# THE ECOLOGY OF THE TAMAR ESTUARY V. UNDER-WATER ILLUMINATION. REVISION OF DATA FOR RED LIGHT

# By L. H. N. Cooper, D.Sc. and A. Milne, M.A., Ph.D.

From the Plymouth Laboratory

When the Pulfrich photometer is used to measure absorption of light by a turbid estuarine or sea water, the true extinction coefficient for light of a given wave-length

 $\mu_{SW} = \mu_M + \mu_{DW},$ 

where  $\mu_M$  is the extinction coefficient as measured and  $\mu_{DW}$  is the extinction coefficient of the distilled water used in the comparison tube of the photometer. When our paper on under-water illumination in the Tamar Estuary was prepared (Cooper & Milne, 1938), we knew of no modern data on the extinction coefficient of distilled water in the deep red, except those of Lange & Schusterius (1932) which we unjustifiably rejected. James (1938) has since published exhaustive data of his own and gives, in full, results obtained by Prof. J. R. Collins and Dr Ernst Ganz. These leave no doubt that between 700 and 760 m $\mu$  the coefficient rises very steeply from about 0.6 to 2.6, and that the distilled water correction which we applied to our results was in serious error. This spectral region is also one in which visual sensitivity is decreasing rapidly. These two factors much favour the transmission and visual perception of light of shorter wave-length. However, the two red filters, S 72 and S 75, equally strongly favour the transmission of light of longer wavelength. Average extinction coefficients cannot be evaluated by simple inspection. Using the data of James, of Collins and of Ganz we have worked them out for the conditions of our experiments.

# EFFECTIVE EXTINCTION COEFFICIENT OF DISTILLED WATER MEASURED BY THE RED FILTERS OF THE PULFRICH PHOTOMETER

Consider a light source of known energy distribution and let  $J_{\lambda}$  represent the relative energy at a given wave-length. After passing through a filter having a transmission  $c_{\lambda}$ , the relative energy will be reduced to  $c_{\lambda}J_{\lambda}$ . If  $V_{\lambda}$ represents relative visibility (i.e. the relation between luminous sensation as perceived by the eye and radiant energy), the visual sensation produced by this transmitted light will be  $c_{\lambda}J_{\lambda}V_{\lambda}$ . The transmission coefficients of the filters expressed as percentages have been supplied by Messrs Carl Zeiss (cf. Cooper & Milne, 1938, Fig. 1). Relative visibility has been taken from

JOURN. MAR. BIOL. ASSOC. vol. XXIII, 1939

391

26

Smithsonian Physical Tables (1933, 8th revised ed.), Table 346;  $c_{\lambda}J_{\lambda}V_{\lambda}$  may be plotted against wave-length, and the area  $\int_{\lambda_1}^{\lambda_2} c_{\lambda}J_{\lambda}V_{\lambda} d\lambda$  enclosed by the curve and the wave-length axis will provide a measure of the light sensation perceived by the observer's eye. It may be evaluated most conveniently with a planimeter.

Since in our work we used 0.25 m. absorption tubes, we have also calculated the relative energy of a narrow wave band, having extinction coefficient  $\mu_{\lambda}$ , after passing through 0.25 m. of distilled water,  $J_{\lambda}e^{-0.25\mu_{\lambda}}$ . This intensity also may be transmitted by the spectral filter and perceived by the eye so that the measure of light sensation will be  $c_{\lambda}J_{\lambda}e^{-0.25\mu_{\lambda}}V_{\lambda}$ . This may be plotted on the same graph as  $c_{\lambda}J_{\lambda}V_{\lambda}$ . The ratio of the areas suffices to evaluate the logarithmic term in the expression

$$0.25 \,\mu_{\mathcal{A}} = 2.303 \,\log_{10} \frac{\int_{\lambda_1}^{\lambda_2} c_{\lambda} J_{\lambda} V_{\lambda} \,d\lambda}{\int_{\lambda_1}^{\lambda_2} c_{\lambda} J_{\lambda} e^{-0.25 \mu_{\lambda}} V_{\lambda} \,d\lambda}.$$

The value obtained for  $\mu_A$  is the effective extinction coefficient of distilled water for a given filter, which is required as our correction factor. It applies only for 0.25 m. tubes.

