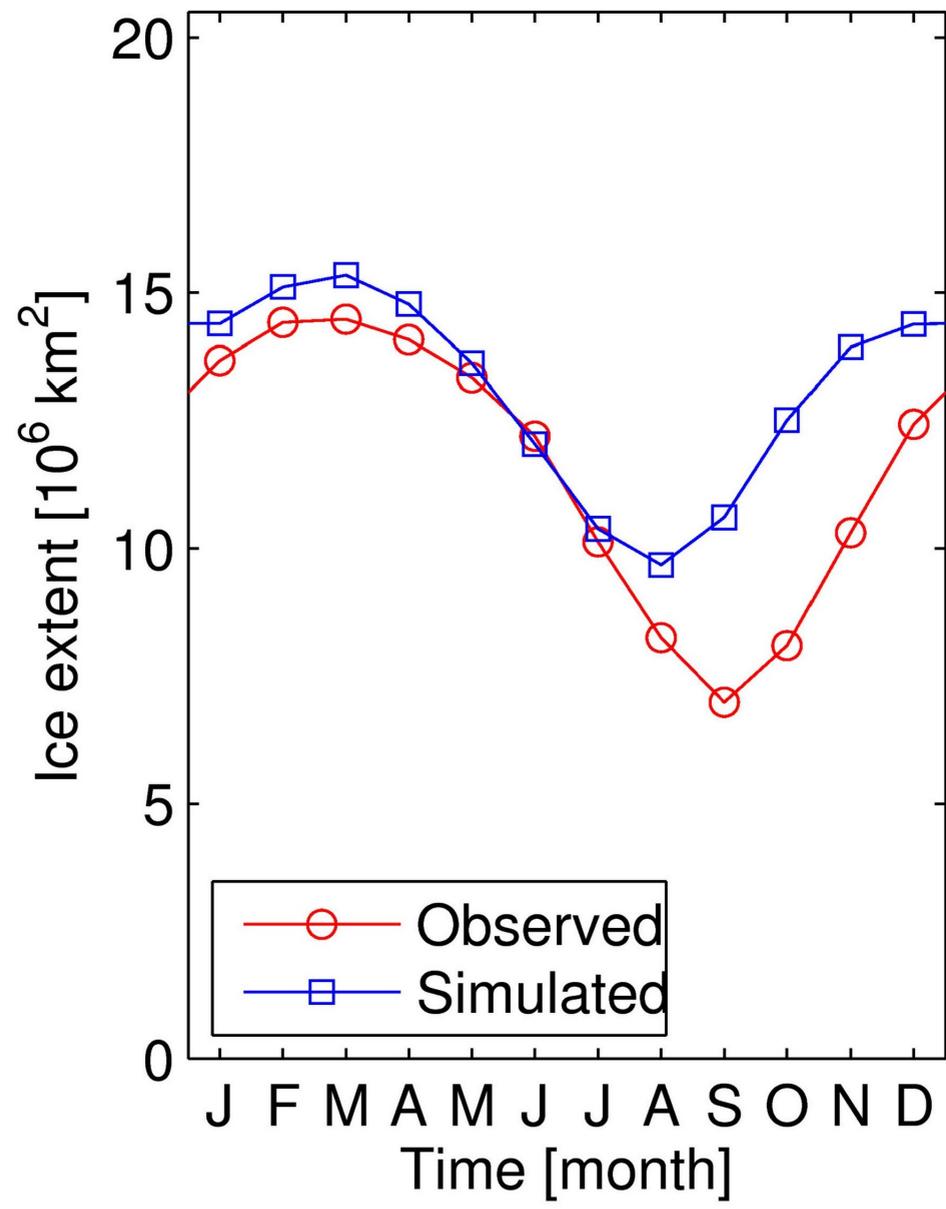


Figure 10.

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Arctic



Antarctic

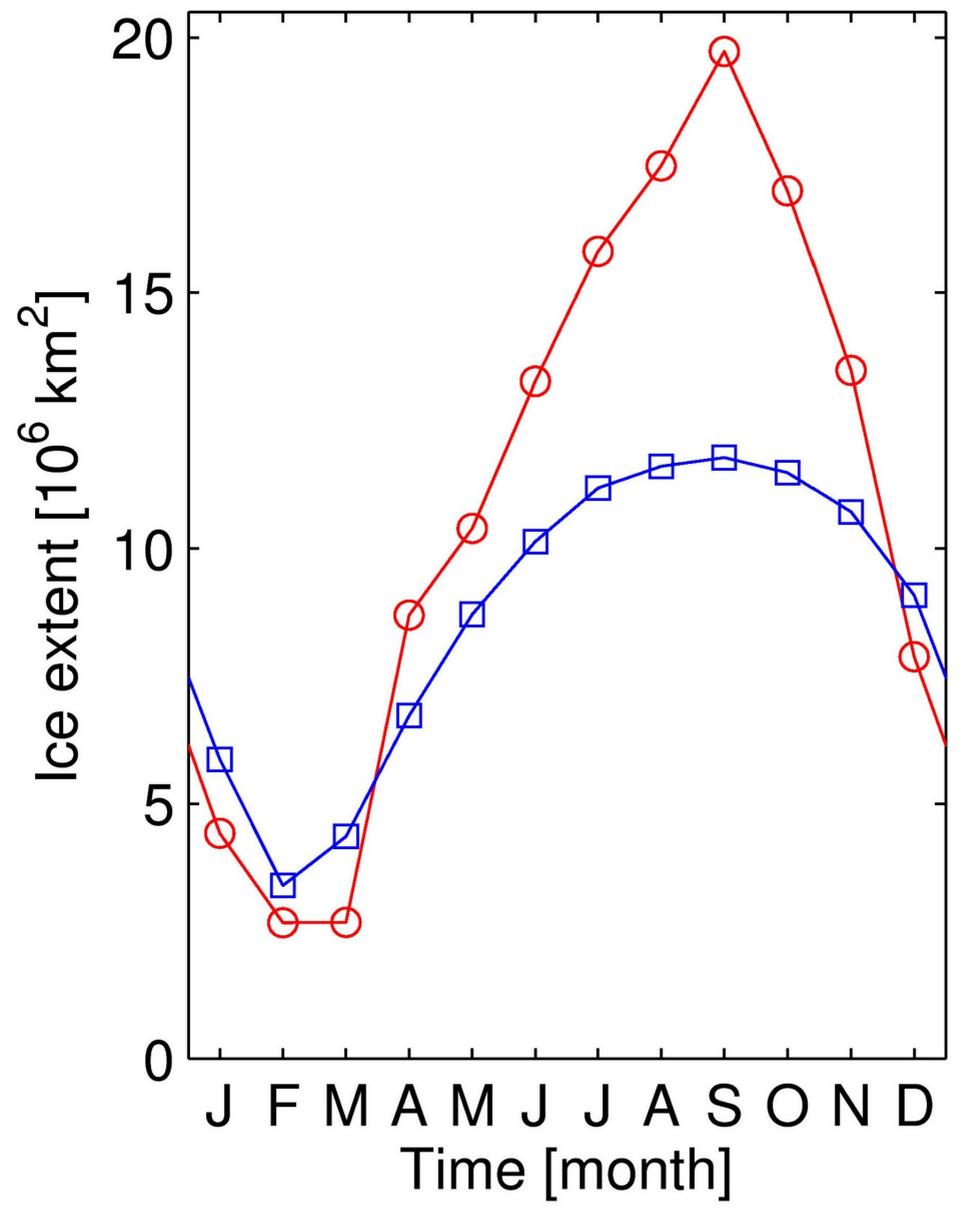
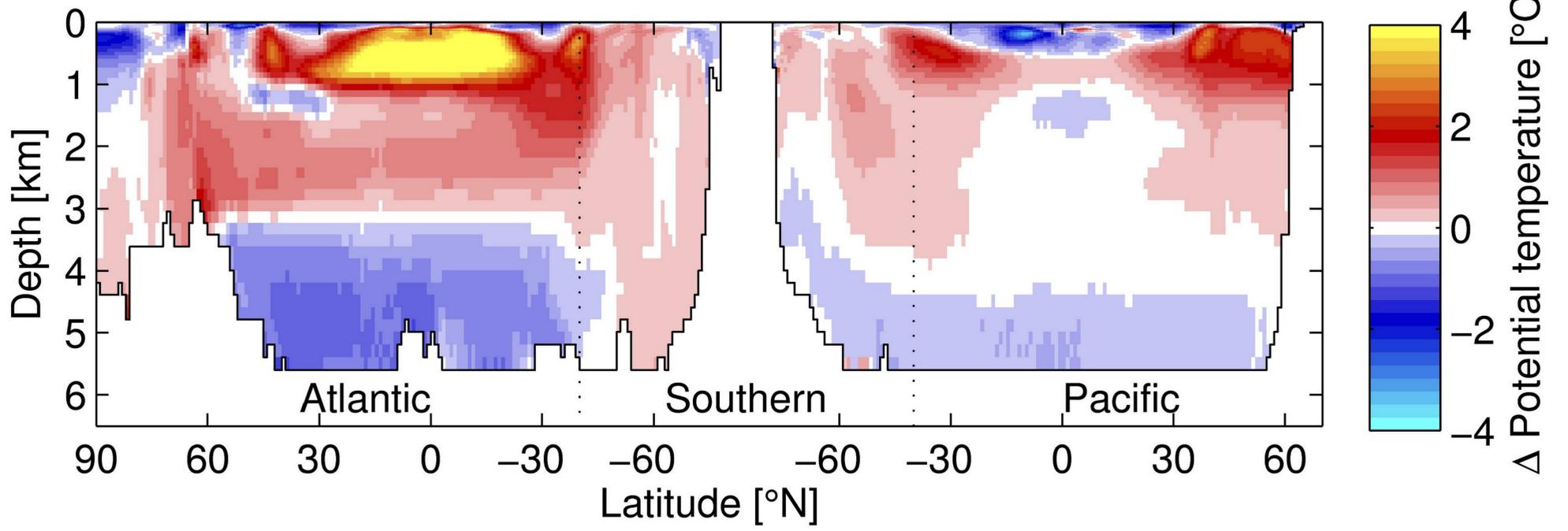


Figure 11.