Since we do not know the spectral energy distribution of the Nitra lamp supplied with the Pulfrich photometer, we have calculated the effective extinction coefficient, first from Abbot's mean noon sunlight data (Seventh Internat. Congress of Photography, 1929), secondly for light having an equal energy spectrum over the range of wave-length considered, and thirdly for a light source having a colour temperature of 2800° K. (Table I). Considerable latitude in the colour composition of the light source may be tolerated. Much more important are the great variations of the extinction coefficients based on the determinations of different observers and of relative visibility with wave-length. Furthermore, in this region considerable personal idiosyncracy in relative visibility is to be expected.

James (1938, p. 37) discusses the variable coefficients for distilled water found by different observers. These cannot be attributed to errors of measurement and must be due to some variable property of distilled water itself, possibly minute traces of impurities or variable degree of association of the water molecules dependent on the previous history of the water. Whatever the cause, any ordinary laboratory sample of distilled water is likely to show variability at least as great. Below 600 m $\mu$  the coefficient is small compared with the measurements on a fairly turbid estuarine water, so that some degree of uncertainty is of little consequence. Above 650 m $\mu$  matters are very different. The best we can do is to apply an average effective extinction coefficient as correction factor, as suggested in Table I.

392

#### THE ECOLOGY OF THE TAMAR ESTUARY

In relatively transparent ocean water, accurate measurement of the extinction coefficients for blue, green and yellow light requires a water column much longer than can conveniently be accommodated by the Pulfrich photometer. Even if this difficulty could be overcome, the correction terms for the absorp-

### TABLE I. EFFECTIVE EXTINCTION COEFFICIENT FOR RED LIGHT

# Transmitted by 0.25 m. of distilled water and measured with the filters of the Pulfrich photometer

	perature ° C.	Туре с	of spectrum	Particulars of filters		
Filter no. Effective wave-band*, $m\mu$ Centre of gravity of filter, $m\mu$				S 75 730-770 750	S 72 700–760 729	S 61 600–650 619
				Effective extinction coefficient		
Calculated from data of:						
James (1938)	Room	Colour 2800°	temperature K.	2.24	1.74	
55	22	Equal e	nergy	2.23	1.28	0.28
22	22	Mean n	oon sunlight	2.29	1.76	
Collins (cit. James)	0.5	22	>>	2.61	2.02	0.291
22 22	26	22	22	2.79	2.14	
Ganz (cit. James)	12	25	>>	2.54	1.55	
Values adopted as correction terms				2.5	I.8	0.28

\* These are the limits within which the transmission exceeds one-half the maximum. For purposes of summation, transmissions exceeding one-tenth were included.

† Approximate value obtained by inspection (strictly for equal energy spectrum).

tion of light by ordinary laboratory samples of distilled water remain too uncertain for the method there to be of much value. Its utility is confined to fairly turbid inshore or estuarine waters.

#### **REVISION OF THE ORIGINAL DATA**

A list of the numerical revisions made necessary by the change in our correction factors will be found at the end.

The revised values of the ratio  $\mu_{\lambda}/\mu_{470}$  for the S 61, S 72 and S 75 filters still show that very turbid waters transmitted red light as well as or better than blue. When  $\mu_{\lambda}$  is graphed against  $\mu_{470}$ , the relationship may be expressed surprisingly well by the linear equations given in Table II. The constant terms for the red filters are almost the same as those applied to correct for the absorption by distilled water. For the remaining filters similar terms would be quite small. The views expressed in the earlier paper as to the transmission of red light still apply up to about 680 or 700 m $\mu$ . Above this the molecules of water itself manifest strong absorption of light, and in fairly clear waters this is all important, but that part of the extinction coefficient,

26-2

which was due to dissolved and suspended solids, was always least in the deep red. These factors work in opposition, and as a result the minimum extinction coefficient or maximum transmission for our more turbid waters was to be found around  $600 \text{ m}\mu$ .