Accepted Article

Difference Atlantic–Pacific zonal mean



Difference Atlantic–Pacific zonal mean

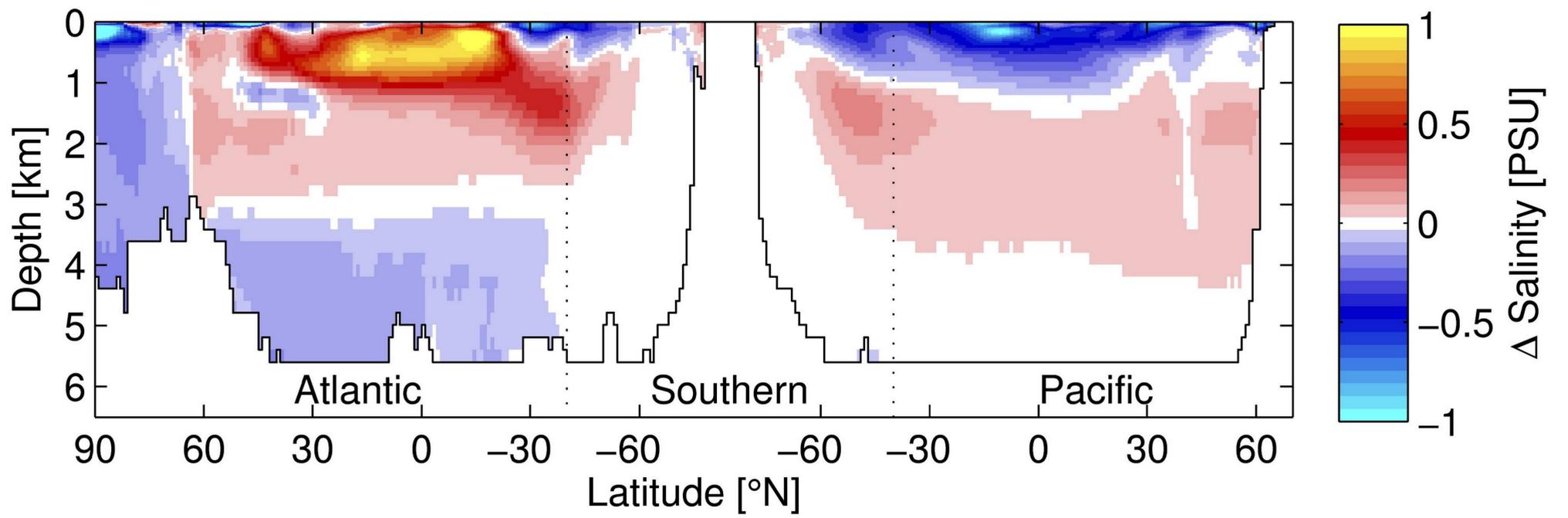
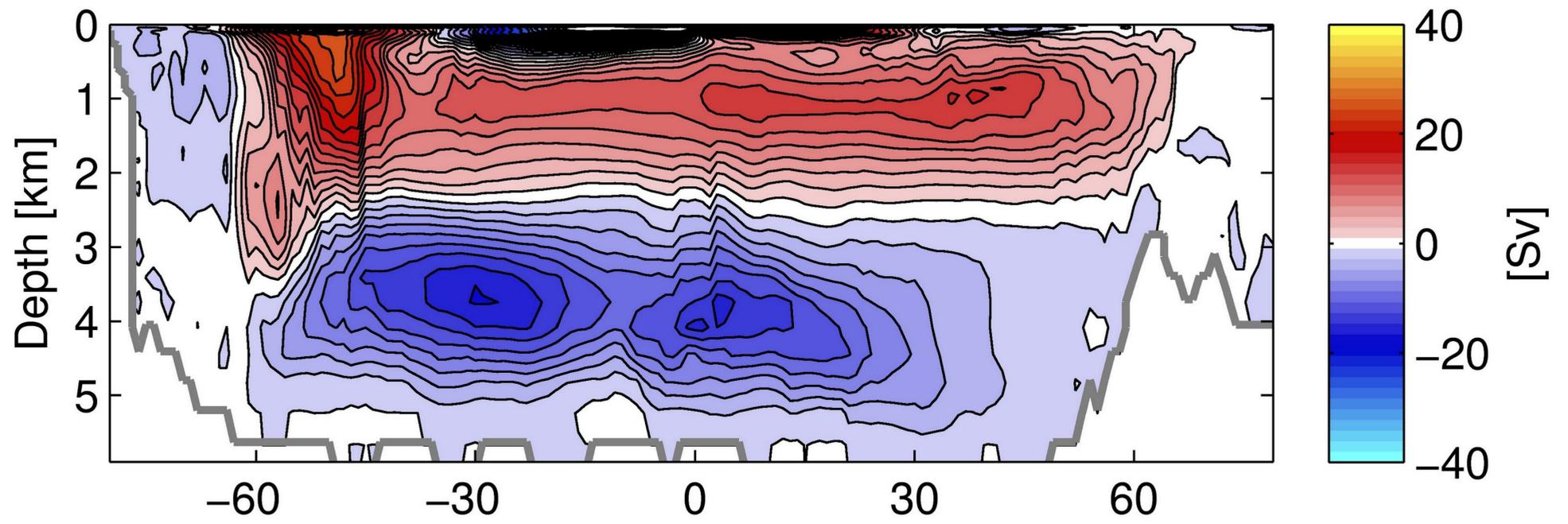


Figure 12.

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Global meridional streamfunction, observed



Global meridional streamfunction, UKESM1

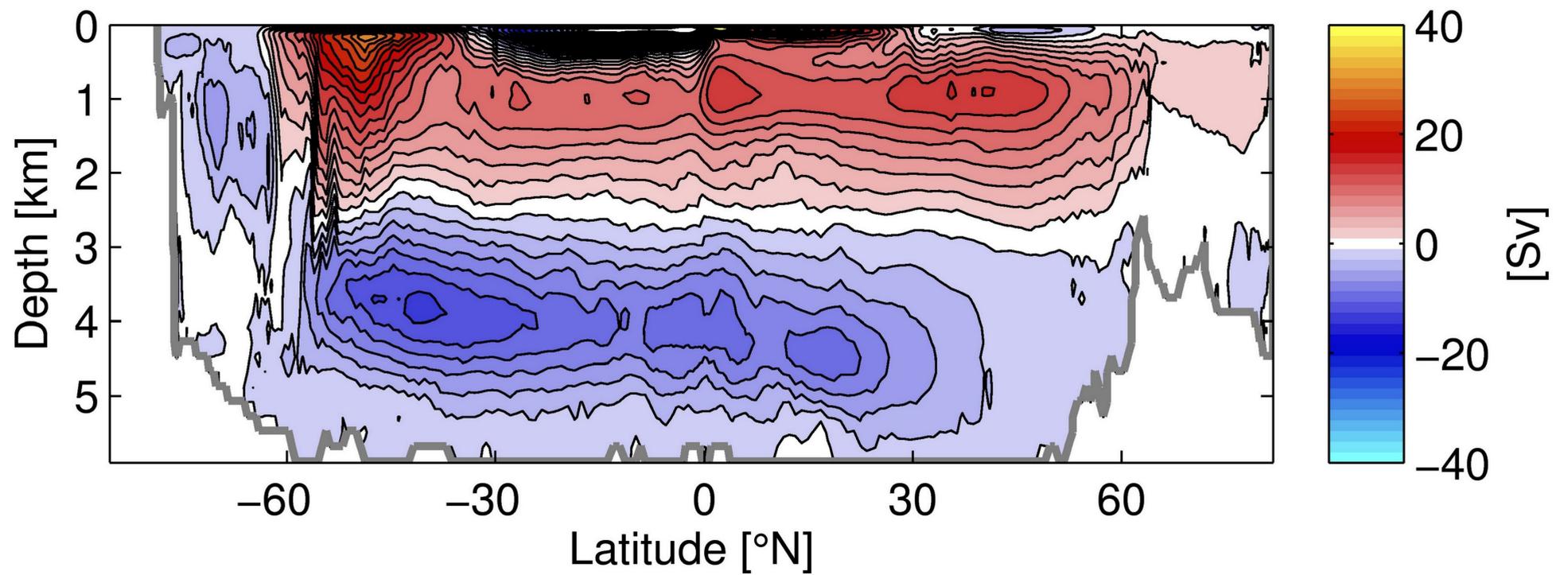
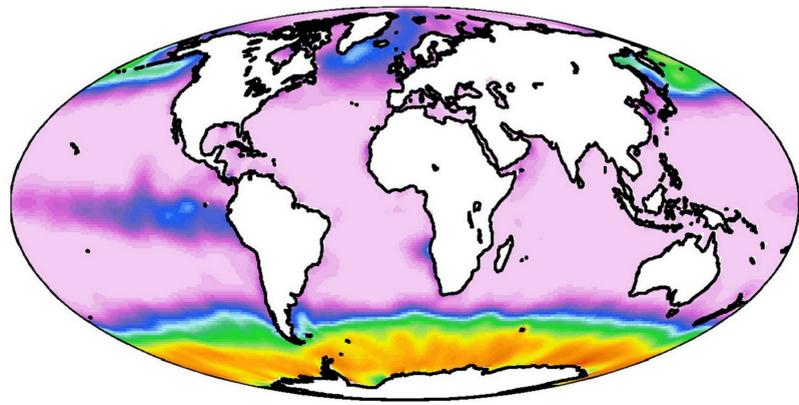


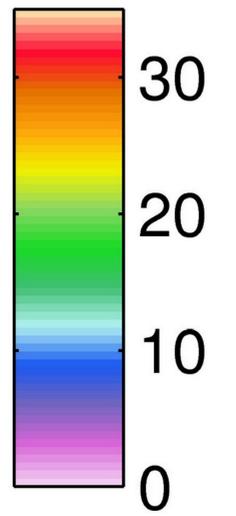
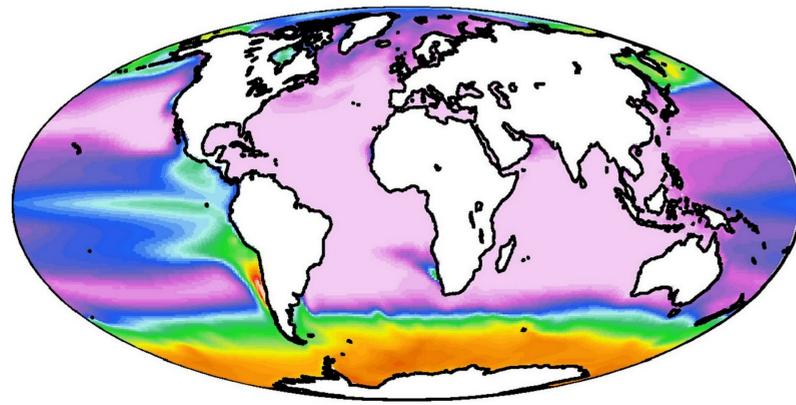
Figure 13.

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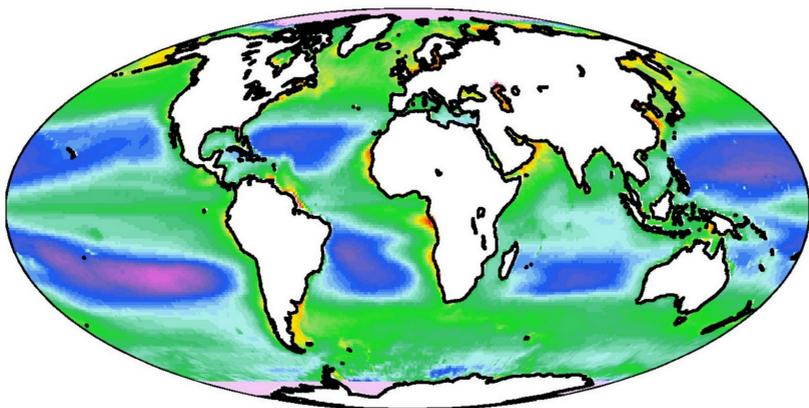
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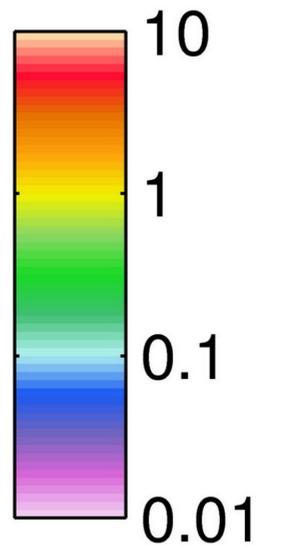
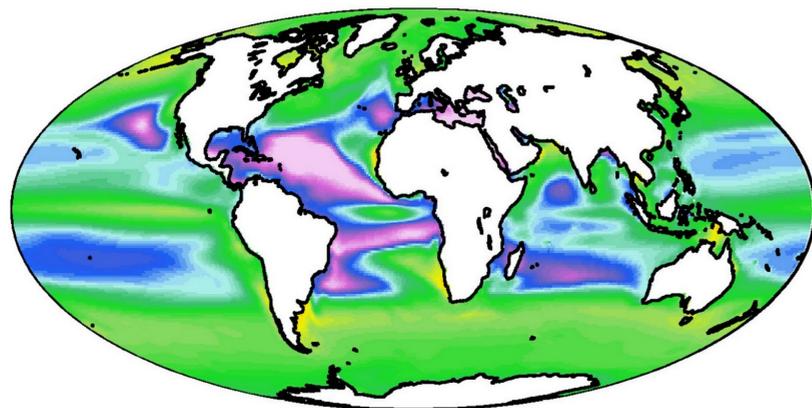
Simulated