The ratios of the wave-length at the centre of gravity of the blue S47 filter compared with the corresponding wave-lengths of the other filters, also given in Table II, show that that part of the extinction coefficient due to

#### TABLE II. RELATION BETWEEN EXTINCTION COEFFICIENTS AND WAVE-LENGTH

#### (Revision of original Table III)

Filter	Equation connecting $\mu$ and $\mu_{463}$	Ratio of wave- length of blue filter to that of the others (at centres of gravity) $\frac{4^{63}}{\lambda}$
S 43 S 47	$\mu_{434} = \mathbf{I} \cdot \mathbf{I} \mathbf{I} \ \mu_{4b3} \\ \mu_{462} = \mathbf{I} \cdot \mathbf{OO} \ \mu_{462}$	1.07 1.00
S 50 S 53	$\mu_{494} = 0.92 \ \mu_{463}$ $\mu_{590} = 0.86 \ \mu_{163}$	0.94
S 57	$\mu_{572} = 0.82 \ \mu_{463}$	0.81
S 61	$\mu_{619} = 0.35 + 0.73 \mu_{463}$	0.75
S 72	$\mu_{729} = 1.80 + 0.68 \mu_{463}$	0.64
\$ 75	$\mu_{750} = 2.54 + 0.64 \mu_{463}$	0.62

suspensoids and possibly to matter in solution is inversely proportional to the first power of the wave-length. The relationship is sufficiently close that, if only one extinction coefficient were known, that for any other wave-length could be calculated with an accuracy sufficient for most practical purposes.

The absorption index, K, defined by the equation

$$\mu_{SW} - \mu_{DW} = \frac{4\pi K}{\lambda},$$

is a constant for each water;  $\mu_{SW} - \mu_{DW}$  is, in effect, a measure of absorption by suspensoids.

The dependence of scattering of light by particles in sea water upon various powers of the wave-length has been discussed by Kalle (1938).

#### REVISED DATA (COOPER & MILNE, 1938)

P. 510, Table I: Extinction coefficients of distilled water: Accepted values for filters S 61, S 72 and S 75 are 0.28, 1.8 and 2.5 respectively.

P. 514, l. 26: For " $(\mu_{500} = 5.0; \mu_{720} = 3.8)$ " read " $(\mu_{500} = 5.0; \mu_{720} = 5.3)$ ". P. 514, Table II: For last line of figures, read:

 $\mu_{720}$  5.9 8.2 7.5

394

P. 515, ll. 1 and 2 from bottom: *Read*: "the intensity of red and blue light entering the water was cut down to one-thousandth within 0.5 and 0.6 m. respectively ( $\mu = 13.5$  and 16.5)."

P. 516, Fig. 3: The plots at 610, 720 and 750 m $\mu$  should be 0.1, 1.3 and 2.0 units higher respectively. A minimum extinction coefficient occurs at about 600 m $\mu$  in most cases.

P. 517, Table III: See p. 394.

Pp. 519 and 520, Figs. 4 and 5: All curves should show less relative energy in the dark red. The statement about preferential transmission of red light does not apply above about 680 m $\mu$ .