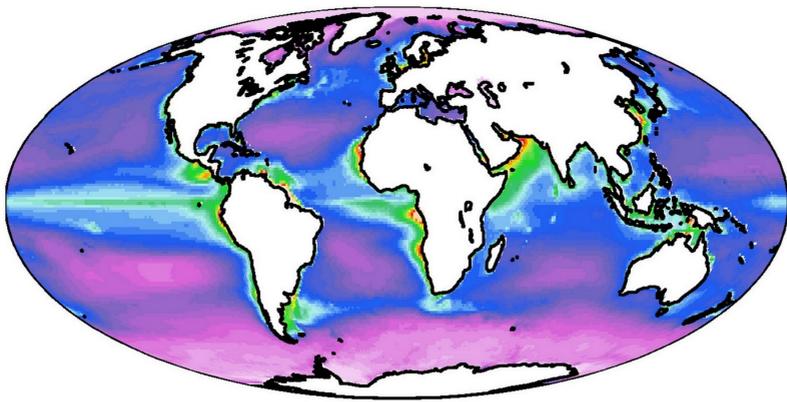
Observed



Simulated



Observed



Simulated

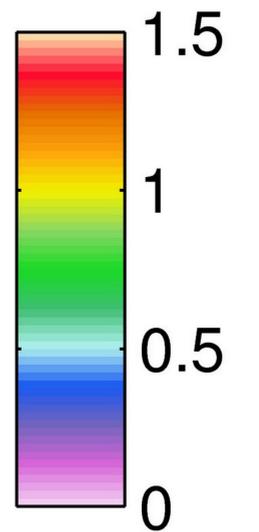
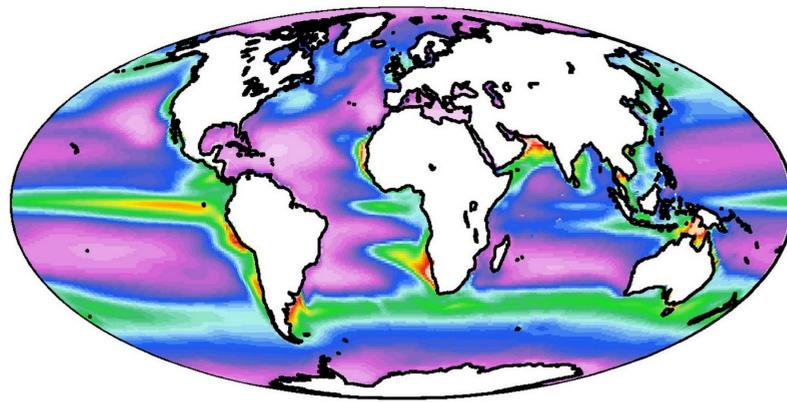
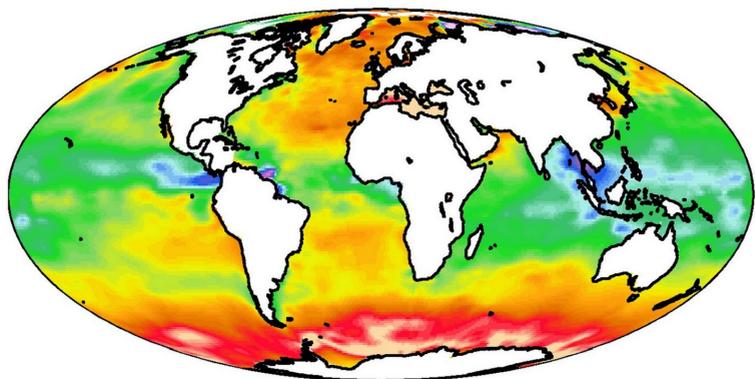


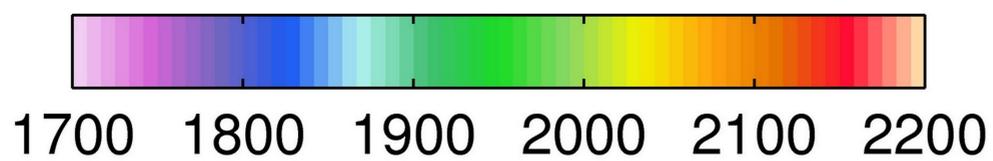
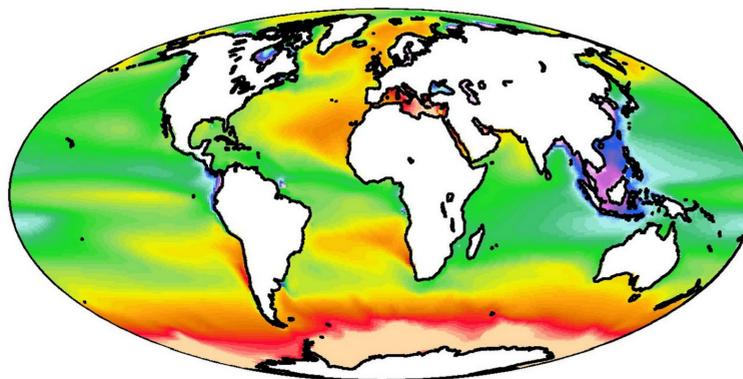
Figure 14.

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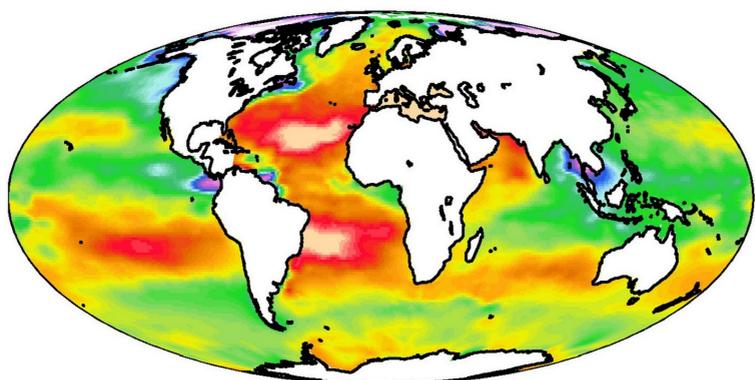
Observed, DIC



Simulated DIC



Observed, alkalinity



Simulated, alkalinity

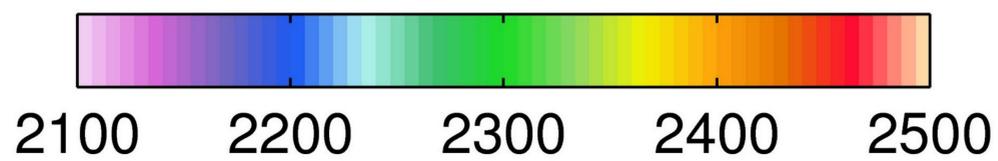
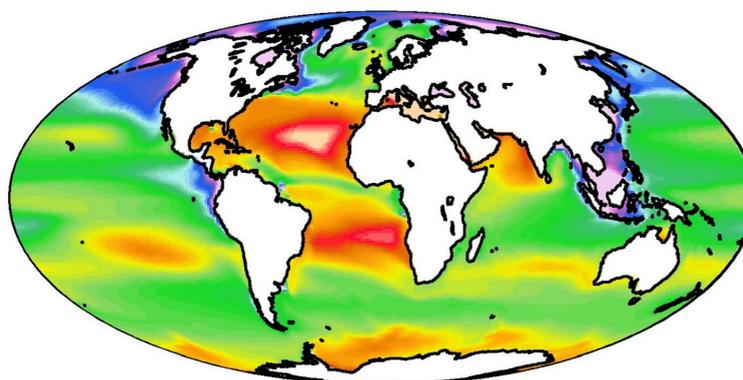
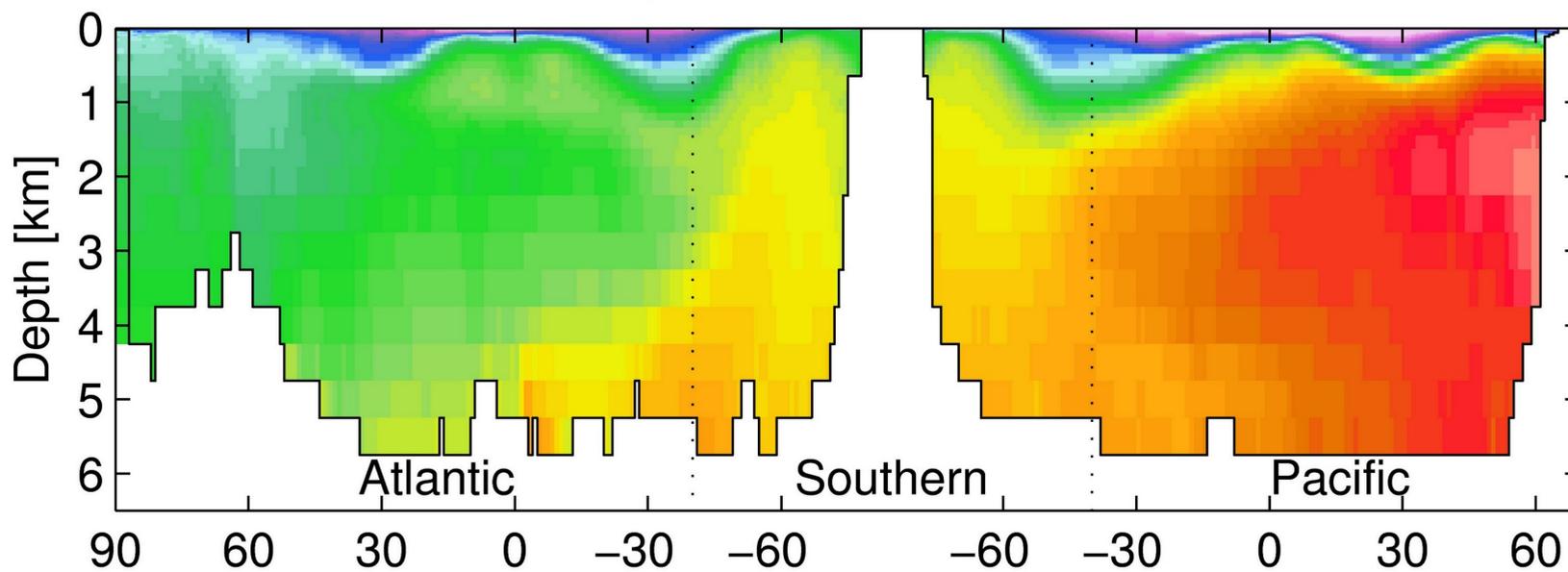


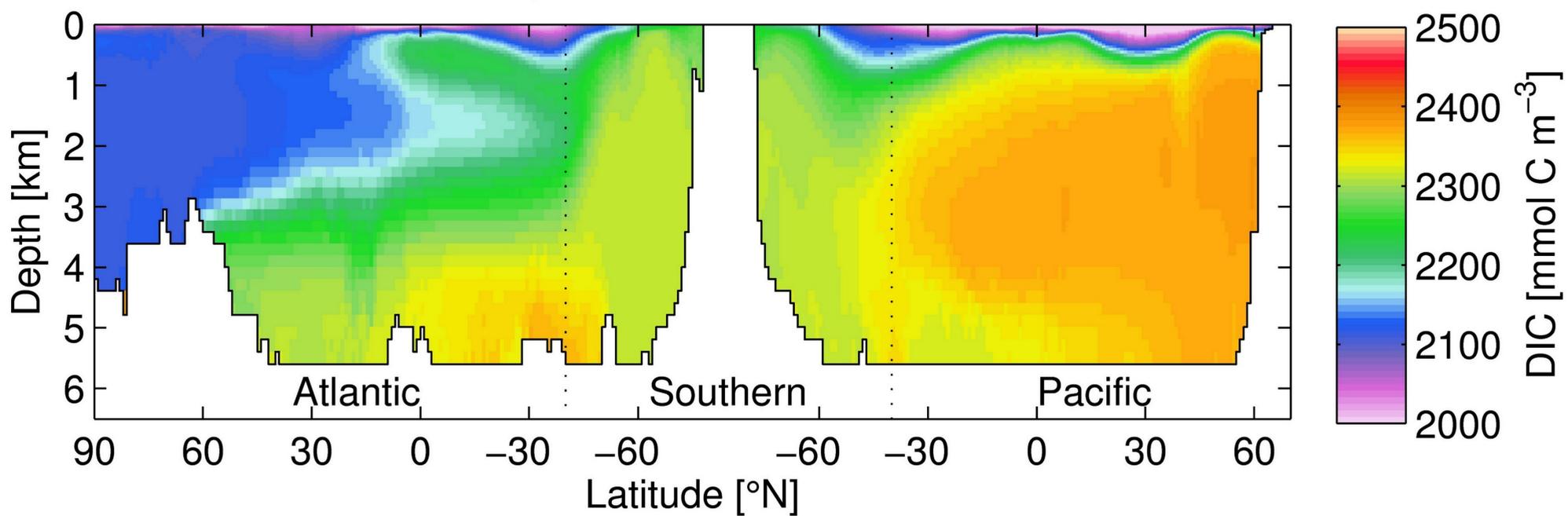
Figure 15.