P. 522, Table IV.

# Depth in Metres at which Intensity of Red Light is Reduced to 1 % of that Incident on the Surface

# Assuming 15 % Surface Loss; $\mu d_1 = 4.44$

Centre of gravity of filter	Breakwater	Drake's		Hamoaze		Neal Point
mμ	No. I Buoy	No. I Buoy	No. I Buoy	No. 7 Buoy	No. 15 Buoy	Buoy
	Fe	ebruary 18 19	37. Spring tie	de. High wat	er	
619	2.35	1.97	1.85	1.41	1.30	
729	1.37	1.26	1.22	1.02	0.99	
750	1.03	1.00	1.04	0.86	0.82	
	F	ebruary 16 19	37. Spring ti	de. Low wat	er	
619	1.93	1.95	1.10	0.78	0.23	0.35
729	1.33	1.25	0.89	0.69	0.48	0.33
750	1.04	1.02	0.76	0.60	0.46	0.32
		June 17 193	7. Neap tide.	High water		
619	3.86	4.27	3.06	2.48	2.54	2.48
729	1.77	1.78	1.60	1.42	1.40	1.39
750	1.38	1.40	1.27	1.18	1.19	1.12
		June 23 1937	7. Spring tide	. Low water		
619	5.00	3.21	2.74	2.04	1.41	0.98
729	1.95	1.60	1.48	1.27	1.00	0.79
750	1.49	1.27	1.16	1.02	0.90	0.71

# P. 525, Table V.

#### DAYLIGHT FACTORS AT BOUNDARIES OF ALGAL ZONES ON BUOYS

No correction for surface loss

	Zone I		Zone II	
	Plymouth Sound	Hamoaze No. 7 Buoy	Plymouth Sound	Hamoaze No. 7 Buoy
Depth of bottom of zone, m Daylight factor, %:	0.18	0.092	0.326	0.223
June 17 (high water) June 23 (low water)	78·3 76·4	80·0 77·8	62·3	61·8
February 18 (high water) February 16 (low water)	65·4 60·9	69·8 52·4	43·7 37·7	45·0 23·3

# TABLE VI. DAYLIGHT FACTORS AT BOTTOM OF TWO BUOYS (WITHIN ZONE III)

Mallard Buoy (Plymouth Sound) Hamoaze No. 7 Buoy

	Depth of bottom	Daylight	Wave-length of maximum	Depth of bottom	Daylight	Wave-length of maximum
Date and state of tide	m.	%	mµ	m.	%	mµ
June 17 (high water)	2.19	7.41	570	2.47	0.97	570
June 23 (low water)	2.19	5.21	570	2.47	0.35	600
February 18 (high water)	2.19	0.85	580	2.47	0.020	600

620

2.47

0.000025

620

P. 524, Table VII. All the extinction coefficients given in the last column should be increased by 1.3.

0.34

2.19

P. 526, l. 3: For "0.36 and 17 m. (ratio 1:48)" read "0.33 and 17 m. (ratio 1:52)."

P. 526, Table in Summary: The values for red light (610-640 and 710-760 m $\mu$ ) should be increased by 0.1 and 1.33 respectively.

#### SUMMARY

New values for the extinction coefficient of red light in distilled water require that published results for this coefficient in waters from the Tamar Estuary shall be increased by 0.1, 1.3 and 2.0 at 610, 720 and 750 mµ respectively. Earlier deductions as to the favourable transmission of red light in turbid estuarine waters apply only up to about 680 m $\mu$ . That part of the extinction coefficient due to suspensoids and colouring matter was inversely proportional to the first power of the wave-length. Other conclusions remain unaffected.

#### REFERENCES

COOPER, L. H. N. & MILNE, A., 1938. The ecology of the Tamar Estuary. II. Underwater illumination. Fourn. Mar. Biol. Assoc., Vol. XXII, pp. 509-28.

JAMES, H. R., with BIRGE, E. A., 1938. A laboratory study of the absorption of light by lake waters. Trans. Wisconsin Acad. Sci., Vol. xxxI, Chapter II.

KALLE, K., 1938. Zum Problem der Meereswasserfarbe. Ann. Hydrogr. Marit. Meteorol., pp. 1-13.

LANGE, B. & SCHUSTERIUS, C., 1932. Die Absorption des Wassers im sichtbaren Spektralgebiet. Zeits. Physikal. Chem., Abt. A, Bd. 159, pp. 303-5; Berichtigung, ibid., Bd. 160, Heft 6.

February 16 (low water)