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Observed, Atlantic–Pacific zonal mean



Simulated, Atlantic–Pacific zonal mean



Difference Atlantic–Pacific zonal mean

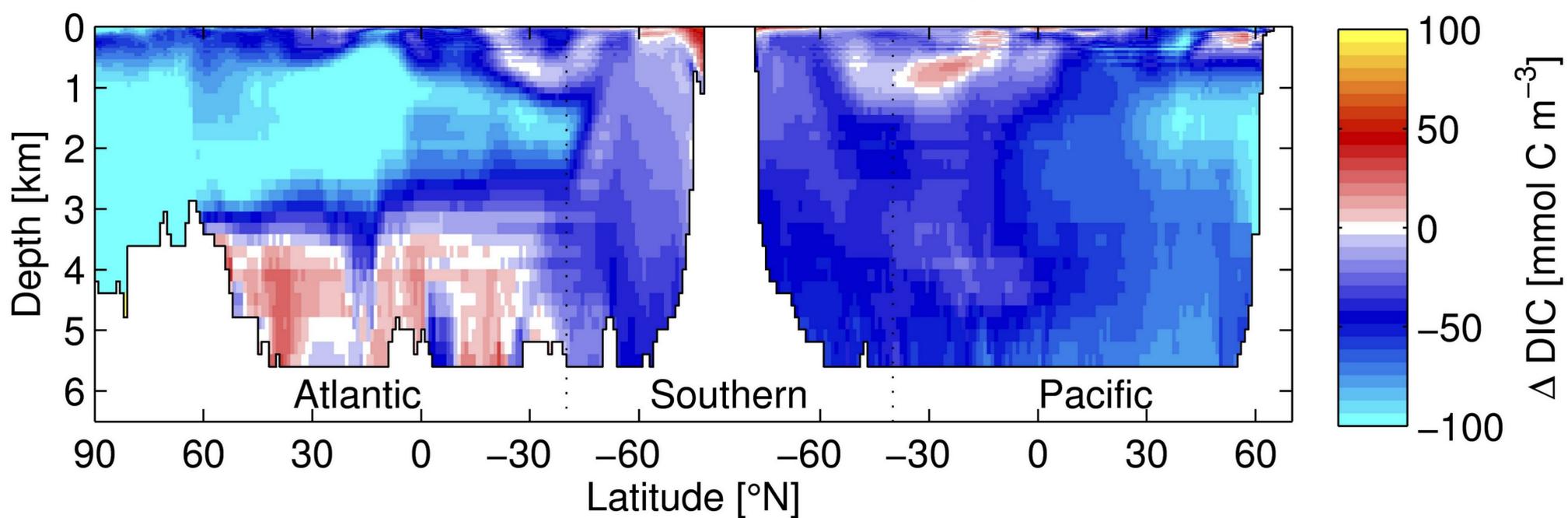
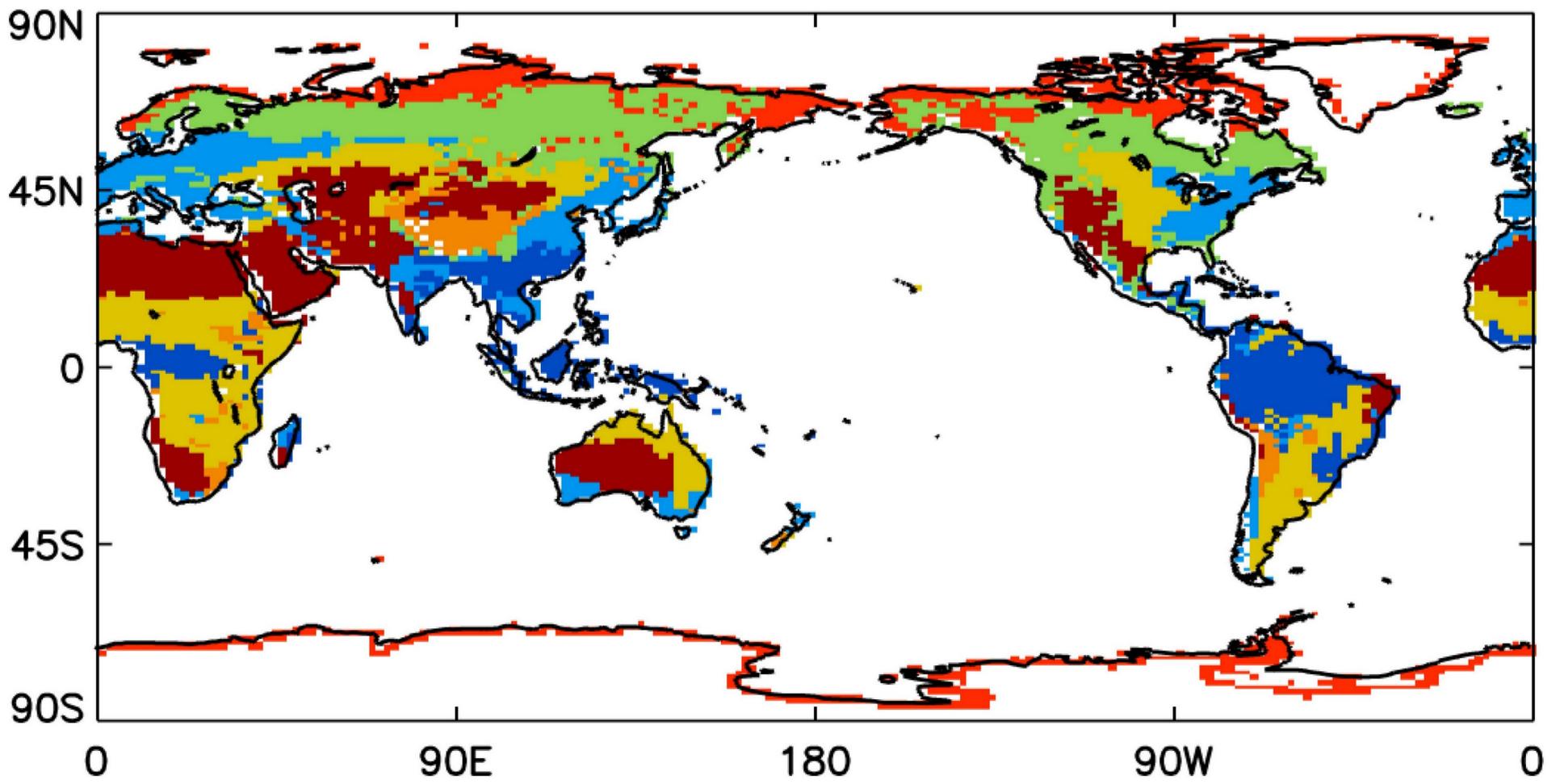


Figure 16.

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Biomes



Tropical Trees Extra Tropical Trees Boreal and Coniferous Savanna Grassland Tundra Desert

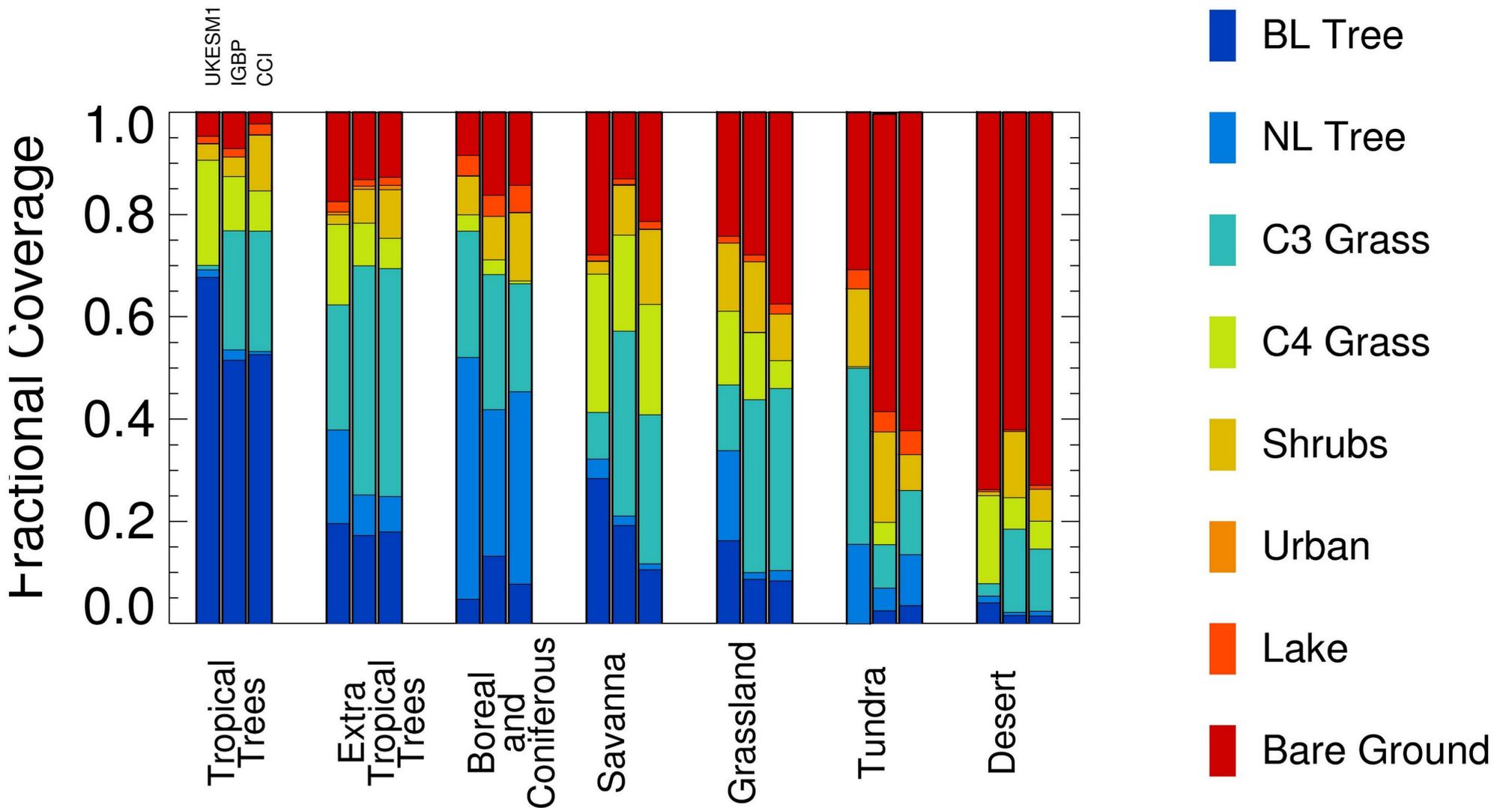


Figure 17.

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